

# Photon Emission near a Critical Point

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Es ist schon Alles gesagt worden,  
aber noch nicht von Allen!

- photon emissivities:  $L(\text{QCD}) \rightarrow \text{LsM}$   
(F. Wunderlich)
- viscosities:  $L(\text{QCD}) \rightarrow \text{Holography}$   
(R. Yaresko, J. Knaute\*)
- $dN_{\text{ch}} / dy$ :  $L(\text{QCD}) \rightarrow \text{Hydro} + \text{CF f.o.}$   
(G. Schlisio)

FAIR & NICA

RHIC & LHC

LHC



**hzdr**

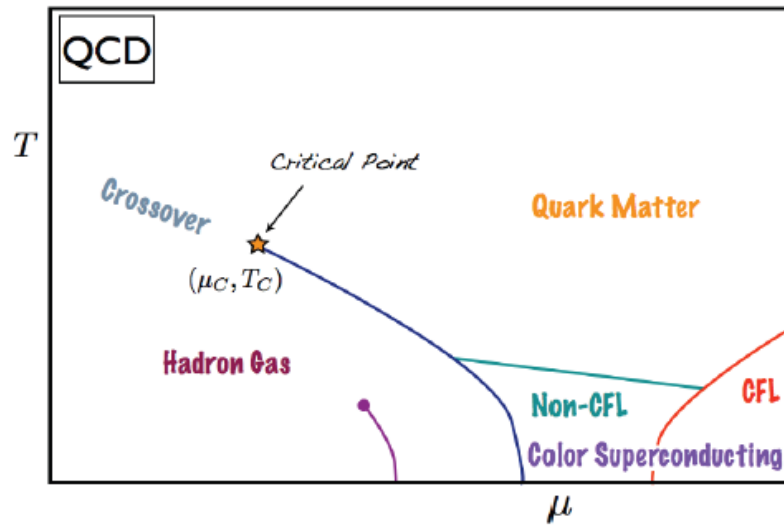


HELMHOLTZ  
ZENTRUM DRESDEN  
ROSSENDORF

\* now in Princeton

# 1. Photon Emissivities and CEP

Hypothesis: a second first-order transition in QCD phase diagram at  $\mu > 0$



deWolfe,  
Gubser,  
Rosen,  
PRD (2011)

reminder: QCD ( $\mu = 0$ ) displays a cross over,  
no reliable lattice data for  $\mu \gg 0$

options for the phase transition:  
hadron-quark vs. gas-liquid  
→ different  $p_c(T)$  curves,  
different isentropic curves

Steinheimer, Randrup, Koch, PRC (2014)  
Hempel, Dexheimer, Schramm, Iosilevskiy, PRC (2013)

# Toy Models (i) CEP

$$s(T, \mu) = s_{\text{reg}}(T, \mu) + s_{\text{sing}}(T, \mu)$$

Bluhm, BK, PoS (2006)  
 based on Nonaka, Asakawa, PRC (2005)  
 based on Giuda, Zinn-Justin, NPB (1997)

3D Ising with proper crit. exps.

