

# Dynamical study of fluctuations in relativistic nuclear collisions

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**H-QM** | Helmholtz Research School  
Quark Matter Studies

- Introduction
- Fluctuations in the number of participants
- Multiplicity fluctuations
- Baryon number and electric charge fluctuations
- Energy dependence of multiplicity fluctuations
- Fluctuations at RHIC

In collaboration with Elena Bratkovskaya (Frankfurt) and Mark Gorenstein (Kiev)

# Study of fluctuations within transport models

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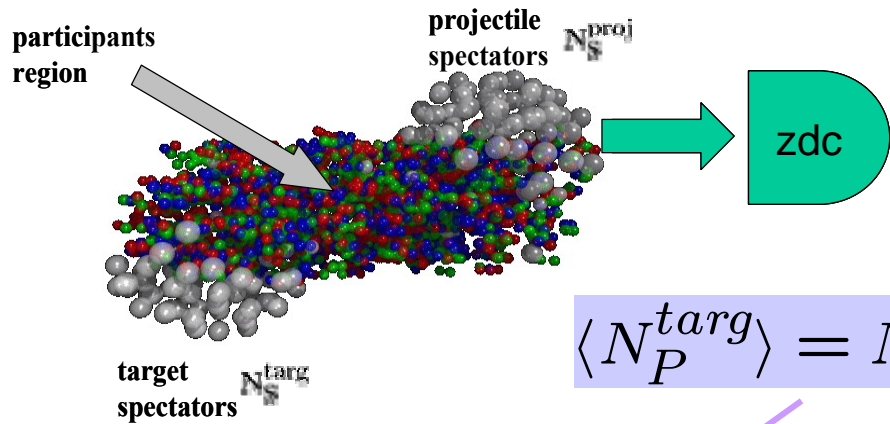
**HSD** – **H**adron-**S**tring-**D**ynamics transport approach

**UrQMD** – **U**ltra-**r**elativistic-**Q**uantum-**M**olecular-**D**ynamics

Transport models allow to study:

- ❑ **statistical and dynamical** fluctuations
- ❑ **event-by-event** analysis similar to the experiment
- ❑ the influence of the experimental acceptance  
on the final results on fluctuations
- ❑ the **centrality** dependence
- ❑ the **energy** dependence of fluctuations

# Fluctuations in the number of participants in HSD & UrQMD

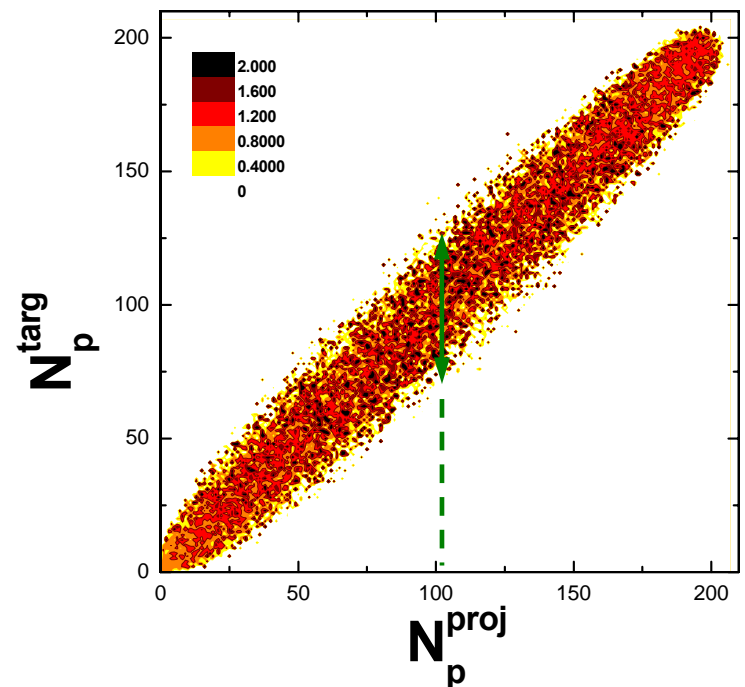
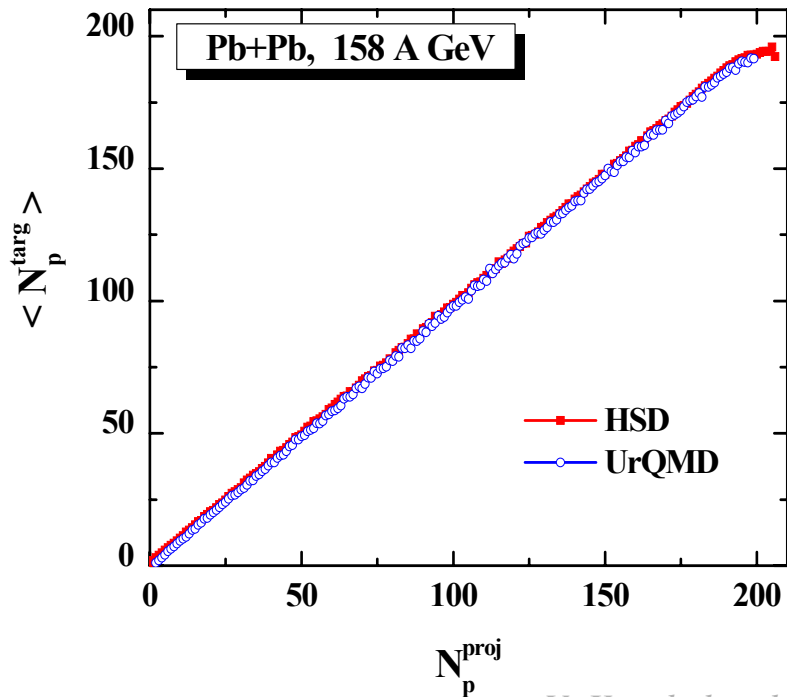


