

# D-meson propagation in hadronic matter and consequences on heavy-flavor observables in heavy-ion collisions

**Vitalii Ozvenchuk,**

in collaboration with

**J.Aichelin, P.B.Gossiaux, J.M.Torres-Rincon**



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Hersonissos, Crete, Greece**



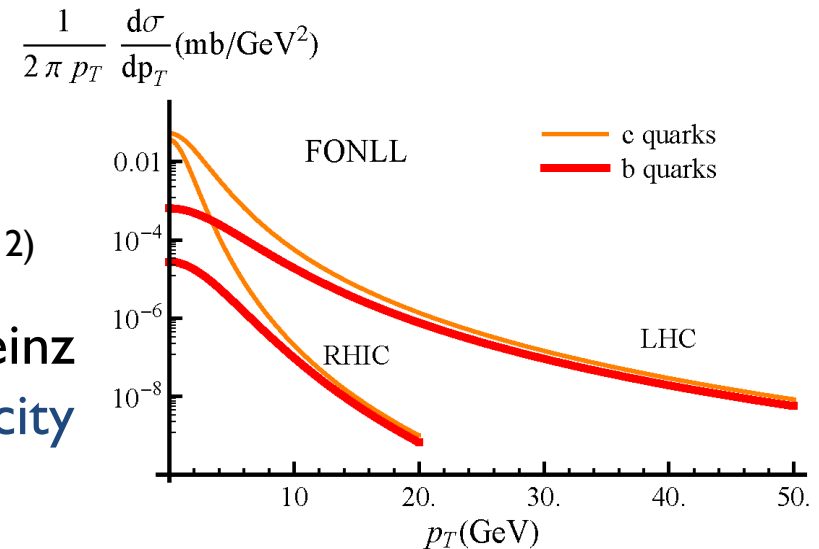
# „Monte Carlo @ Heavy Quark“ generator

- production of heavy quarks at the original NN scattering points according to the FONLL spectra

M.Cacciari et al., Phys. Rev. Lett. **95** (2005), JHEP **1210** (2012)

- bulk evolution: non-viscous Kolb-Heinz hydro; provides temperature and velocity fields

P.F.Kolb, J.Sollfrank, U.Heinz, Phys. Rev. **C62**, 054909 (2000)



- evolution of HQ in the bulk: the Boltzmann equation
- hadronization of HQ: coalescence (low  $p_T$ ) and fragmentation (high  $p_T$ )

$$T_c = 165 \text{ MeV}, \quad \varepsilon_c = 0.45 \text{ GeV/fm}^3$$

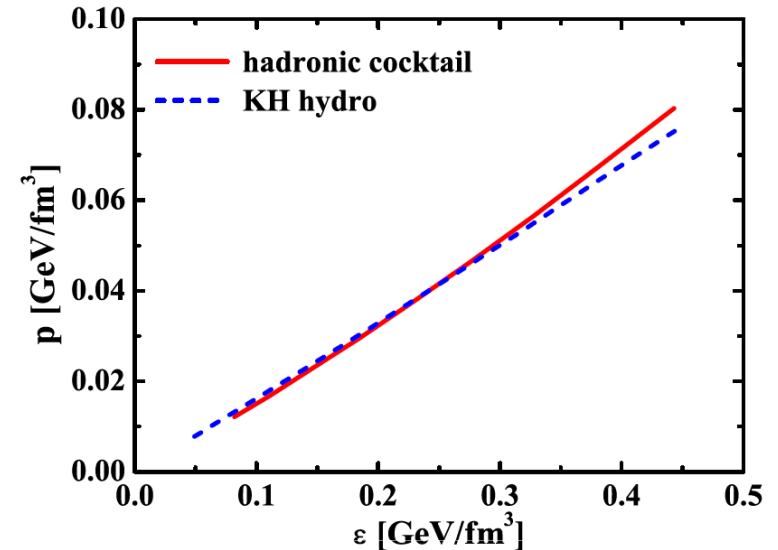
- D-meson propagation in hadronic matter: the Fokker-Planck equation

$$\frac{\partial f(\mathbf{p}, t)}{\partial t} = \frac{\partial}{\partial p_i} \left[ A_i(\mathbf{p}) f(\mathbf{p}, t) + \frac{\partial}{\partial p_j} B_{ij}(\mathbf{p}) f(\mathbf{p}, t) \right]$$

# Hadronic cocktail

## □ Hadron gas **composition:**

- light mesons (up to masses 1.285 GeV)
- strange mesons (K, K\*, K<sub>1</sub>)
- nucleons
- nuclear and  $\Delta$ -resonances (up to masses 1.7 GeV)



## Thermal equilibrium + effective chemical potentials

□ Employ a **specific entropy** of  $S/N_B = 250$  (characteristic value for collisions at top RHIC energy)  
R.Rapp, Phys. Rev. **C66**, 017901 (2002)

□ **Freeze-out** point:  $T_{fo}^{ch} = 170$  MeV,  $\mu_B^{ch} = 28.3$  MeV  
 $\epsilon \approx 0.45$  GeV/fm<sup>3</sup>

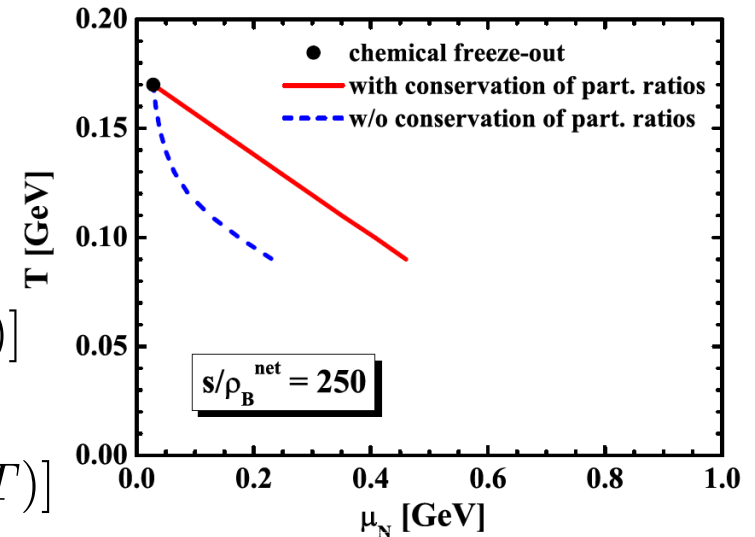
# Thermodynamic trajectories

□ **Thermodynamic trajectory** keeping a specific entropy fixed:

$$s/\rho_B^{\text{net}} = 250$$

$$s = \mp \sum_i d_i \int \frac{d^3k}{(2\pi)^3} [\pm f \ln f + (1 \mp f) \ln (1 \mp f)]$$

$$\rho_B^{\text{net}} = \sum_{B_i} d_{B_i} \int \frac{d^3k}{(2\pi)^3} [f^{B_i}(\mu_{B_i}, T) - f^{\bar{B}_i}(\mu_{\bar{B}_i}, T)]$$



□ Keep a **ratios** of effective stable particle numbers to effective antibaryon number **constant** in a hadronic evolution:

R.Rapp, Phys.Rev. **C66**, 017901 (2002)

$$\frac{N_B^{\text{eff}}}{N_{\bar{B}}^{\text{eff}}}, \frac{N_\pi^{\text{eff}}}{N_{\bar{B}}^{\text{eff}}}, \frac{N_\eta^{\text{eff}}}{N_{\bar{B}}^{\text{eff}}}, \frac{N_K^{\text{eff}}}{N_{\bar{B}}^{\text{eff}}}, \frac{N_\omega^{\text{eff}}}{N_{\bar{B}}^{\text{eff}}}, \frac{N_{\eta'}^{\text{eff}}}{N_{\bar{B}}^{\text{eff}}}, \frac{N_\phi^{\text{eff}}}{N_{\bar{B}}^{\text{eff}}}$$

$$N_{\bar{B}}^{\text{eff}} = V_{FB} \sum_{\bar{B}_i} n_{\bar{B}_i}(T, \mu_{\bar{B}_i})$$

$$N_\pi^{\text{eff}} = V_{FB} \sum_i N_\pi^{(i)} n_i(T, \mu_i)$$

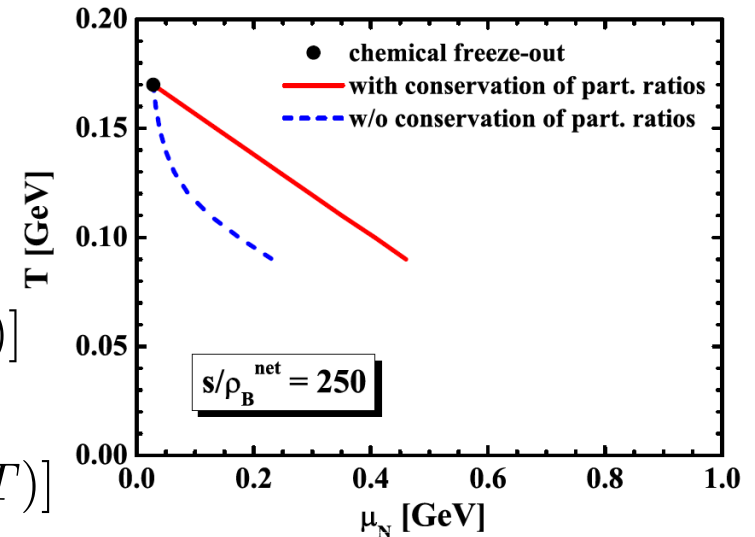
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$$N_\pi^{(\rho)} = 2, N_\pi^{(\Delta)} = 1, N_\pi^{(N(1520))} = 0.55 * 1 + 0.45 * 2$$