special construction

$G(r, h)$ : Gibbs free energy

$M(r, h)$ : magnetization

$h$ : external mag. field

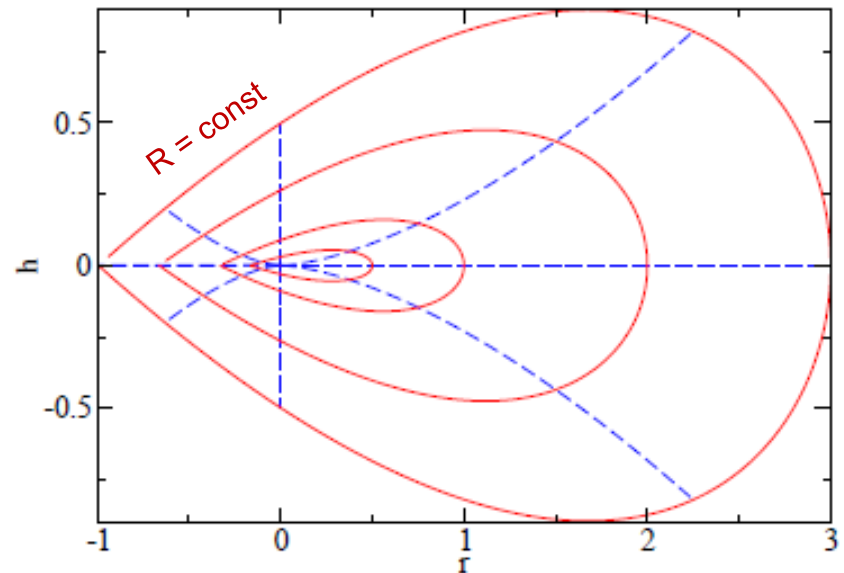
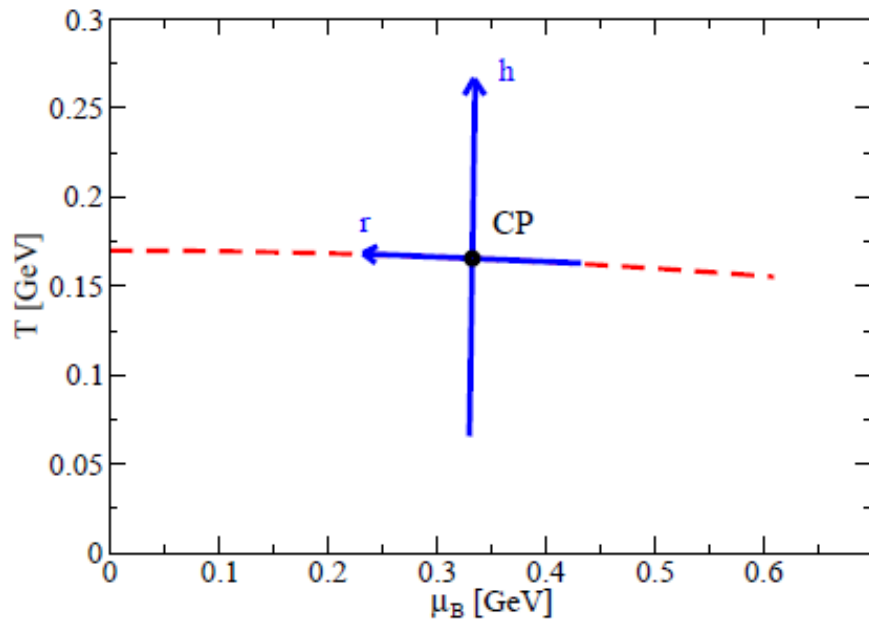
$r = (T - T_c) / T_c$

$$G = h_0 M_0 R^{2-\alpha} g(\theta) - Mh$$

$$r = R(1 - \theta^2),$$

$$h = h_0 R^{\beta\delta} \sum_{i=0}^2 a_{2i+1} \theta^{2i+1},$$

$$\sum_{i=0}^2 a_{2i+1} \theta^{2i+1} (1 - \theta^2 + 2\beta\theta^2) = 2(2 - \alpha)\theta g(\theta) + (1 - \theta^2)g'(\theta)$$