$$N_P^{proj} = A - N_S^{proj} - \text{projectile participants}$$

$$N_P^{targ} = A - N_S^{targ} - \text{target participants}$$

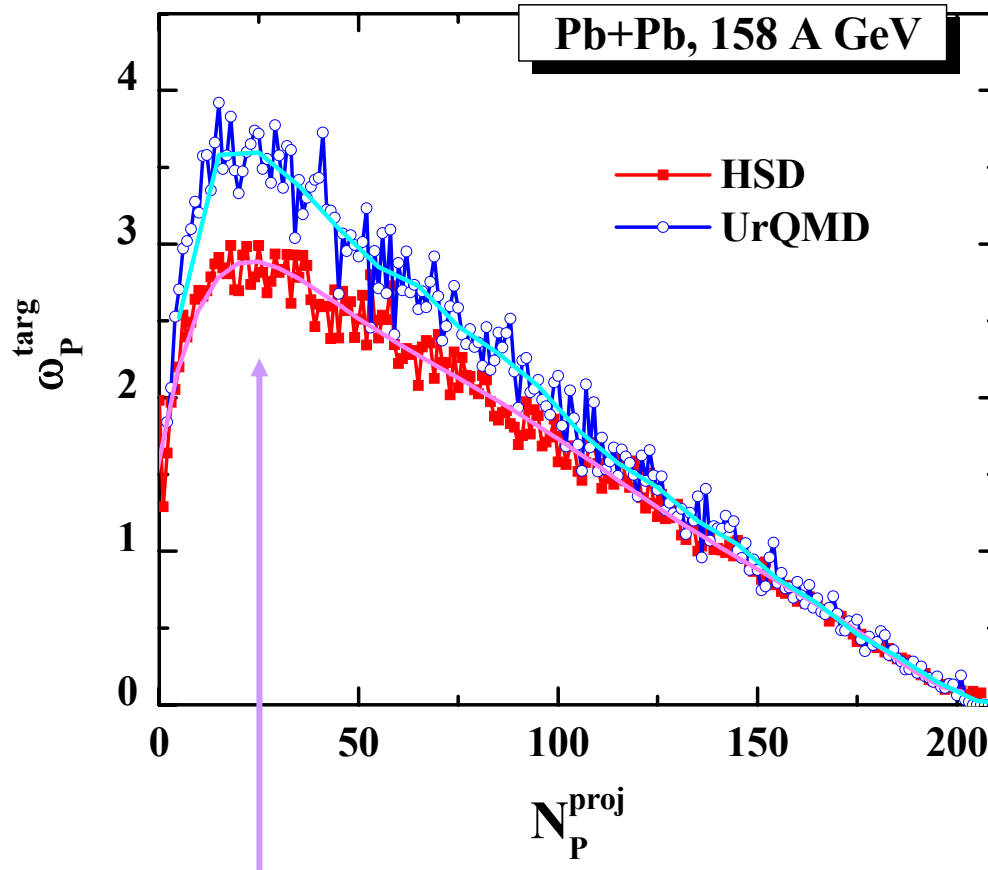
$$\langle N_P^{targ} \rangle = N_P^{proj}$$

$$N_P^{targ} \neq N_P^{proj}$$



# Fluctuations in the number of participants in HSD & UrQMD

PRC 73 (2006) 034902



The scaled variance of “X”:

$$\omega_X = \frac{\langle X^2 \rangle - \langle X \rangle^2}{\langle X \rangle}$$

$\omega_P^{\text{targ}}$  is the scaled variance of the number of target participants  $N_P^{\text{targ}}$

in each sample with  $N_P^{\text{proj}} = \text{const}$  the number of target participants  $N_P^{\text{targ}}$  fluctuates considerably:  $\omega_P^{\text{targ}} \sim 3$  at  $N_P^{\text{proj}} \sim 25$

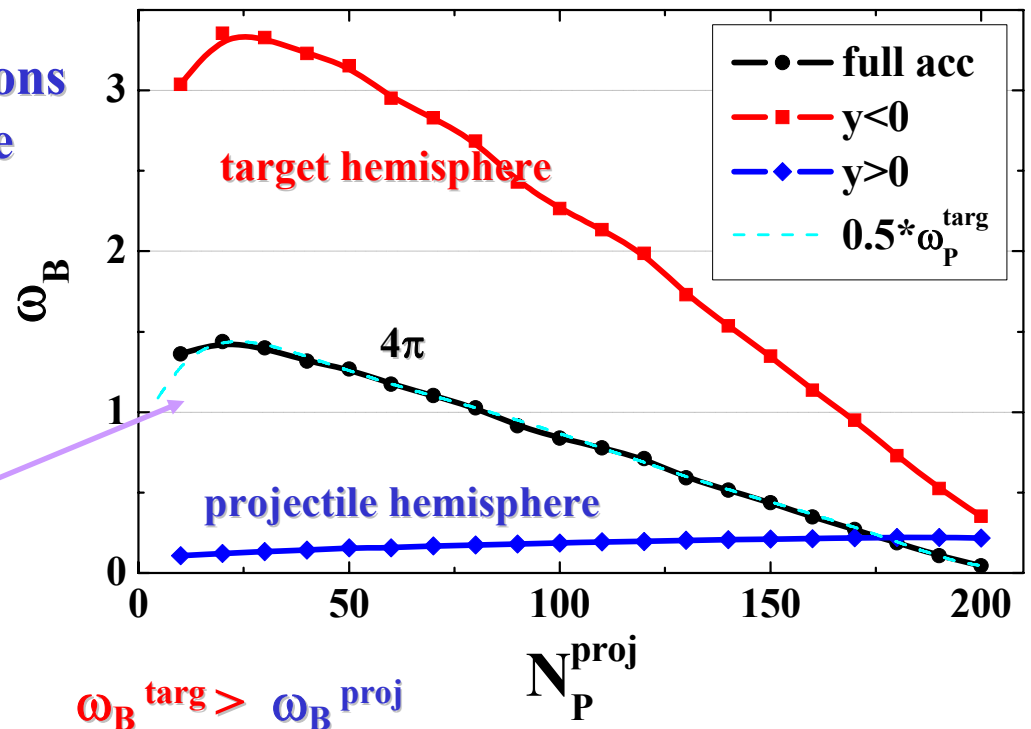
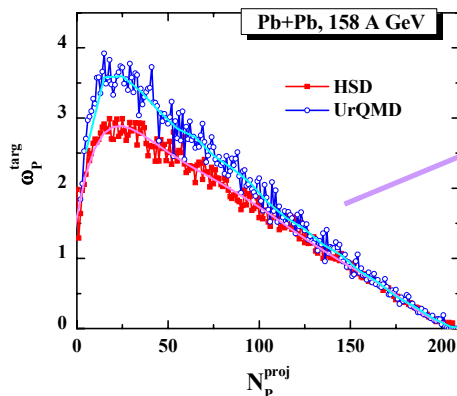
# Baryon number fluctuations (HSD)

Baryon number is a conserved quantity:

the **net baryon number in the full phase space**  $B \equiv N_B - N_{\bar{B}}$   
 is equal to the total number of participants  $N_p = N_p^{\text{proj}} + N_p^{\text{targ}}$

at fixed  $N_p^{\text{proj}}$  the  $N_p$  number  
 fluctuates due to the fluctuations  
 of  $N_p^{\text{targ}}$  in the full phase space