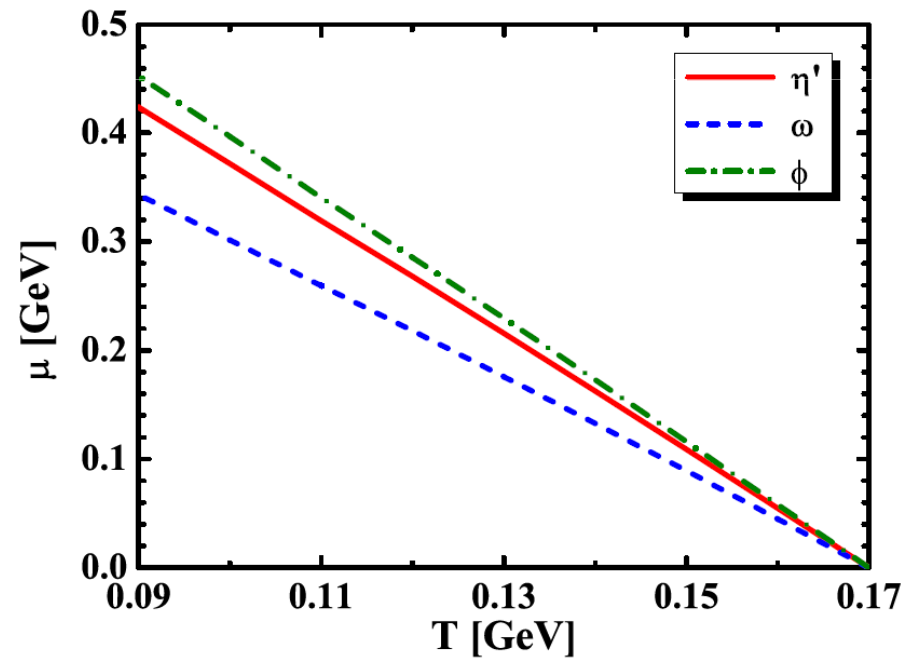
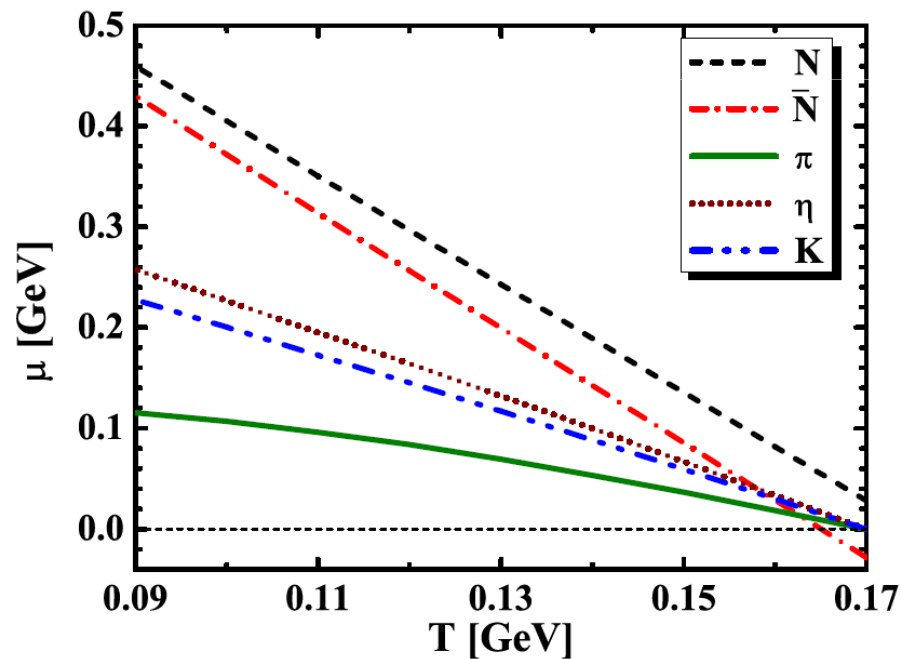
$$\Rightarrow \mu_\rho = 2\mu_\pi, \mu_\Delta = \mu_N + \mu_\pi, \mu_{N(1520)} = \mu_N + 1.45\mu_\pi$$

$$N_\pi^{\text{eff}} = V_{FB} \sum_i N_\pi^{(i)} n_i(T, \mu_i)$$

# Effective chemical potentials

□ To conserve the ratio of effective baryon to antibaryon number we introduce **antibaryon effective ch. potential**,  $\mu_{\bar{B}}^{\text{eff}}$ , e.g.,  $\mu_{\bar{N}} = -\mu_N + \mu_B^{\text{eff}}$ .

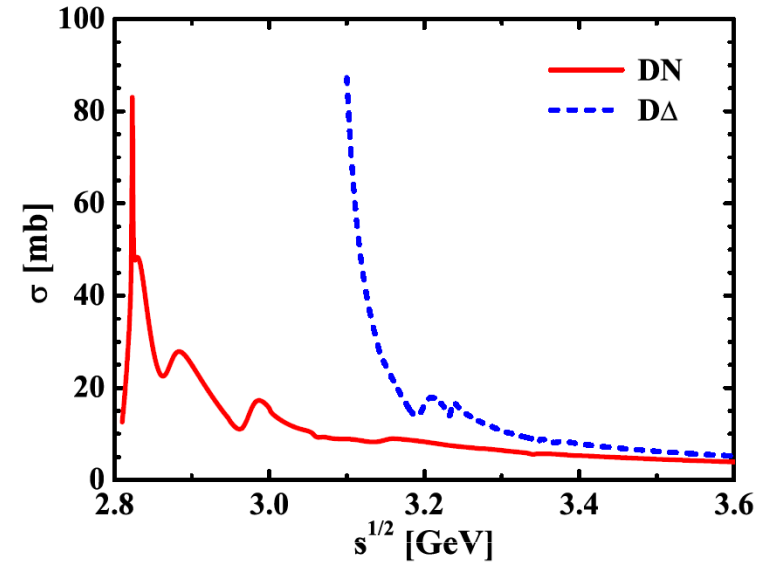
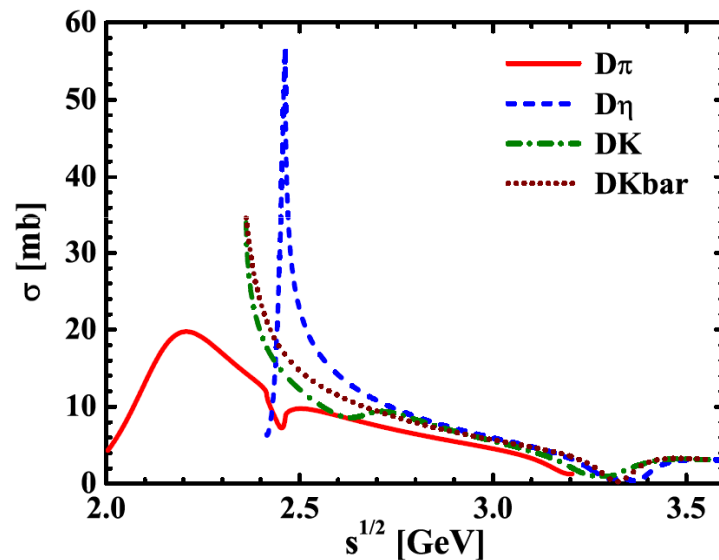
□ At chemical freeze-out temperature all meson effective chemical potentials are **zero**



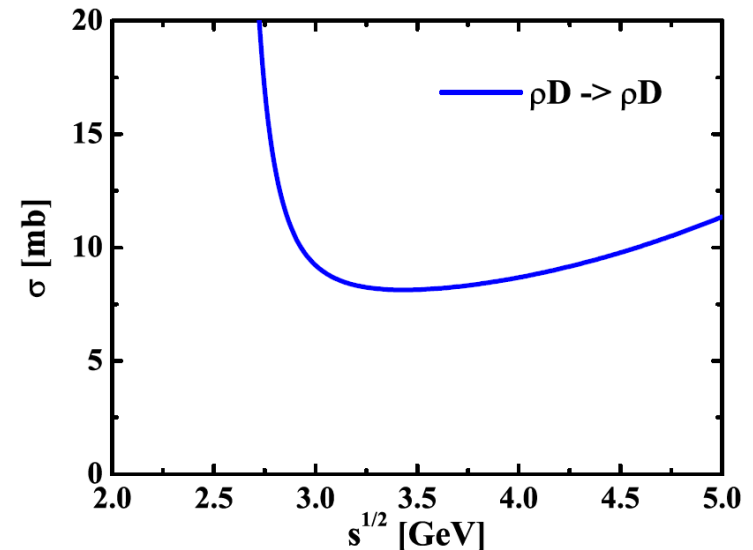
# Elastic cross sections (scenario A)

- Implement the **cross sections** (as in the vacuum) for the interaction of a D-meson with hadrons (effective models):

L.Tolos, J.M.Torres-Rincon, Phys. Rev. **D88**, 074019 (2013)



Z.Lin, T.G.Di, C.M.Ko, Nucl. Phys. **A689**, 965 (2001)



- Other **elastic processes**:

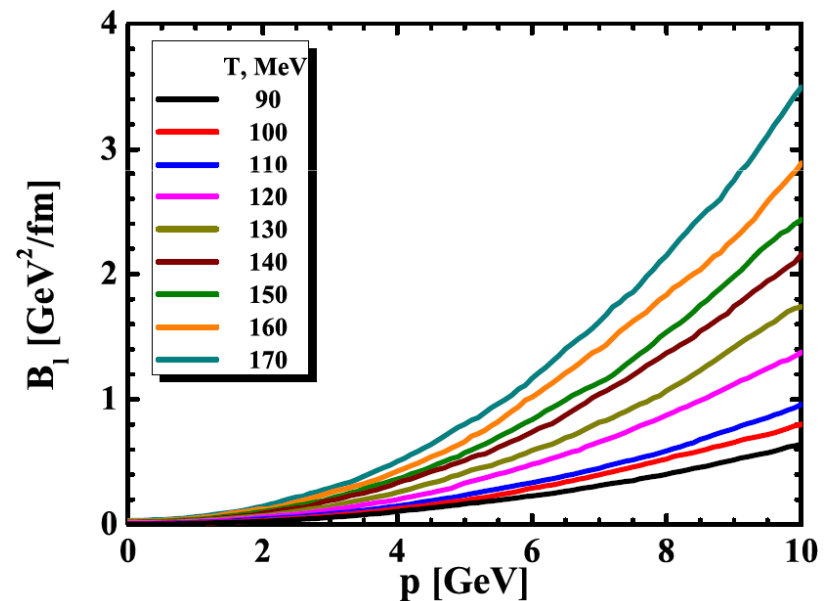
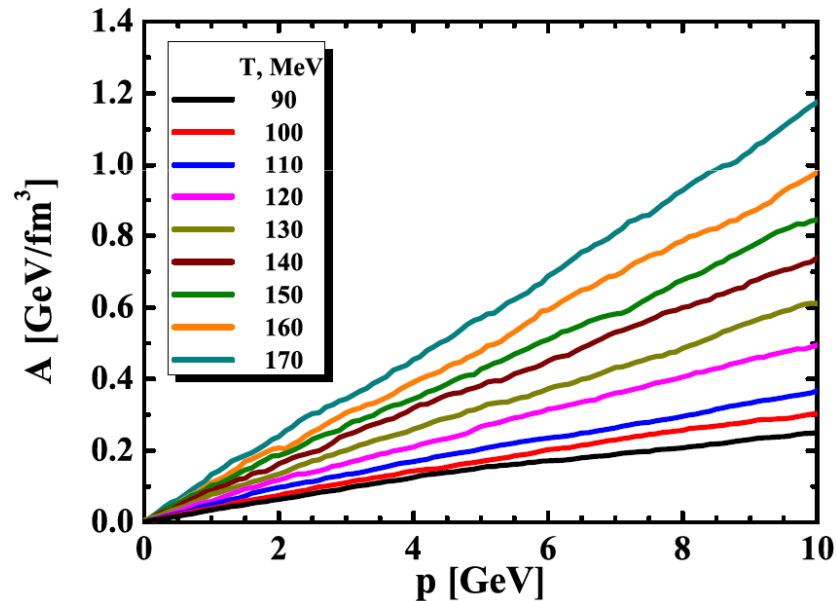
$$Dm \rightarrow Dm \Rightarrow \sigma = 10 \text{ mb}$$

$$DB(\bar{B}) \rightarrow DB(\bar{B}) \Rightarrow \sigma = 15 \text{ mb}$$

# D-meson transport coefficients

- Calculate the following average quantities, which can be related to the **drag, longitudinal and transverse diffusion coefficients**:

$$A = -\left\langle \frac{dp_z}{dt} \right\rangle, \quad B_l = \frac{1}{2} \frac{d(\langle p_z^2 \rangle - \langle p_z \rangle^2)}{dt}, \quad B_T = \frac{1}{4} \left\langle \frac{dp_T^2}{dt} \right\rangle$$