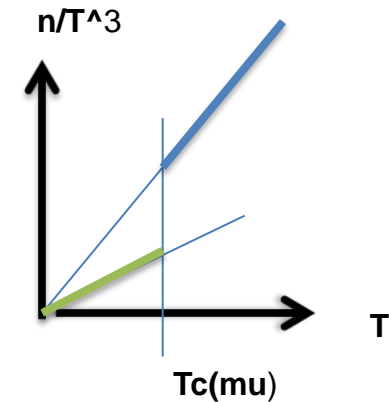
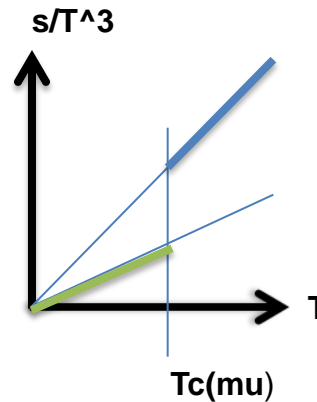
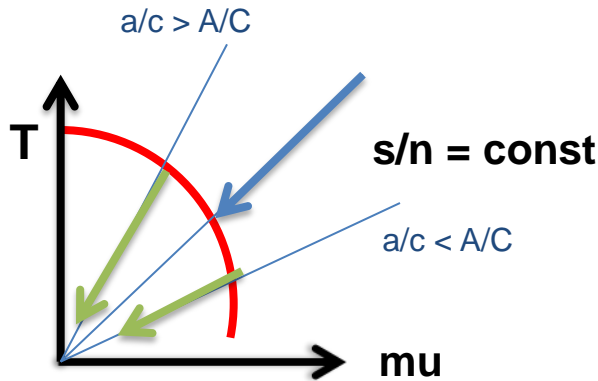
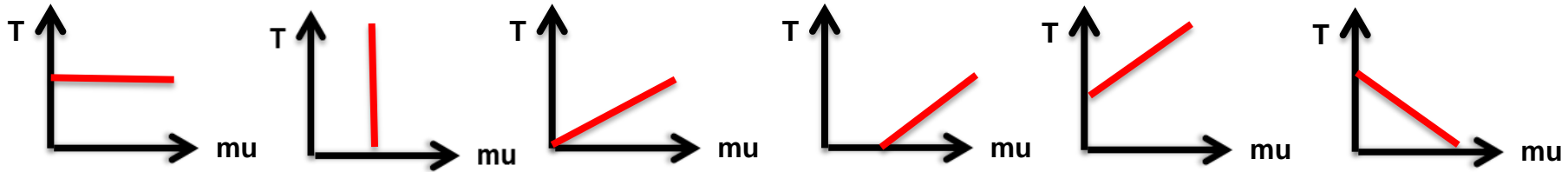
# Toy Models (ii) 1st order transition

ad hoc construction: two-phase model

$$p_1 = a T^4 + c \mu^4 - b$$

$$p_2 = A T^4 + C \mu^4 - B$$

phase border curve/coexistence region:  $p_1 = p_2 \rightarrow T_c(\mu)$ ,  $p_c(T)$  etc.