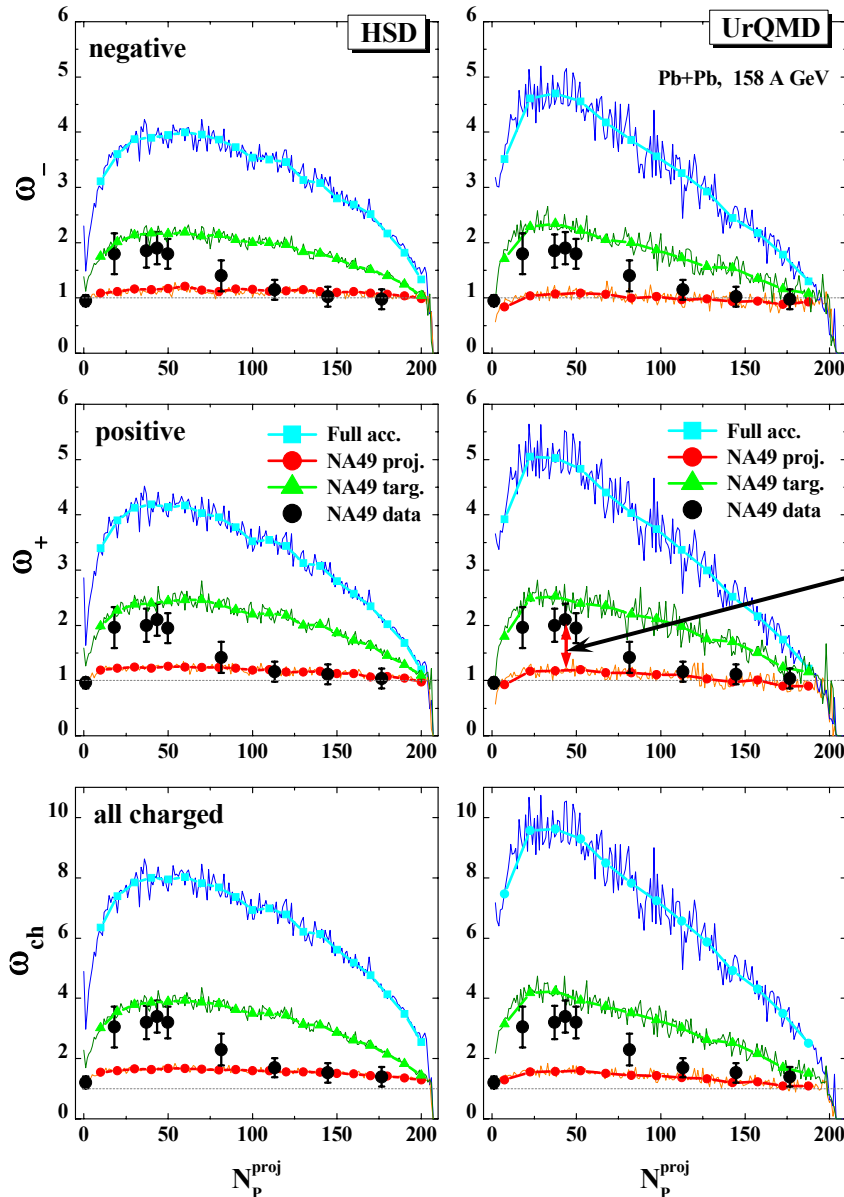
$$\omega_B \approx \frac{1}{2} \omega_P^{\text{targ}}$$



$$\omega_B^{\text{targ}} > \omega_B^{\text{proj}}$$

PRC 74 (2006) 064911

# Multiplicity fluctuations in projectile and target hemispheres

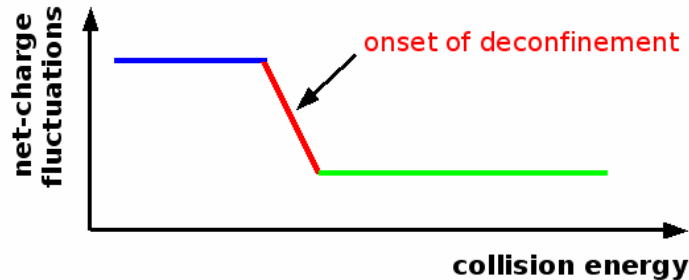


- HSD and UrQMD show **strong multiplicity fluctuations** of dynamical origin;
- HSD and UrQMD show **strong multiplicity fluctuations in  $4\pi$  ,full' acceptance**, however, the observed (by NA49) non-trivial system size dependence of multiplicity fluctuations is not reproduced by HSD and UrQMD
- $\phi$ - $p_T$  acceptance cuts strongly decrease fluctuations !
- Fluctuations in **target hemisphere are larger than in projectile hemisphere**

PRC 73 (2006) 034902

# Net-electric charge fluctuations

sensitive to the **EoS** at the early stage of the collision and to its changes in the deconfinement phase transition region



Jeon, Koch, PRL85 (2000) 2076

Asakawa, Heinz, Muller PRL85 (2000) 2072

net-charge fluctuations are smaller in QGP than in a hadron gas

## Definition of $\Phi$ -measure of fluctuations:

Gazdzicki, Mrowczynski, Z. Phys. C 54 (1992) 127:

$$\Phi_x = \sqrt{\frac{\langle Z^2 \rangle}{\langle N \rangle}} - \sqrt{\overline{z^2}}, \quad z_i = x_i - \bar{x}, \quad Z = \sum_{i=1}^N (x_i - \bar{x}_i)$$

$X$  – single-particles variable  
 $N$  – number of particles in 1 event  
 $\langle \dots \rangle$  - averaging single-particle distribution over events

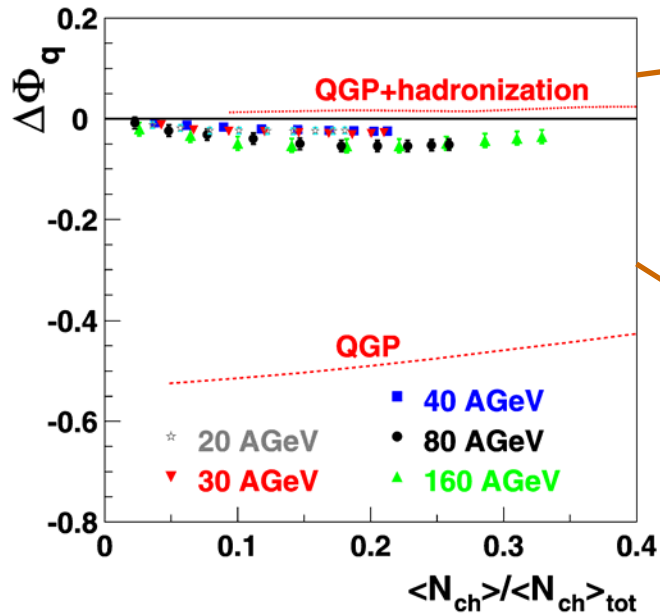
$$\Delta\Phi_q = \Phi_q - \Phi_{q,GCC}$$

$$\Phi_{q,GCC} = \sqrt{1-P} - 1, \quad P = \frac{\langle N_{ch} \rangle}{\langle N_{ch} \rangle_{tot}}$$