- almost **linear rise** with the momentum;
- contributions from **heavier hadrons** become **important** at higher temperatures

in the **static limit**:

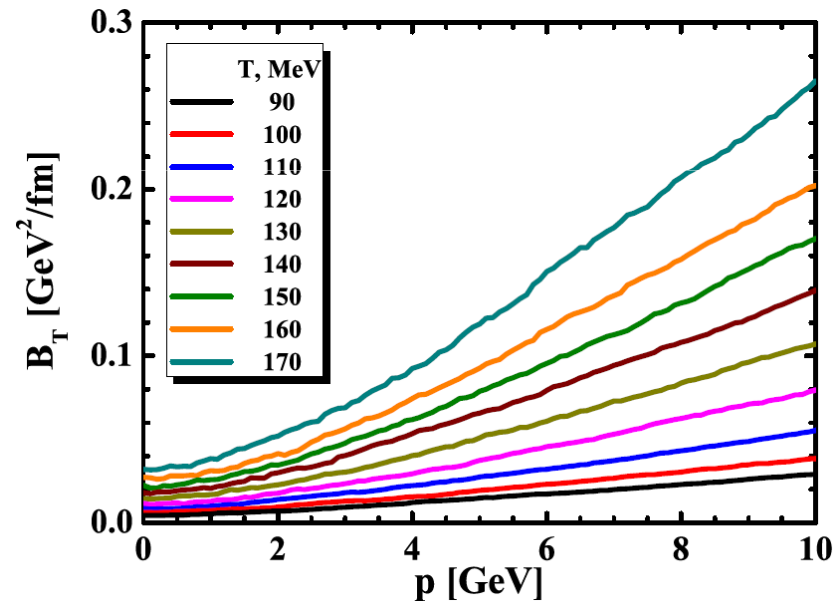
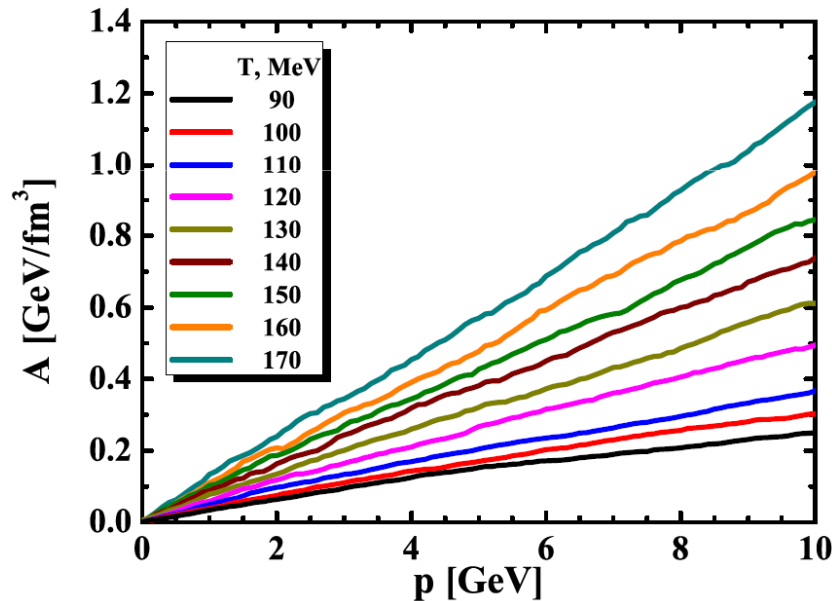
$$\lim_{p \rightarrow 0} [B_l(p) - B_T(p)] = 0$$



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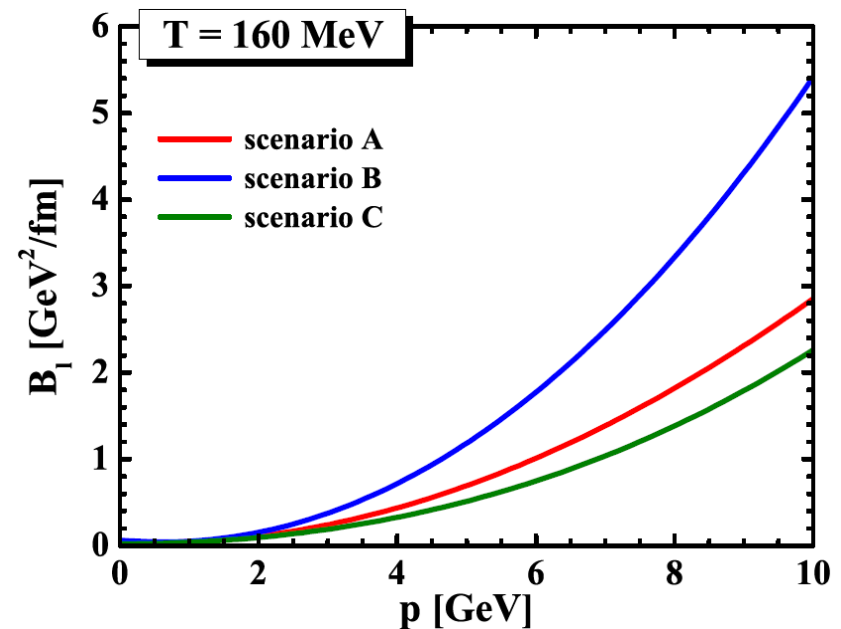
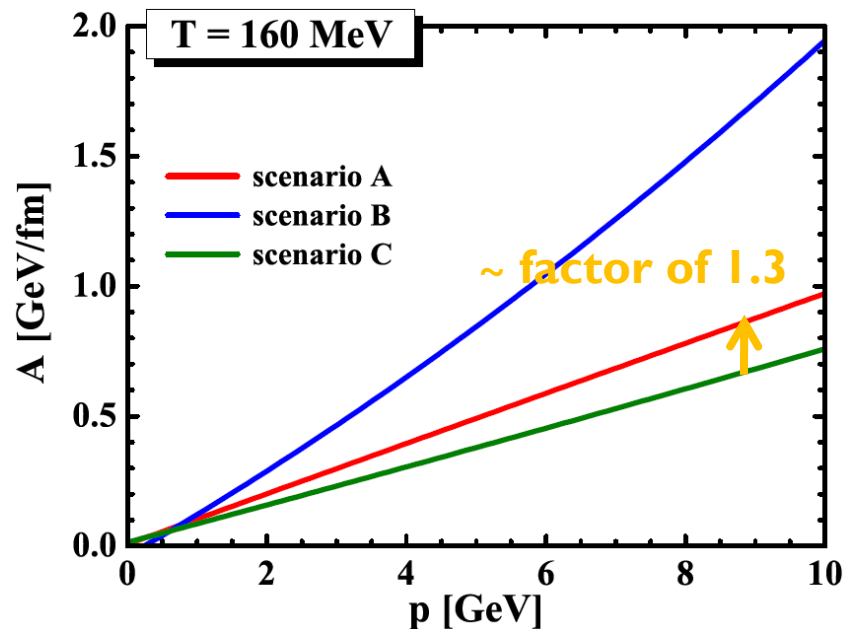
# Other scenarios

❑ **Scenario A** – discussed above...

❑ **Scenario B** – use constant cross sections + effective chemical potentials:

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❑ **Scenario C** – do not implement effective chemical potentials:  $\mu_i^{\text{eff}}(T) = 0$