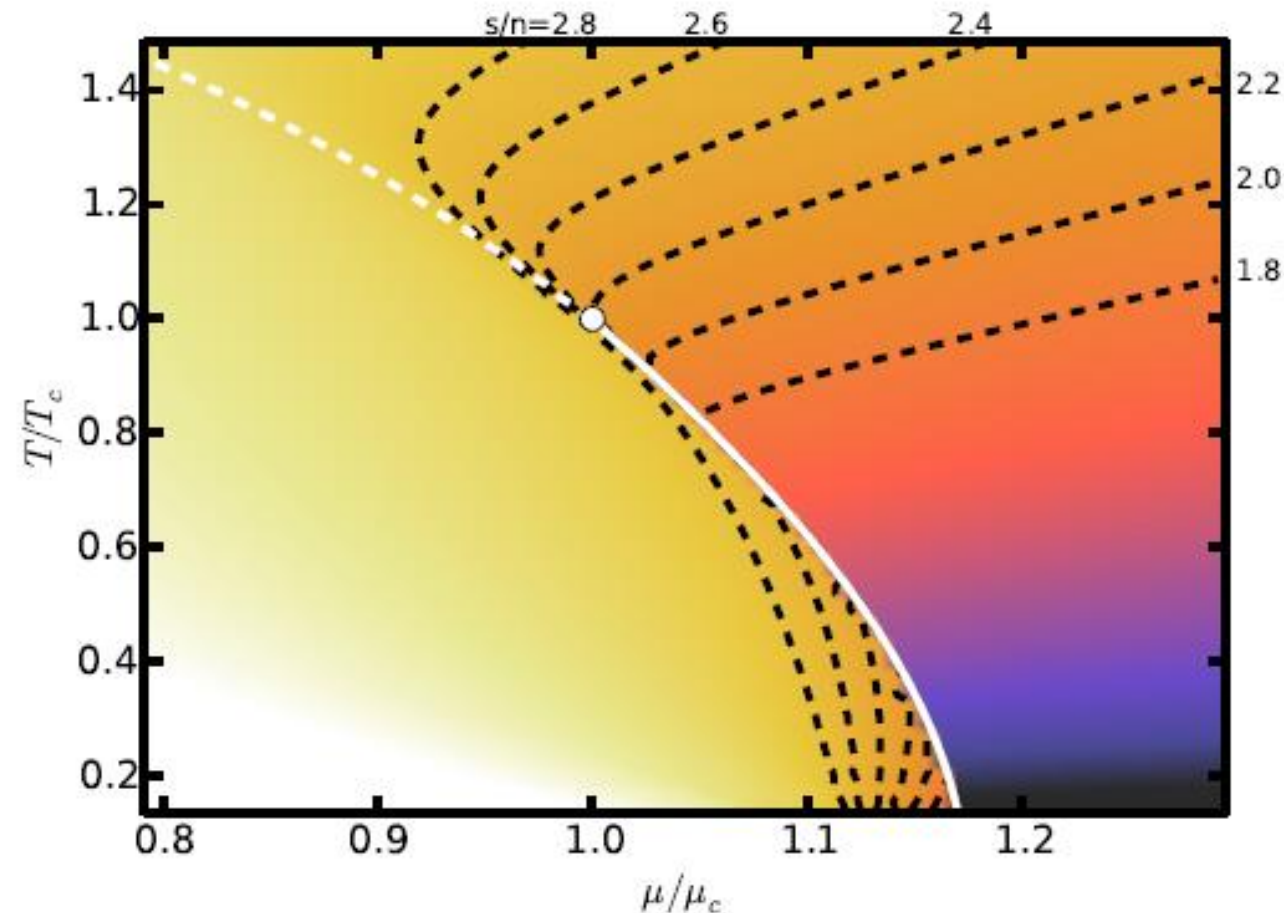
$$\mathcal{L}_{L\sigma M} = \bar{q}(i\partial - g(\sigma + i\gamma_5\tau\pi))q - \mathcal{L}_{km} - U(\sigma, \pi),$$

one alternative:  
PNJL: J.M. Torres-Rincon

$$U(\sigma, \pi) = \frac{\lambda}{4}(\sigma^2 + \pi^2 - \zeta)^2 - H\sigma,$$

$$\mathcal{L}_{km} = \frac{1}{2}(\partial_\mu\sigma\partial^\mu\sigma + \partial_\mu\pi\partial^\mu\pi),$$

CEP discussion:  
Scavenius, Mocsy, Mishustin, Rischke, PRC (2001)  
Schaefer, Wambach, NPA (2005)  
Schaefer, Pawłowski, Wambach, PRD (2007)  
Herbst, Pawłowski, Schaefer PLB (2011)



$$\Omega_{\text{MFA}} = \Omega_{qq} \Big|_{m_q=gv} + U(v, 0),$$

↑ ↑  
sigma pi

$$m_q^* = gv$$

$$m_\sigma^{*2} = \frac{\partial^2 \Omega_{qq} \Big|_{m_q=gv}}{\partial v^2} + \lambda(3v^2 - \zeta),$$

$$m_\pi^{*2} = 2g^2 \frac{\partial \Omega_{qq}}{\partial m_q^2} \Big|_{m_q=gv} + \lambda(v^2 - \zeta),$$

$$0 = \frac{\partial \Omega_{qq} \Big|_{m_q=gv}}{\partial v} + \lambda(v^3 - v\zeta) - H,$$