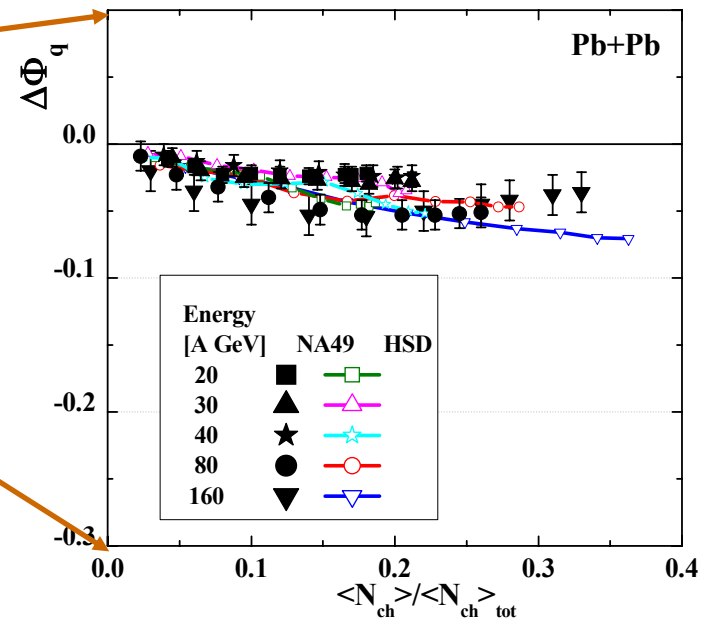
**GCC** – global charge conservation

$\Delta\Phi_q = 0$ , if particles are correlated by global charge conservation **only!**

# NA49: Net-electric charge fluctuations



NA49, PRC 70 (2004) 064903



PRC 74 (2006) 064911

**The decay of resonances strongly modifies the initial QGP fluctuations!**

# Energy Dependence of Multiplicity Fluctuations in N+N and central A+A

PLB, in press [nucl-th/0703052]

➤ HSD gives good description of experimental data for charged hadron multiplicities in A+A.

➤ Fluctuations in p+p and A+A are very close within HSD.

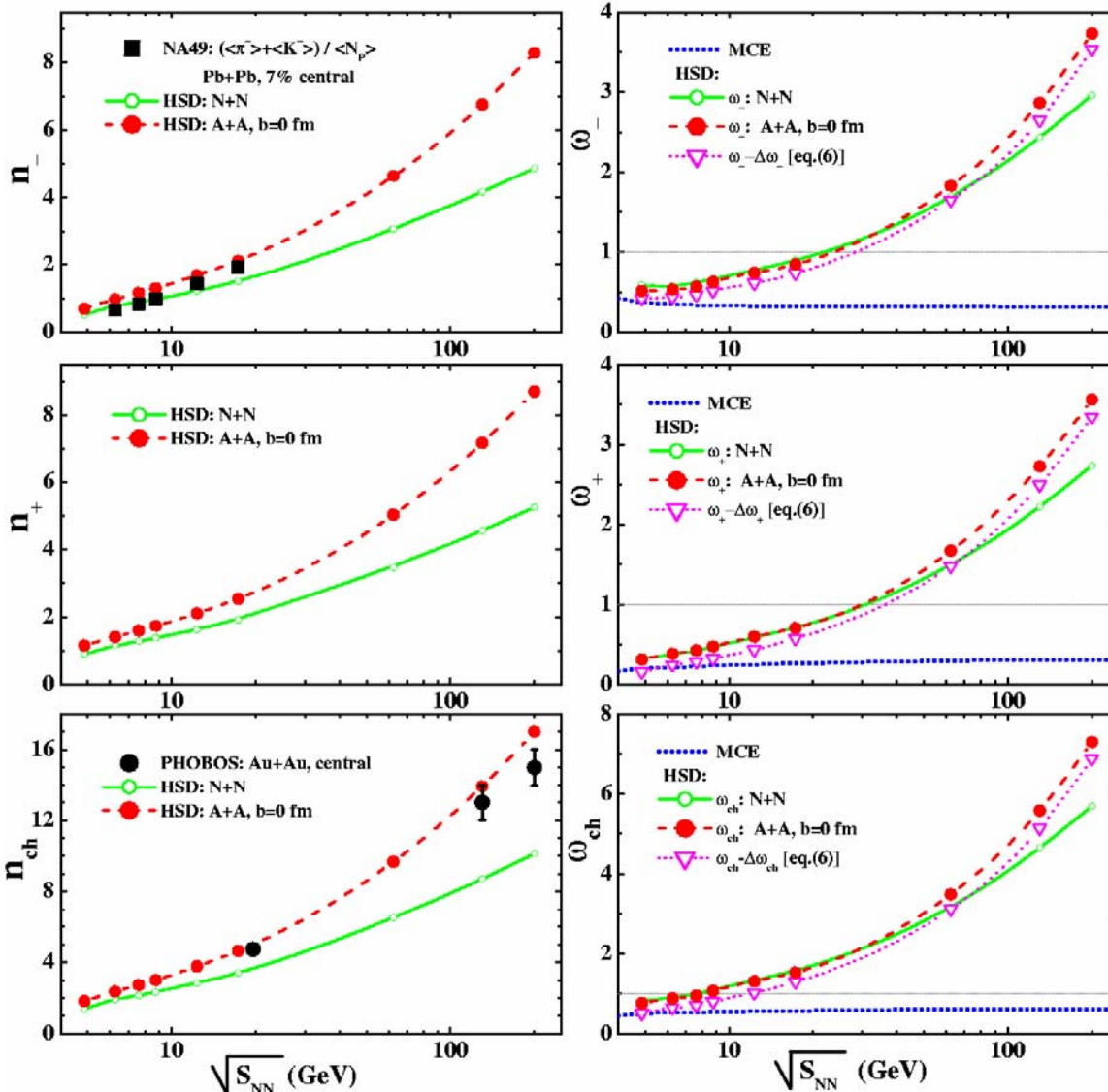
➤ Fluctuations of the number of participants do not explain excess over p+p fluctuations

$$\omega_i = \omega_i^\nabla + n_i \omega_P \quad (6)$$

➤ Statistical model shows very small and energy independent fluctuations and contradicts to the transport calculations where  $\omega$  reaches significant values for the large energies.