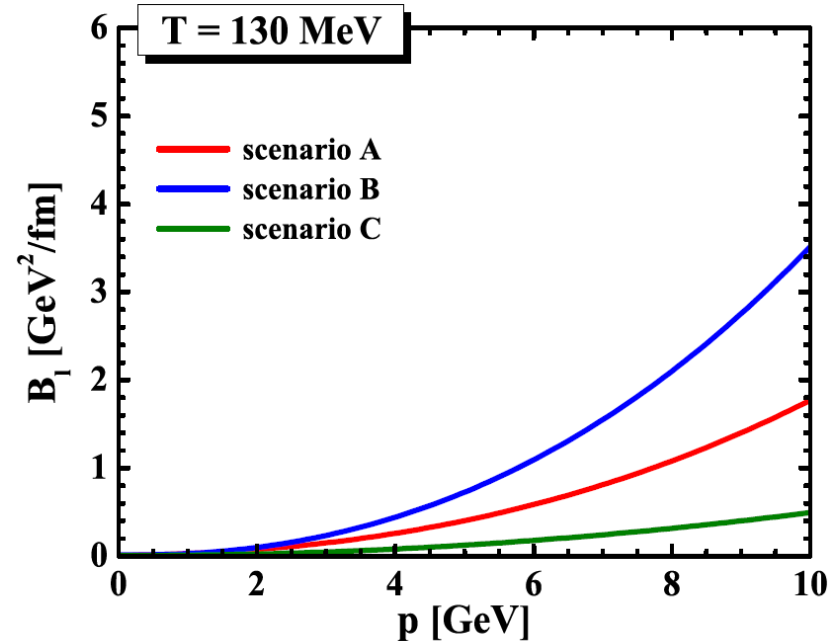
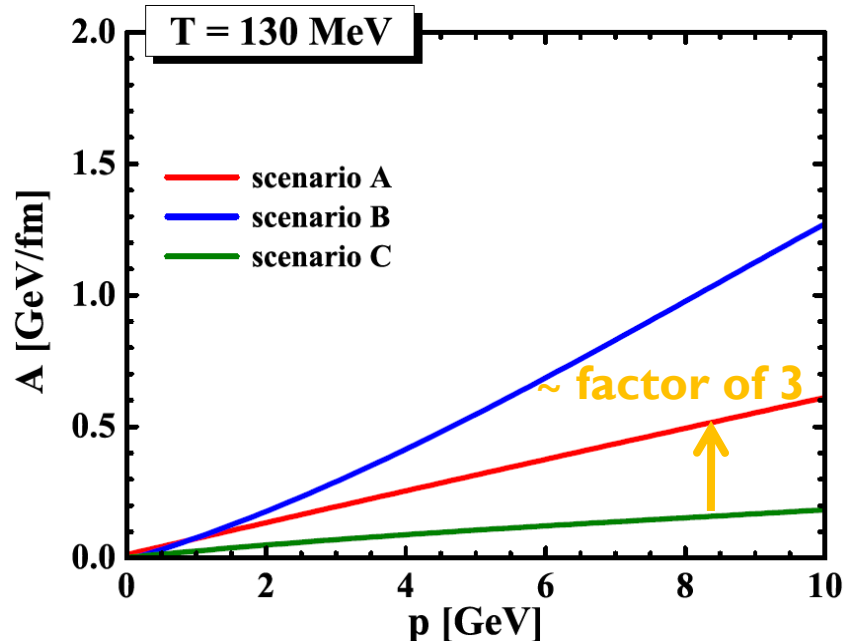
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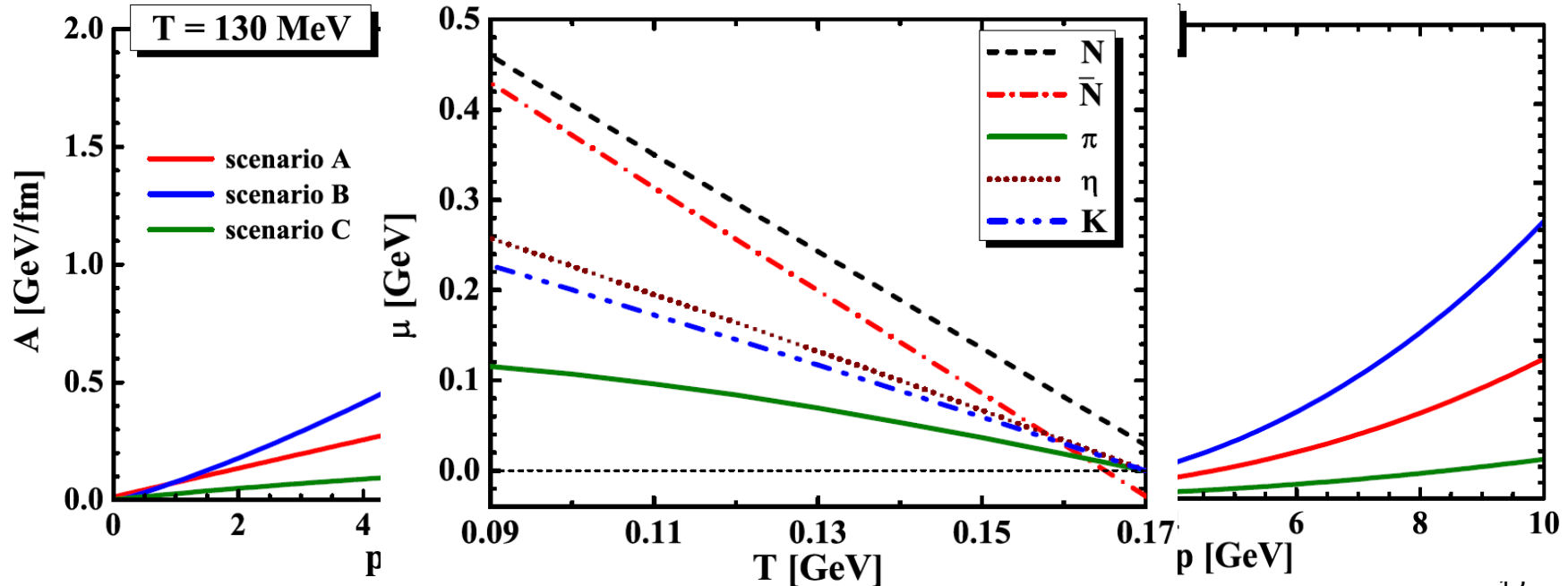
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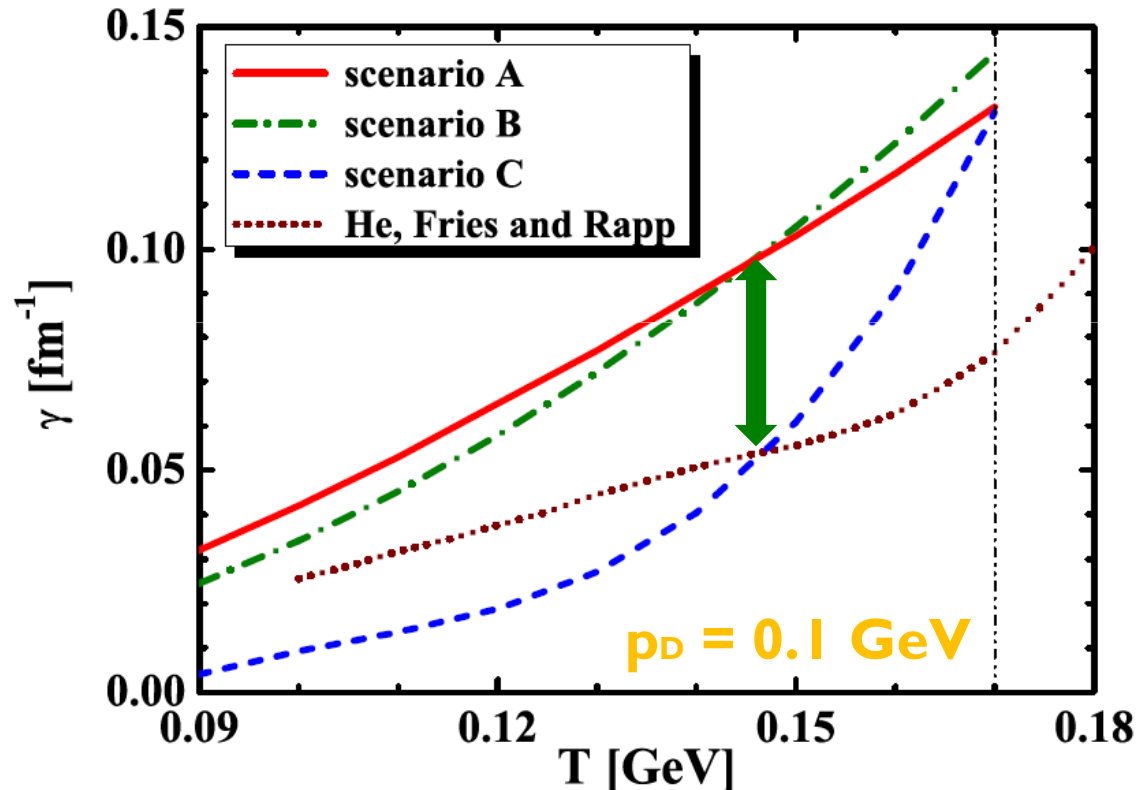


# D-meson thermal relaxation time

□ Evaluate the **D-meson thermal relaxation time**:

M.He, J.Fries, R.Rapp, Phys. Lett **B701**, 445 (2011)

$$\gamma = \frac{A}{p_D} \frac{E_D}{m_D}$$



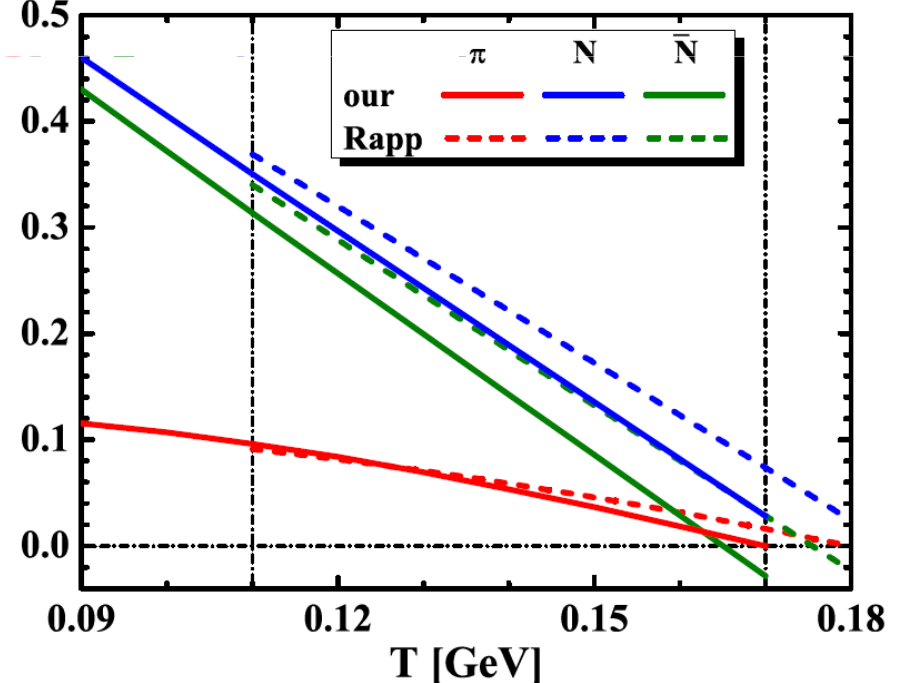
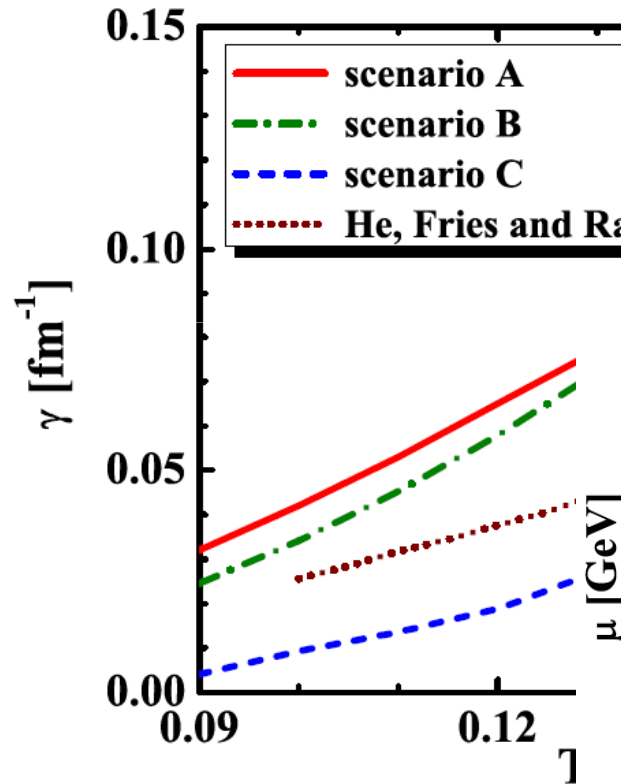
□ **Increase by a factor of 2** of the thermal relaxation time due to the different **hadronic cocktail** and different **cross sections**