$$\begin{aligned} \Omega_{\text{LFA}} = & \langle U(v + \Delta, \boldsymbol{\pi}) \rangle + \langle \Omega_{qq}(m_q) \rangle \\ & - \frac{1}{2} m_\sigma^2 \langle \Delta^2 \rangle - \frac{1}{2} m_\pi^2 \langle \boldsymbol{\pi}^2 \rangle + \Omega_\pi + \Omega_\sigma, \end{aligned}$$

$$m_q^* = g \left\langle \sqrt{\sigma^2 + \boldsymbol{\pi}^2} \right\rangle,$$

$$m_\sigma^{*2} = \left\langle \frac{\partial^2 \Omega_{qq}}{\partial \Delta^2} \right\rangle + \lambda(3v^2 + 3\langle \Delta^2 \rangle + \langle \boldsymbol{\pi}^2 \rangle - \zeta),$$

$$m_\pi^{*2} = \left\langle \frac{\partial^2 \Omega_{qq}}{\partial \pi_a^2} \right\rangle + \lambda(v^2 + \langle \Delta^2 \rangle + \frac{5}{3} \langle \boldsymbol{\pi}^2 \rangle - \zeta),$$

$$0 = \lambda v(v^2 + 3\langle \Delta^2 \rangle + \langle \boldsymbol{\pi}^2 \rangle - \zeta) - H,$$

$$\langle \Delta^2 \rangle, \langle \pi_a^2 \rangle = 2\partial \Omega_{\sigma, \pi} / \partial (m_{\sigma, \pi}^2).$$

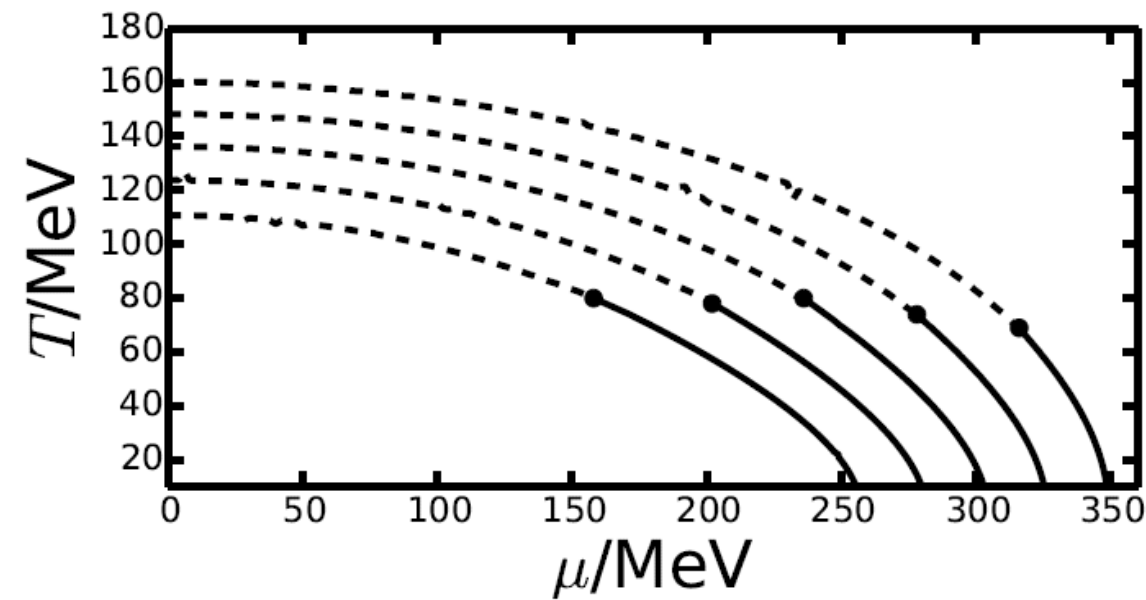
Mocsy, Mishustin, Ellis, PRC (2004)