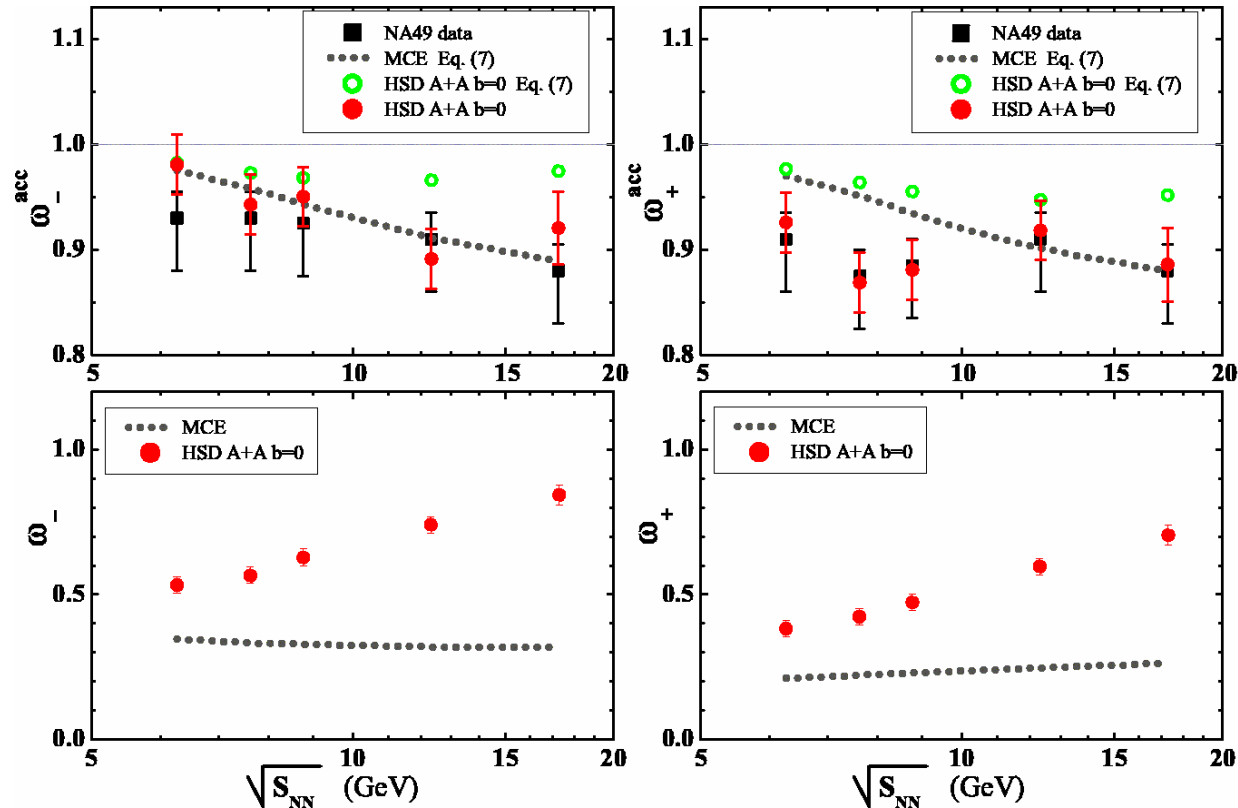
Begun et al, PRC (2006) 044903

Similar results for  $\omega^\pm$  from UrQMD (see talk of M.Bleicher)



# Energy Dependence of Multiplicity Fluctuations: NA49 data

Acceptance scaling formulae:  $\omega_i^{acc} = 1 - q_i + q_i \omega_i^{4\pi}$  (7)



NA49 cannot clearly distinguish between statistical and transport models because of small acceptance  $q=0.04\dots 0.16$  and small differences between the model predictions in this range of energy !

PLB, in press [nucl-th/0703052]

# Fluctuations at RHIC

0704.1831[nucl-th]

➤ Fluctuations of the participant number within different centrality classes in PHENIX for Au+Au @  $\sqrt{s} = 200$  GeV are **significant!**

➤ Using **the model of independent sources:**

$$\omega_i = \omega_i^* + n_i \omega_P$$

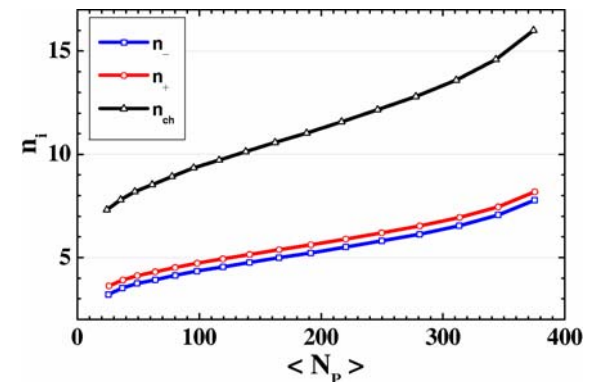
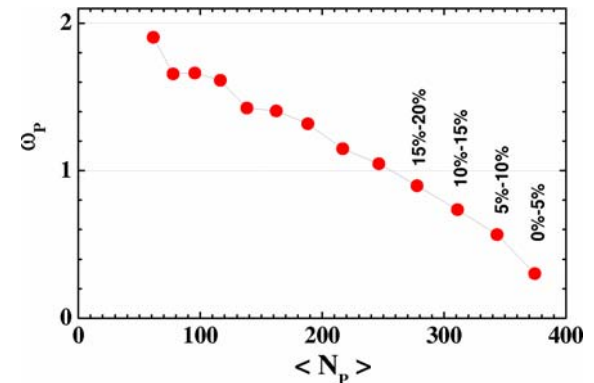
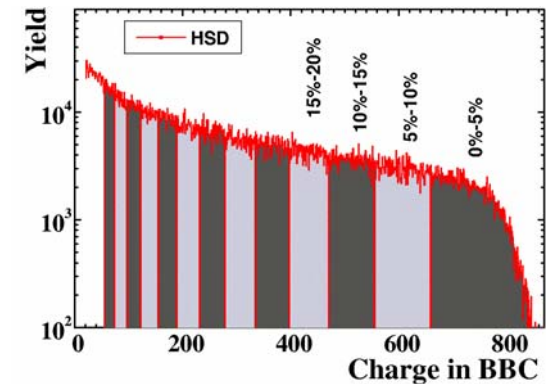
and acceptance scaling formula:

$$\omega_i^{acc} = 1 - q_i + q_i \omega_i$$

we obtain equation:

$$\omega_i^{acc} = 1 - q_i + q_i \omega_i^* + q_i n_i \omega_P$$

➤ From HSD we took only:  $\omega_p$ ,  $q_i$ ,  $n_i$



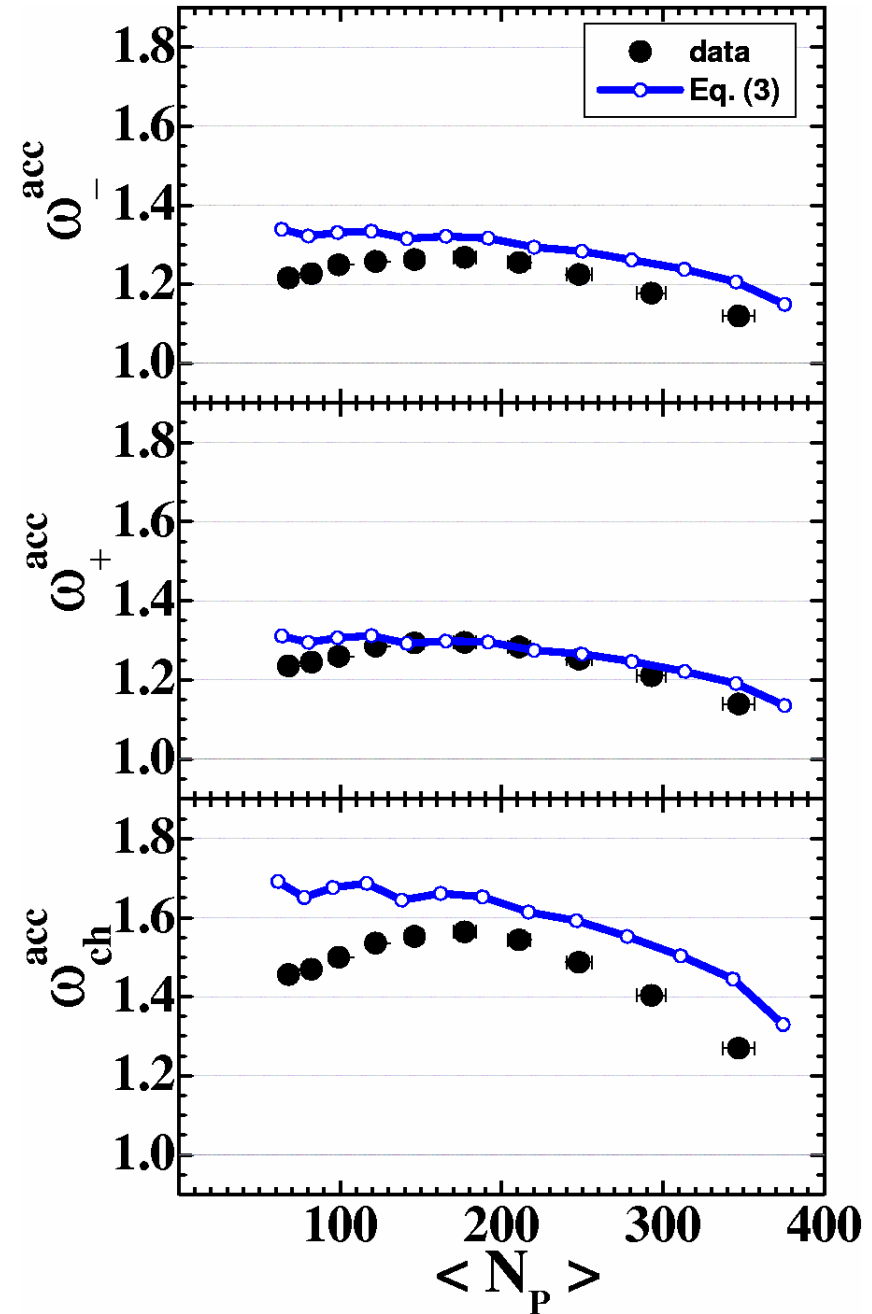
# Fluctuations at RHIC

0704.1831[nucl-th]

$$\omega_i^{acc} = 1 - q_i + q_i \omega_i^* + q_i n_i \omega_P \quad (3)$$

eq.(3) describes exp. data from PHENIX  
for Au+Au @  $\sqrt{s} = 200$  GeV

Multiplicity fluctuations represent  
fluctuations of **participants!**



# Conclusions

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- ❑ **The fluctuations in the number of target participants – for fixed projectile participants - strongly influence all observable fluctuations**
- ❑ **The measured fluctuations of the electric charge in different acceptance windows are consistent with HSD results.**
- ❑ **Statistical and transport models show different results in central A+A collisions for multiplicity fluctuations versus energy. New measurements at higher energies and with larger acceptance are needed!**
- ❑ **In collider-type experiments the fluctuations of the number of participants are significant. To avoid them one has to consider the most central collisions with more rigid events selection.**

## **Thanks to:**

**Mark Gorenstein**

**Michael Hauer**

**Viktor Begun**

**Marek Gazdzicki**

**Benjamin Lungwitz**

**Elena Bratkovskaya**

**Olena Linnyk**

**Horst Stöcker**

**Marcus Bleicher**

**Stephane Haussler**



# Model of independent sources

Gazdzicki, Gorenstein, hep-ph/0511058

**Number of independent sources** is proportional to number of participants:

$$\omega_i = \omega_i^* + \frac{1}{2} \omega_P^{\text{targ}} n_i$$

$\omega_i^*$  is the fluctuation from a single NN source

fluctuation in the number of nucleon participants in A+A

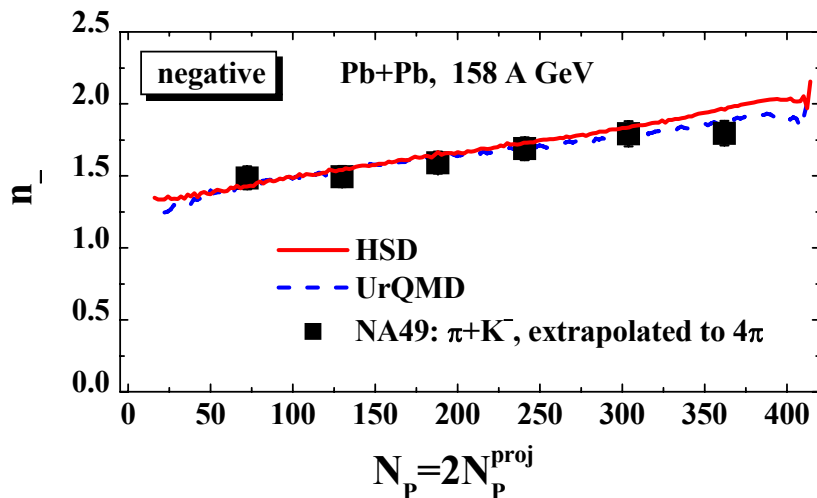
$\omega_-^* = 1.5$   
 $\omega_+^* = 1.1$   
 $\omega_{\text{ch}}^* = 2.5$

**HSD: N+N** (averaged over pp+nn+pn)

$n_i$  is the particle number of  $i$ -th type ( $i=+,-,\text{ch}$ ) per participant  $N_p = N_p^{\text{proj}} + N_p^{\text{targ}}$

$$n_i = \frac{\langle N_i \rangle}{\langle N_P \rangle}$$

$\omega_P^{\text{targ}}$  is the scaled variance for the fluctuations of the number of target participants  $N_p^{\text{targ}}$