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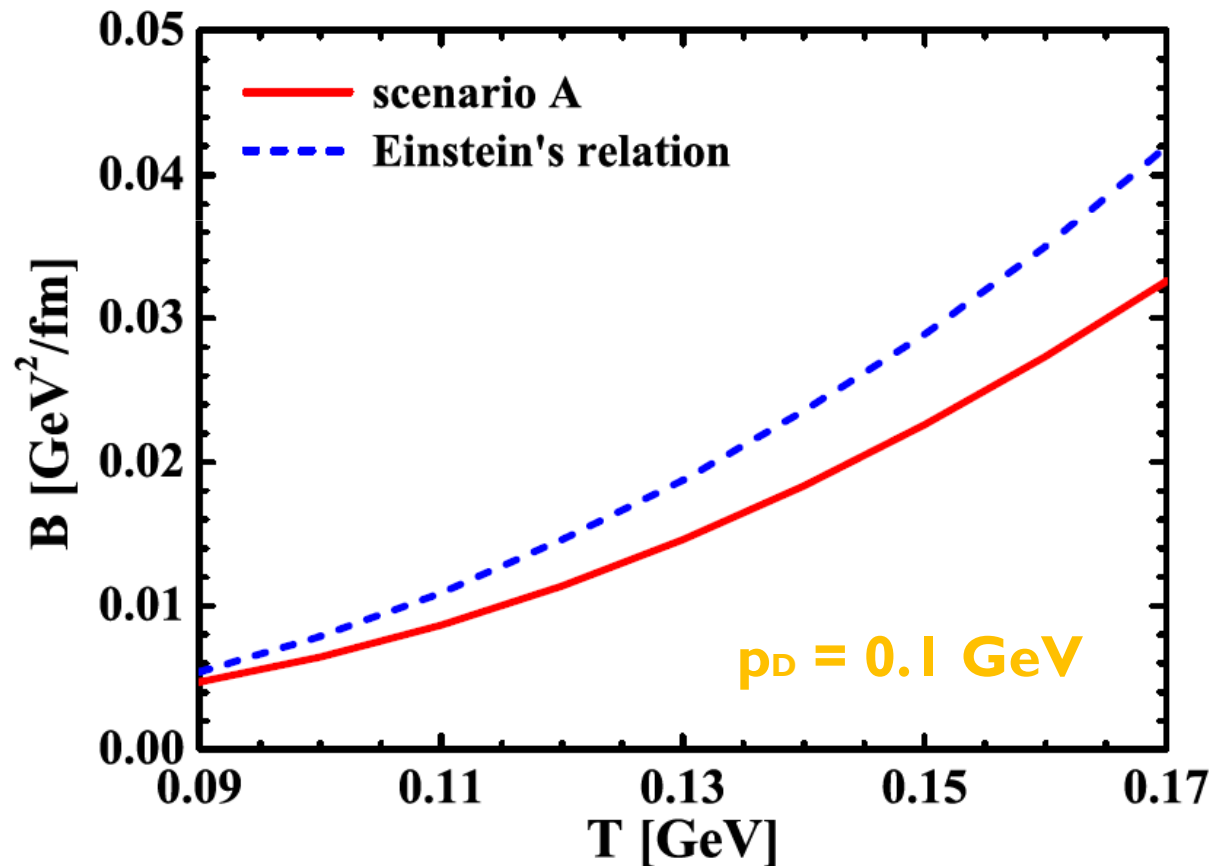


□ Increase by a factor of 2 of the different hadronic cocktail and different

# Einstein relation

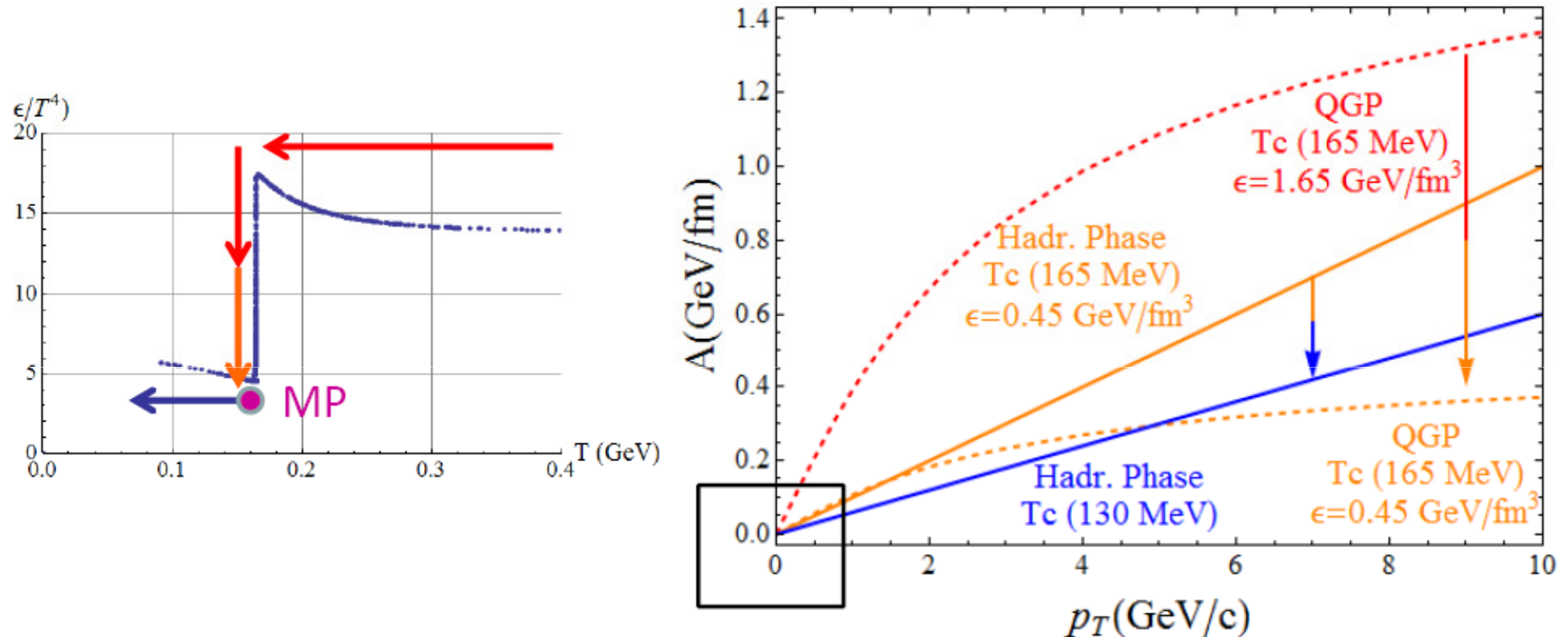
□ In the **static limit**:  $\lim_{p \rightarrow 0} [B_l(p) - B_T(p)] = 0 \Rightarrow B = B_l = B_T$

□ **Einstein relation**:  $B = \gamma m_D T$



# Comparison to HQ in plasma

Slide from Pol's talk

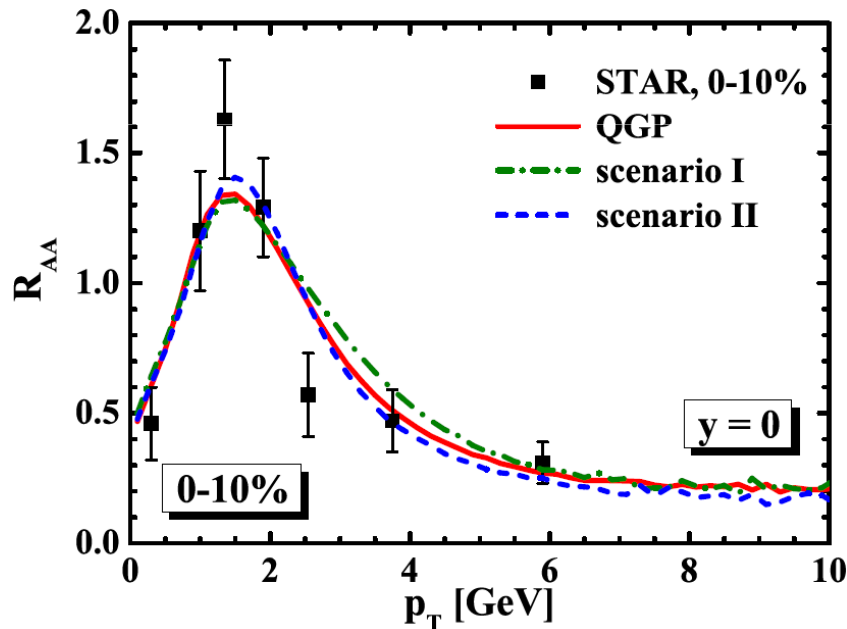


- ❑ **Relaxation times** at the **MP** fairly agrees ( $\approx 10$  fm/c) – it satisfies the “crossover” constrain
- ❑  **$p_T$  dependences** disagree (isotropic cross sections in the HG)

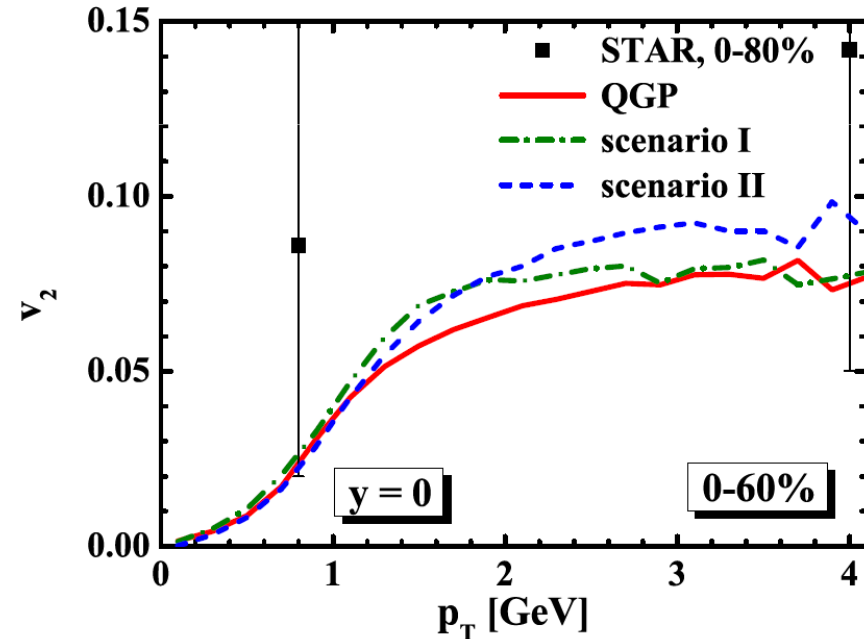


# $R_{AA}$ and $v_2$ of D-meson

- Implement the **obtained results** to “MC@HQ” generator
- Calculate the D-meson **nuclear modification factor and elliptic flow** for two different scenarios:
  - **scenario I**: transport coefficients, drag and diffusion, directly from the simulation
  - **scenario II**: drag – simulation, diffusion – Einstein relation