Bowman, Kapusta, PRC (2009)

Ferroni, Koch, Pinto, PRC (2010)

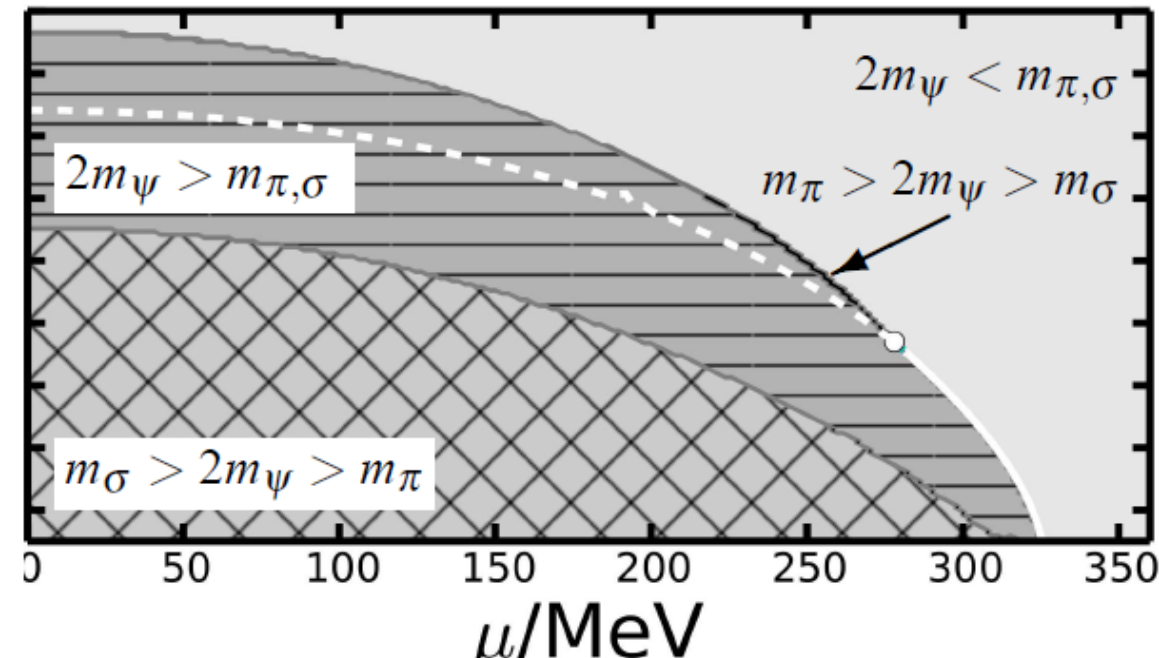
parameters:  $g = 3.387$ ,  $\xi = 7874 \text{ MeV}^2$ ,  $\lambda = 27.8$ ,  $H = 1760000 \text{ MeV}^3$

(from vacuum values of  $m_q$ ,  $m_\pi$ ,  $m_\sigma$ ,  $f_\pi$ )