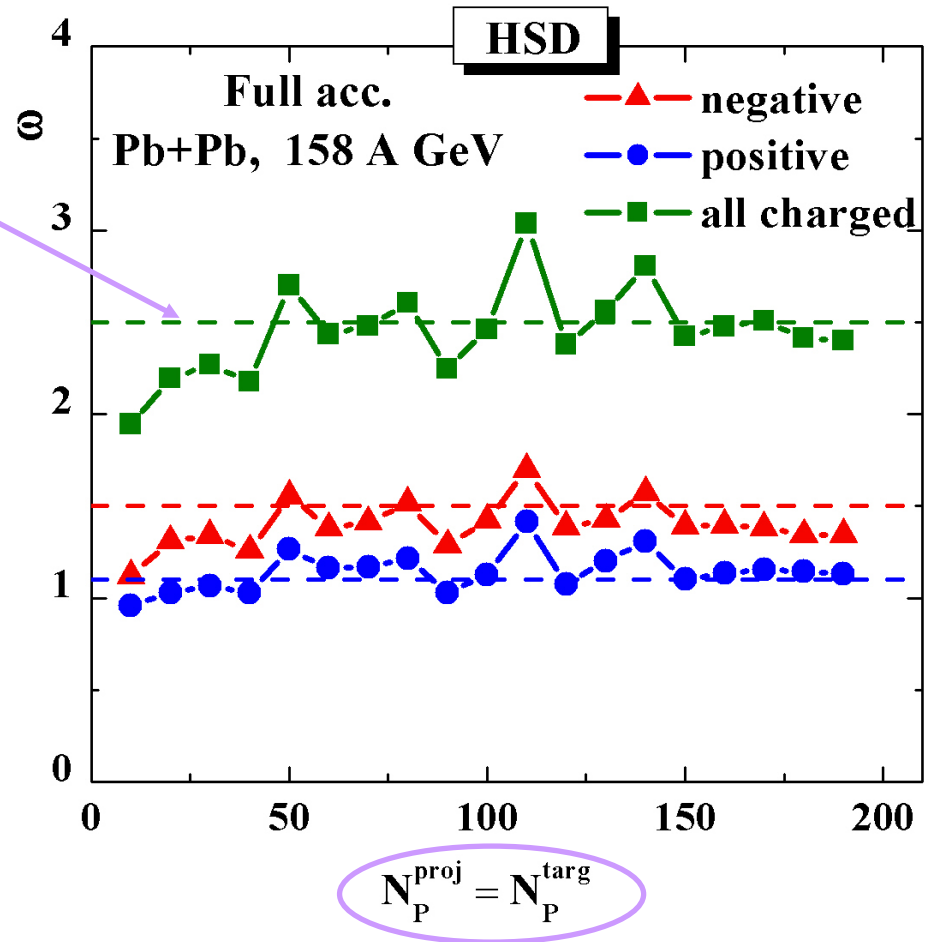


# Model of independent sources

$$\omega_i = \omega_i^* + \frac{1}{2} \omega_P^{\text{targ}} n_i$$

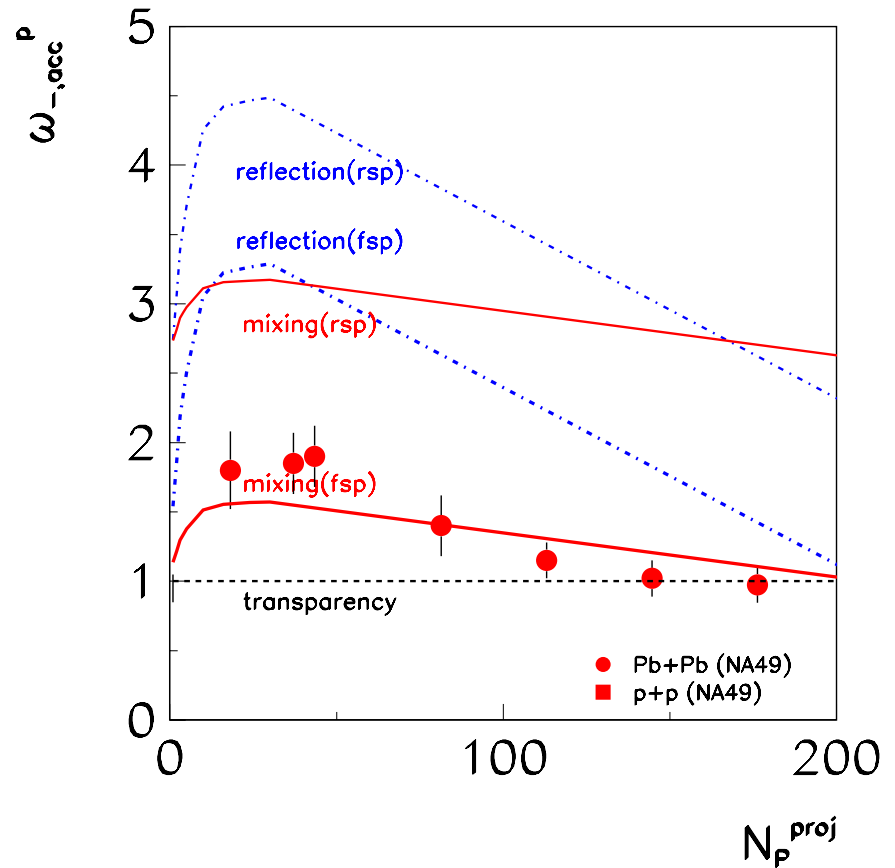
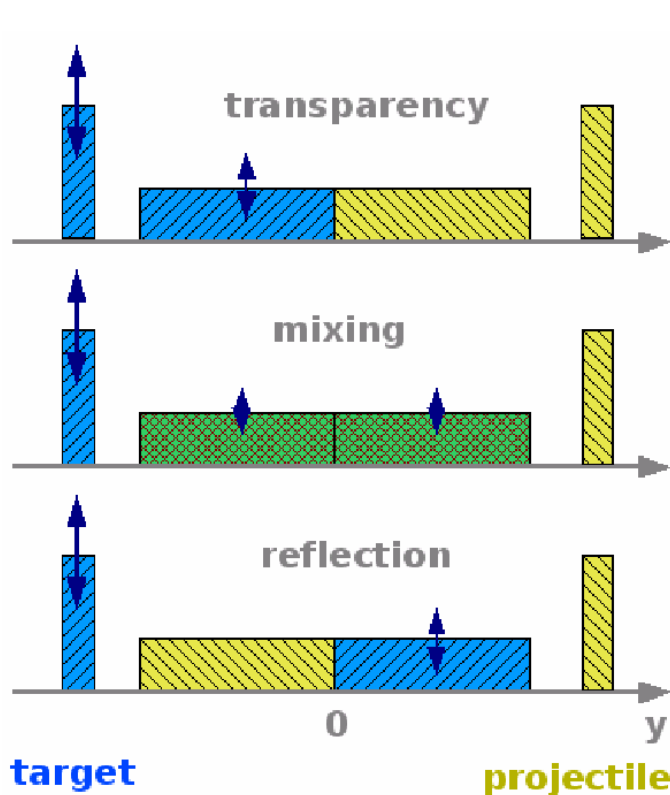
for  $N_p^{\text{proj}} = N_p^{\text{targ}}$ :  $\omega_P^{\text{targ}} = 0$