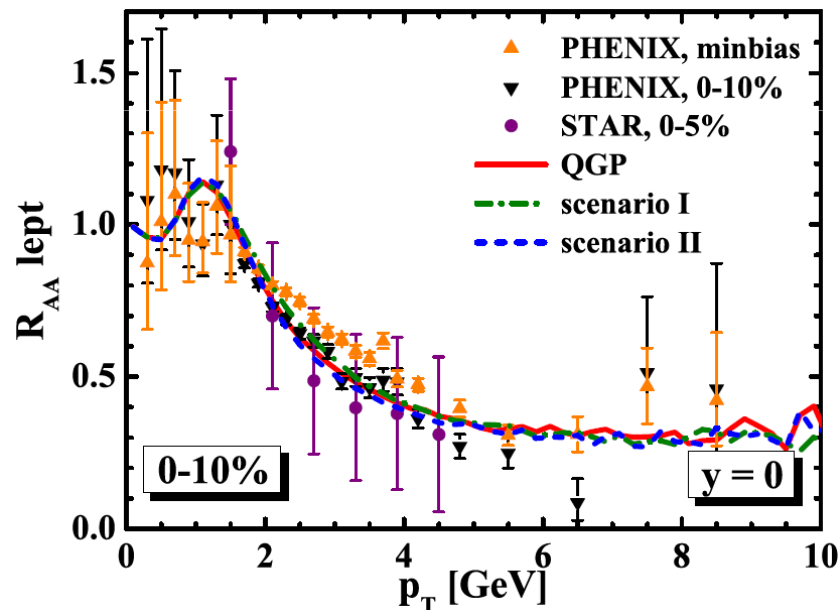
**Almost invisible for  $R_{AA}$**



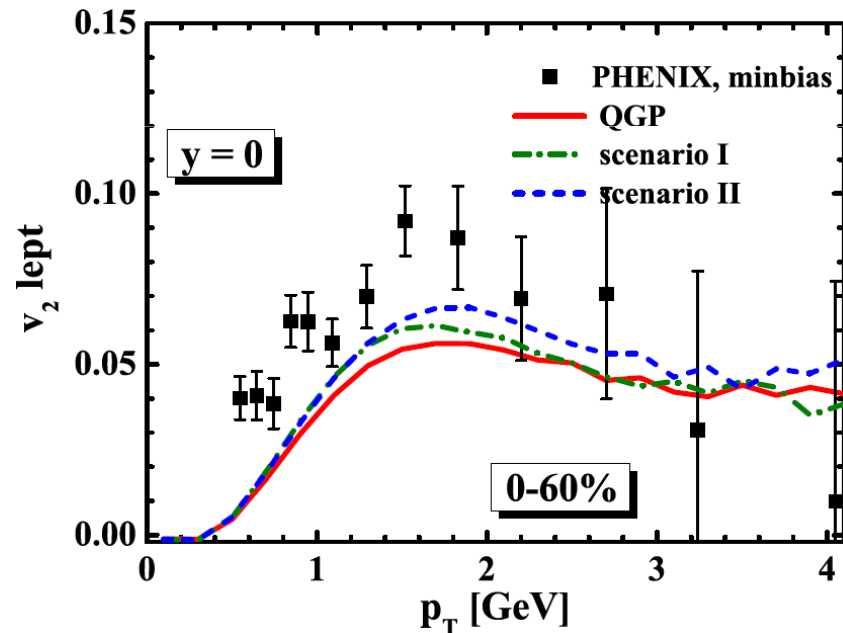
**Moderate effect on  $v_2$ ,  
but systematic**

# $R_{AA}$ and $v_2$ of single nonphotonic leptons

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  - **scenario I**: transport coefficients, drag and diffusion, directly from the simulation
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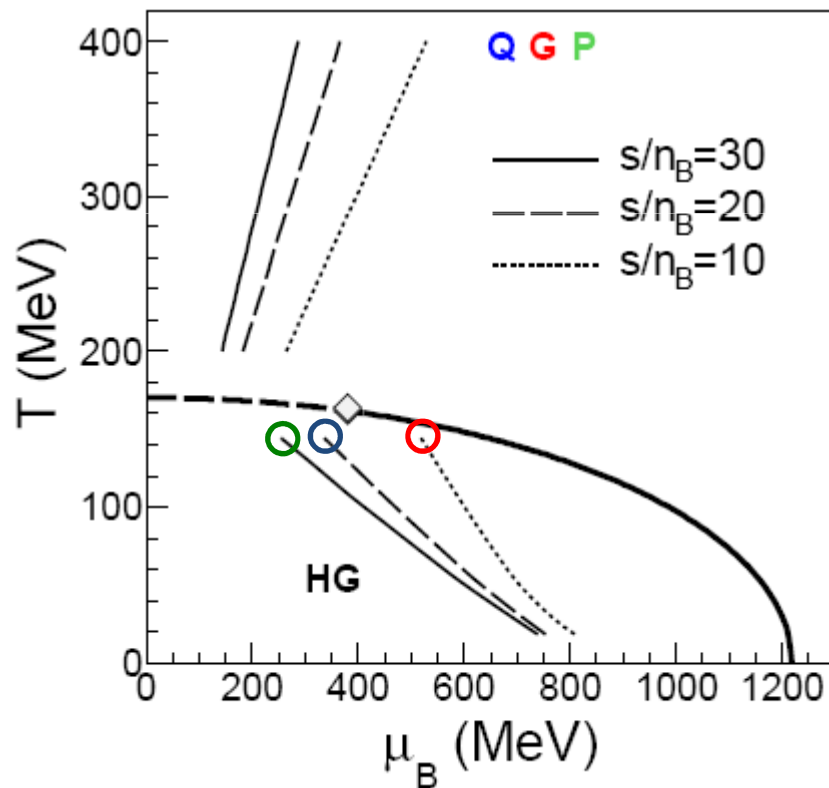
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# Isentropic trajectories (FAIR facility)

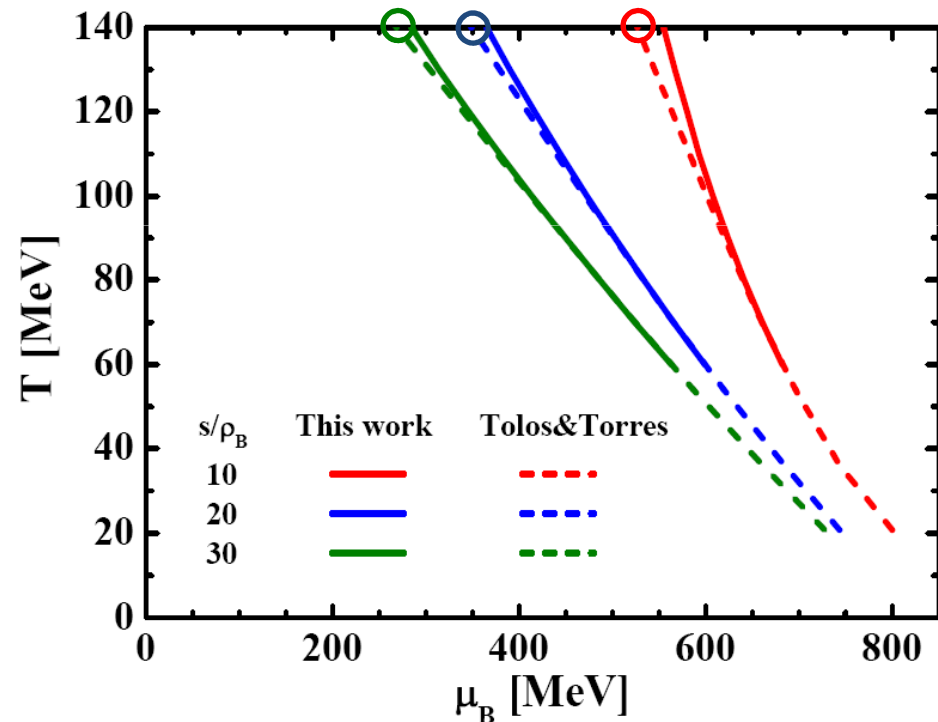
□ Assume a **constant specific entropy** (entropy per net baryon) for **FAIR physics**:

$$\sqrt{s} = 5 - 40 \text{ AGeV} \Leftrightarrow s/n_B = 10 - 30$$

Juan Torres FAIRNESS 2013



L.Tolos, J.M.Torres-Rincon, Phys. Rev. **D88**, 074019 (2013)

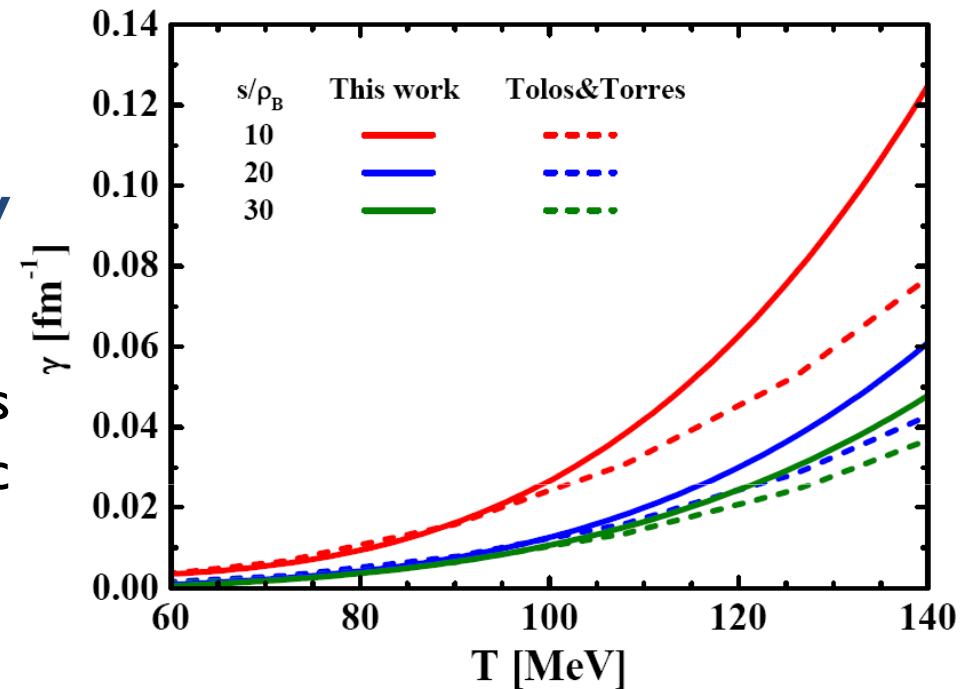


**Small deviation** due to the higher states in our hadronic cocktail

# Thermal relaxation rate (FAIR facility)

- **strong dependence** on the isentropic trajectory
- **baryons** contribute **significantly** for finite baryochemical potential
- **deviation** at higher temperatures due to **higher states** in our hadronic cocktail

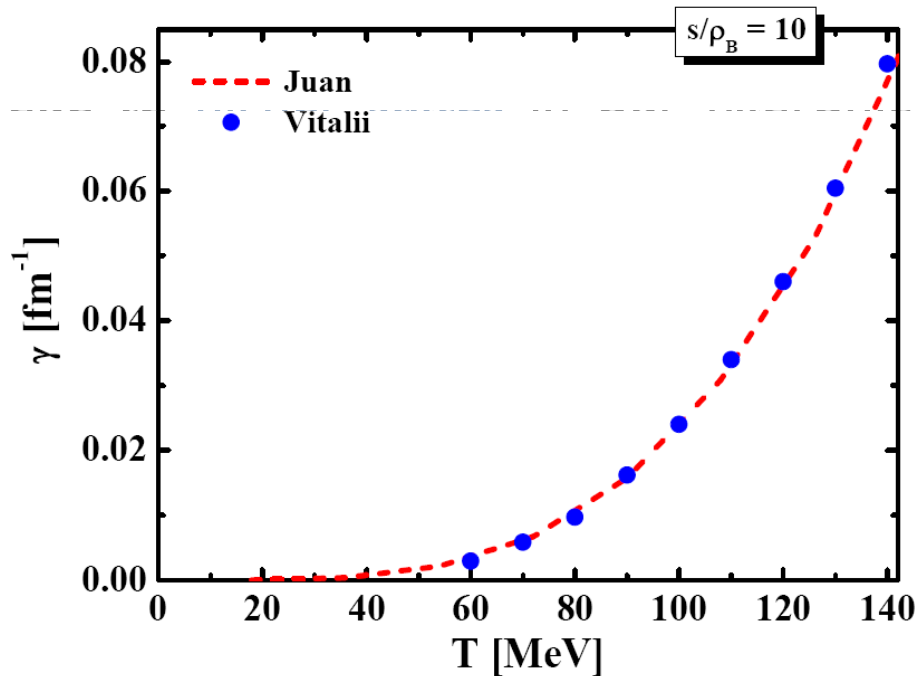
L.Tolos, J.M.Torres-Rincon, Phys. Rev. **D88**, 074019 (2013)



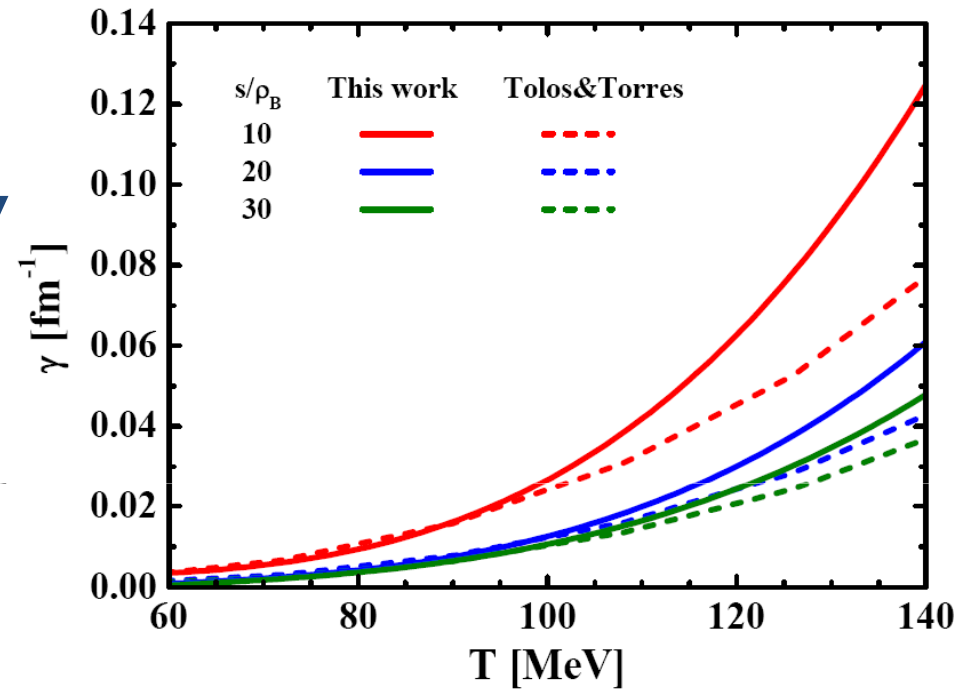
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L.Tolos, J.M.Torres-Rincon, Phys. Rev. **D88**, 074019 (2013)

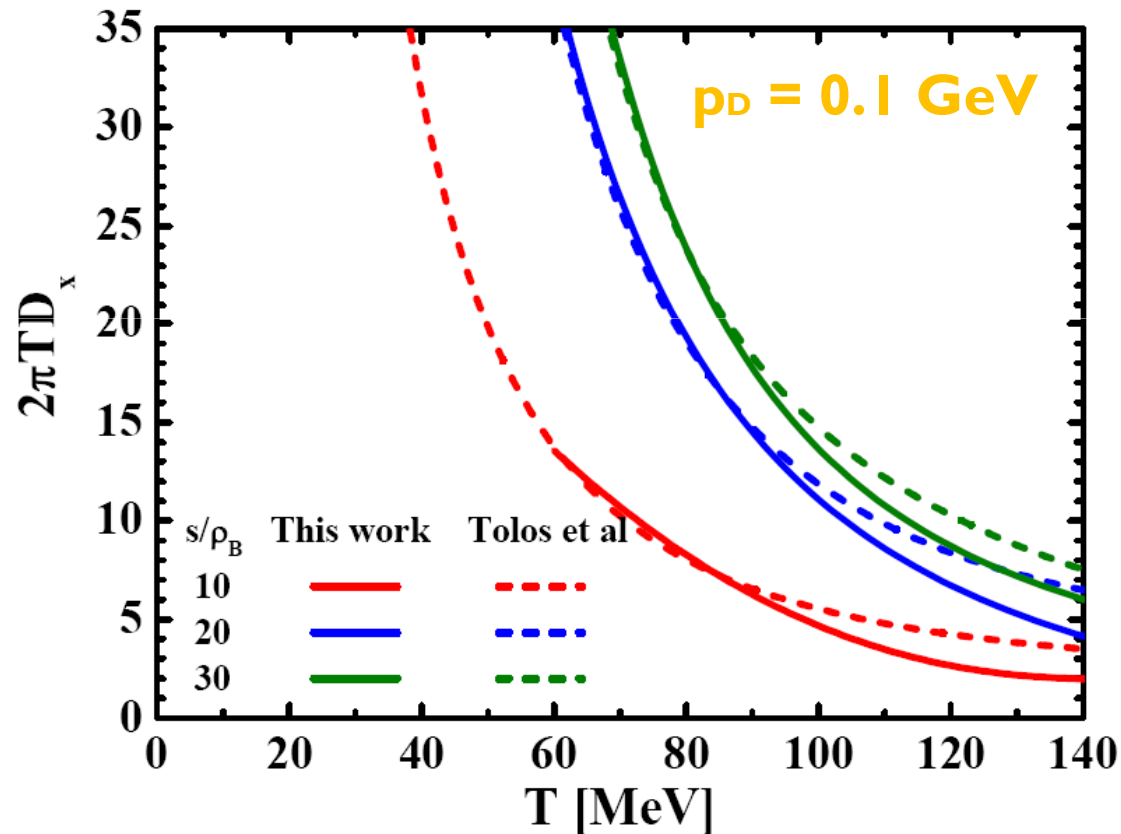


**Perfect agreement**

# Spatial diffusion coefficient (FAIR facility)

□ **Spatial diffusion coefficient:**  $D_x = \lim_{p \rightarrow 0} \frac{B}{m_D^2 \gamma}$

L.Tolos, J.M.Torres-Rincon, Phys. Rev. **D88**, 074019 (2013)



the higher states are **important**, but less known

# Summary

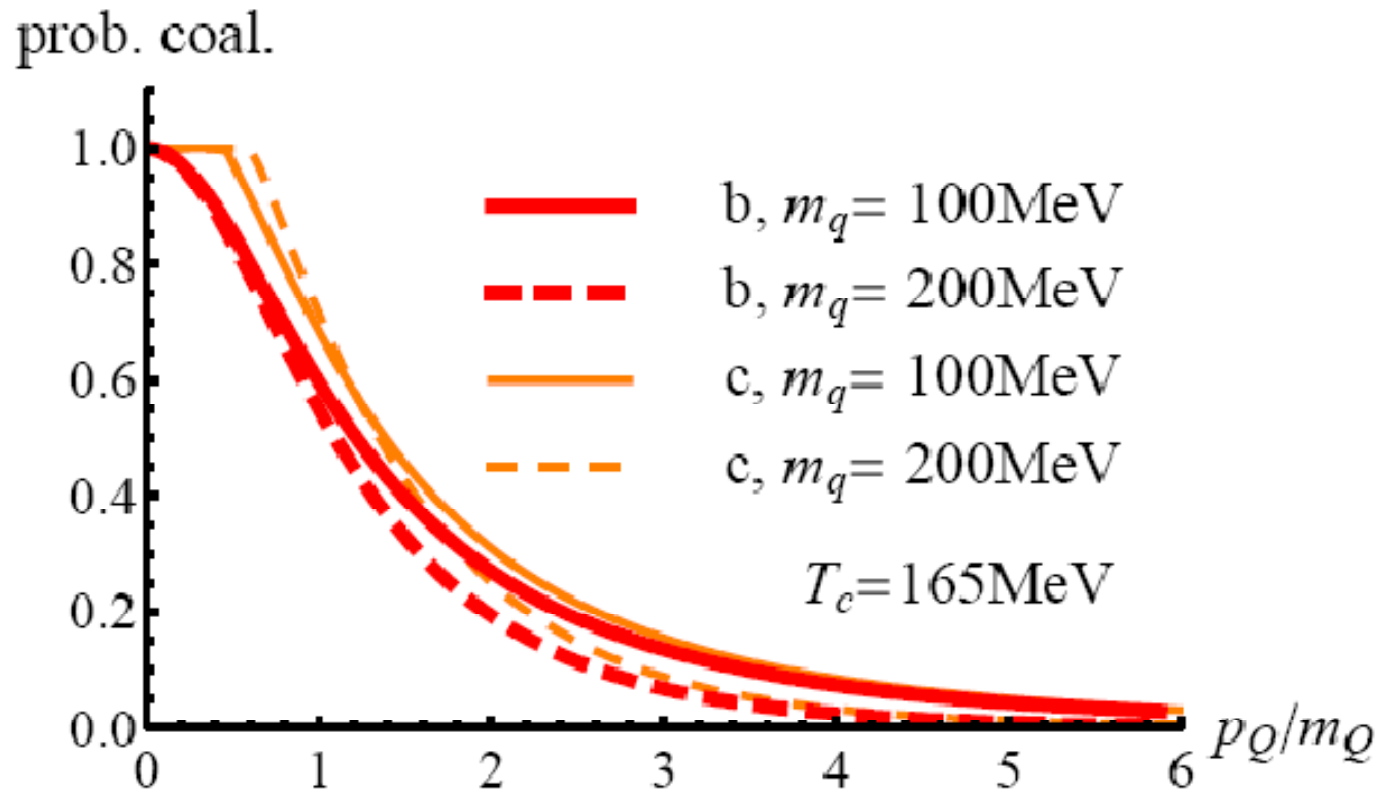
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- Introduce the **effective chemical potentials** for all **hadron species** that are subject to strong decays
- Calculate the **D-meson transport coefficients** as functions of **momentum** and **temperature** for three scenarios
- The **presence** of D-meson rescattering in HG is **almost invisible** for the **RAA**, but shows a **systematic contribution of 1%-2%** to the  **$v_2$**  of D-meson and of single nonphotonic leptons originating from the decays of heavy mesons
- Transport coefficients **strongly depend** on the isentropic trajectory
- The **spatial diffusion coefficient** is **sensitive** to the higher states in the HG at higher temperature (extension of the calculations to **FAIR** case)