**Fluctuations in A+A collisions are dominated by the fluctuations of the particle number in single N+N collisions**



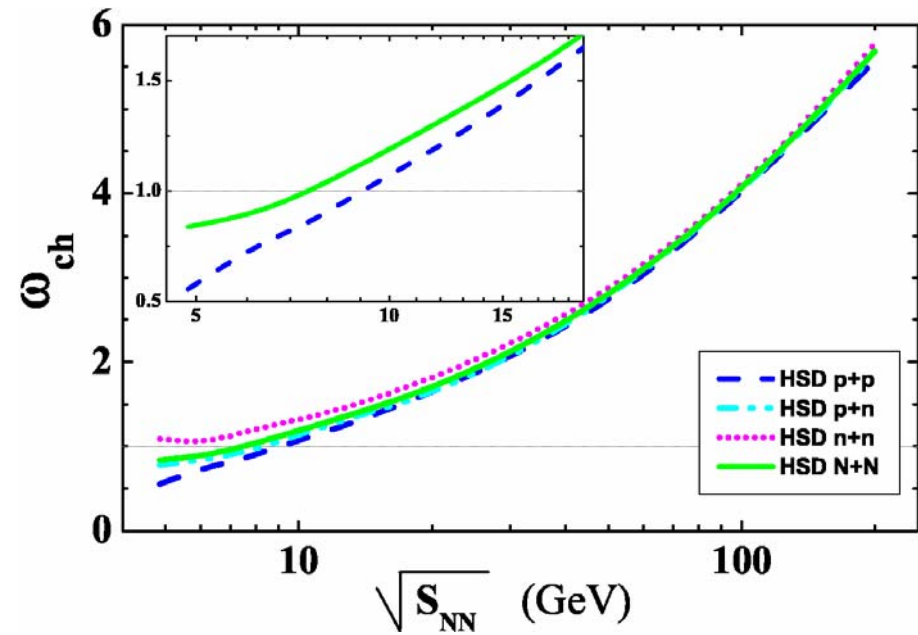
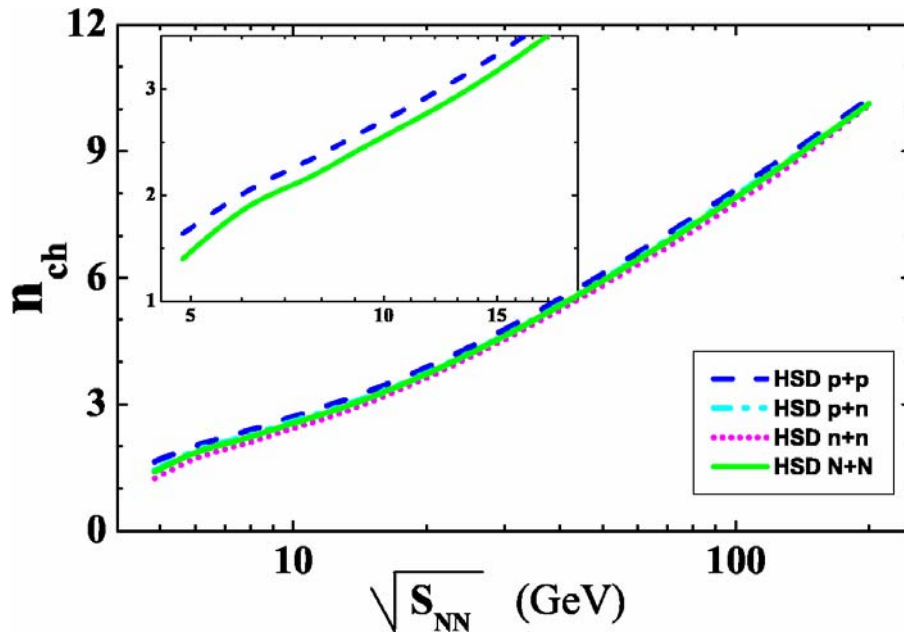
# Multiplicity fluctuations as a probe of transparency, mixing and reflection of initial flows in HIC

Gazdzicki, Gorenstein, hep-ph/0511058:



- Model of independent sources favors **mixing** (but not extreme!) scenario
- HSD & UrQMD show smaller mixing as follows from NA49 data, **i.e. too transparent**

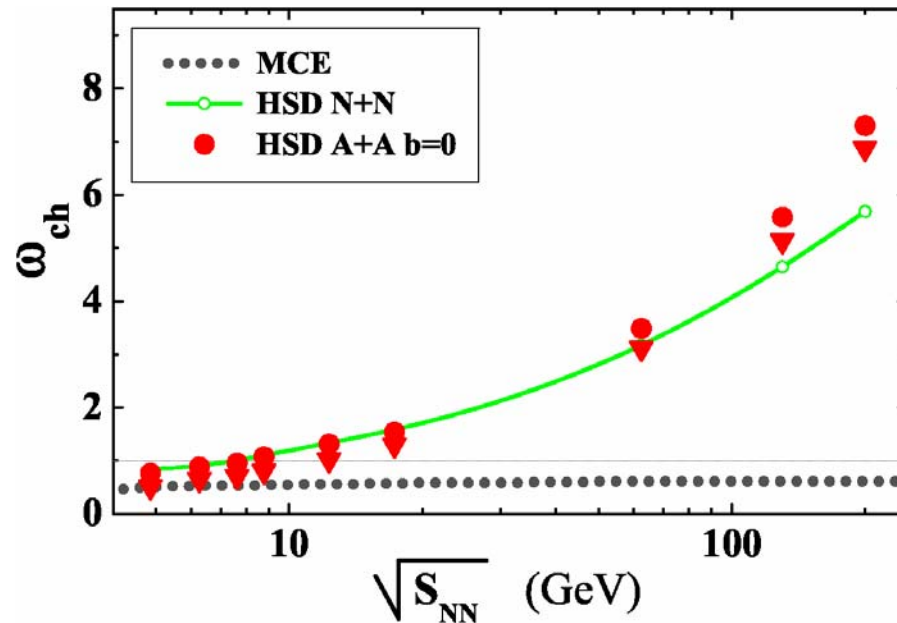
# “Nucleon-nucleon” collisions



$$\langle N_i^{NN} \rangle = \alpha_{pp} \langle N_i^{pp} \rangle + \alpha_{pn} \langle N_i^{pn} \rangle + \alpha_{nn} \langle N_i^{nn} \rangle ,$$

$$\omega_i^{NN} = \frac{1}{\langle N_i^{NN} \rangle} \left[ \alpha_{pp} \omega_i^{pp} \langle N_i^{pp} \rangle + \alpha_{pn} \omega_i^{pn} \langle N_i^{pn} \rangle + \alpha_{nn} \omega_i^{nn} \langle N_i^{nn} \rangle \right]$$

# N+N & A+A central collisions



$$\omega_i = \omega_i^{\blacktriangledown} + n_i \omega_P$$

Pb+Pb @ 158 AGeV,  $b = 0$ :

$$\langle N_P \rangle \simeq 392$$

$$\omega_P \simeq 0.055$$

# Fluctuations at RHIC