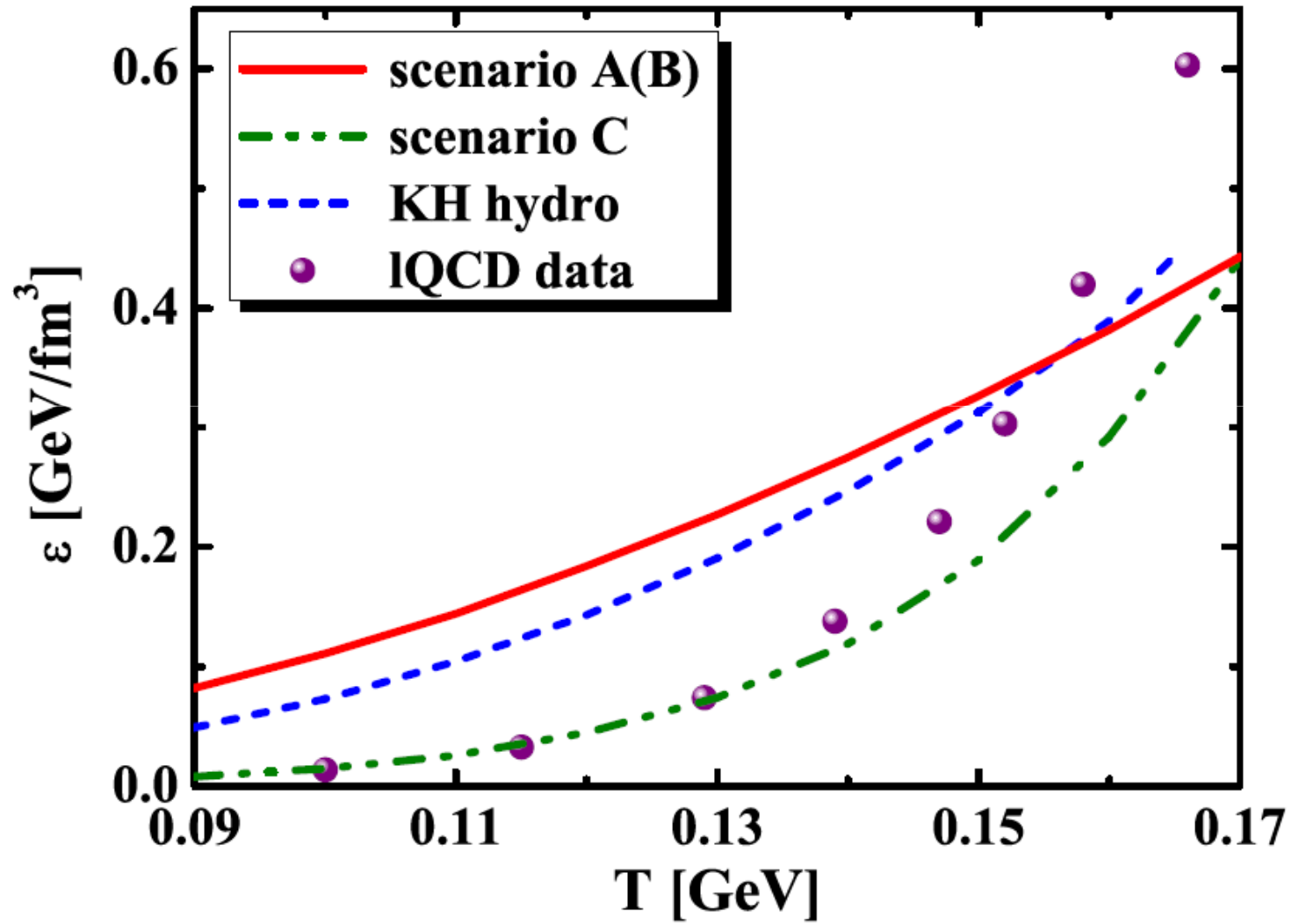
**Back up**



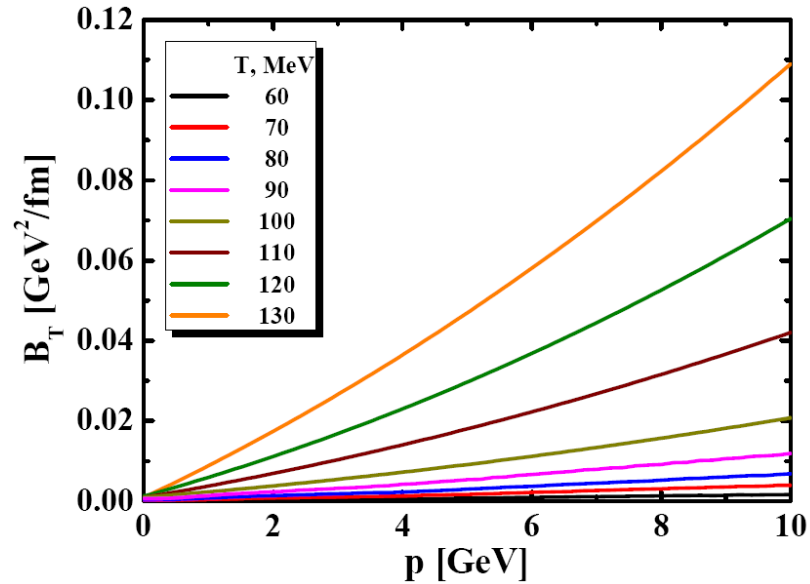
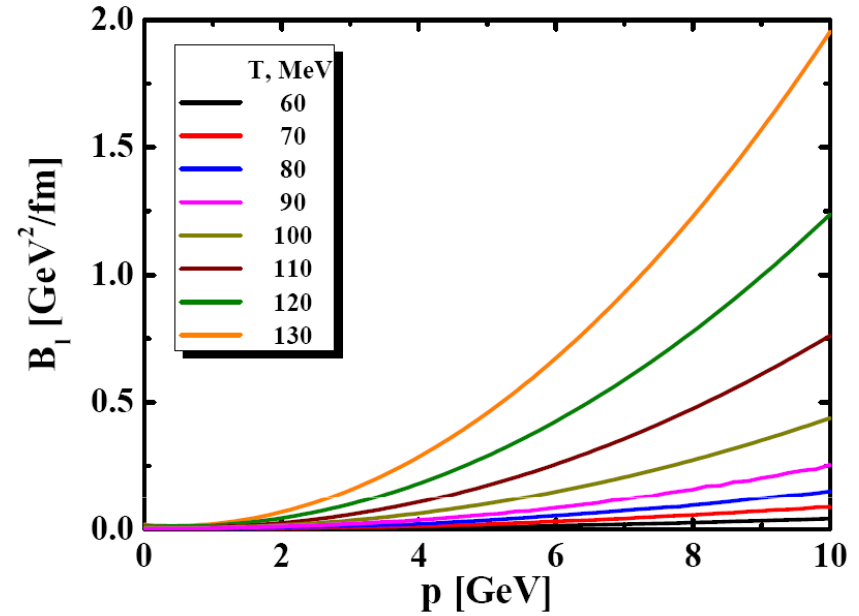
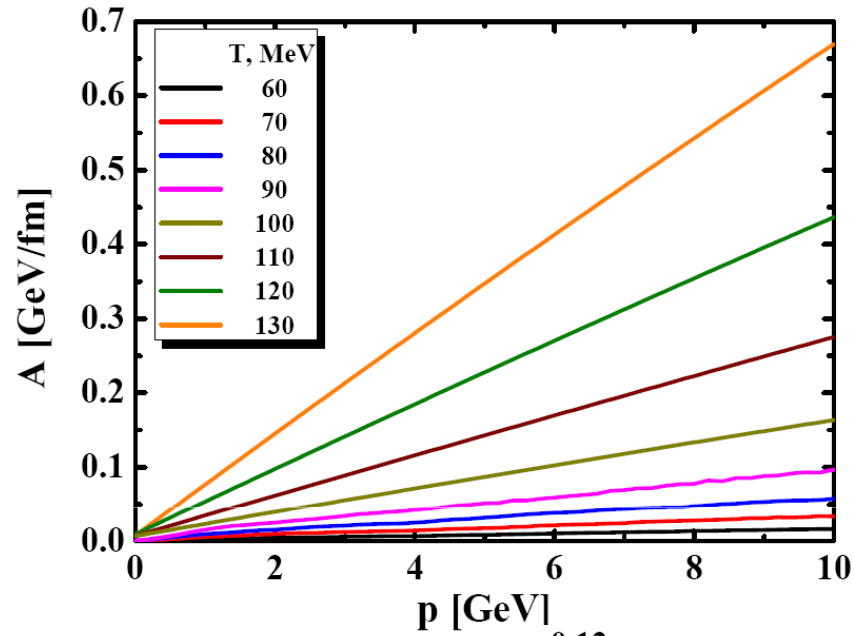
# Hadronization of HQ



# Hadronic equation of state



# D-meson transport coefficient ( $s/n_B = 10$ )



# Spatial diffusion coefficient (Juan)