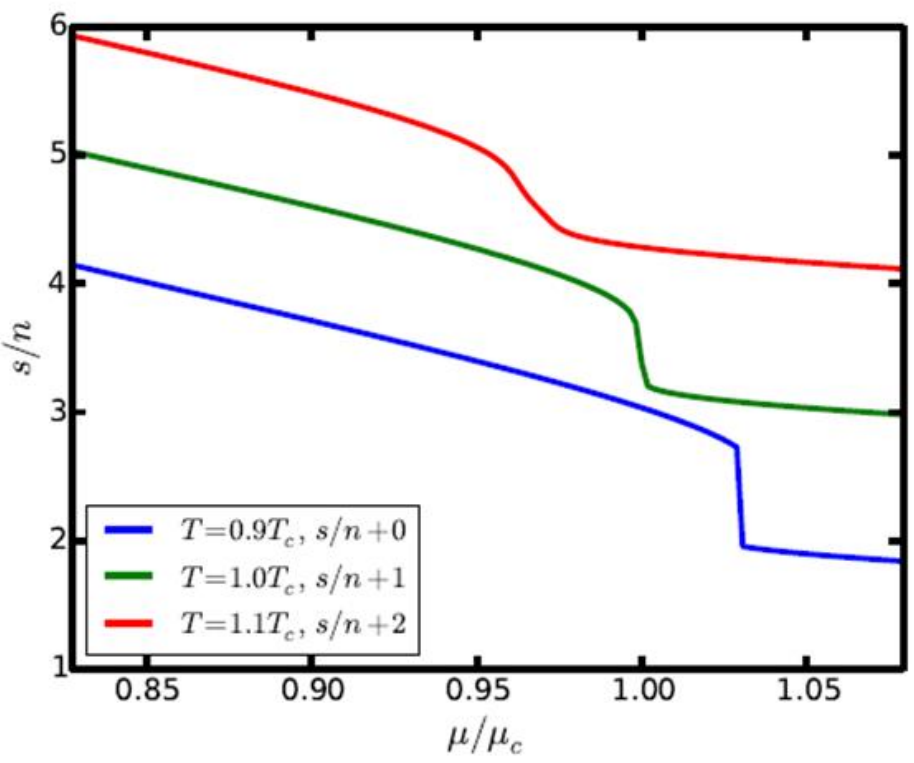
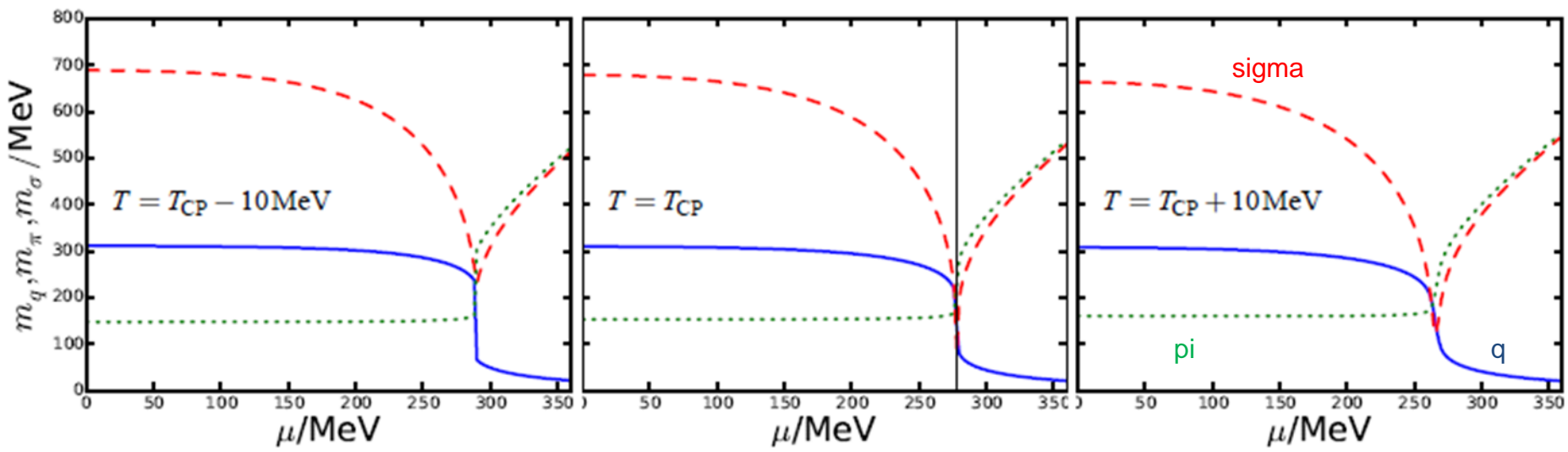


$m_s$  from 400 MeV (top)  
to 800 MeV (bottom)

note: improper CEP  
proper 1st order transition



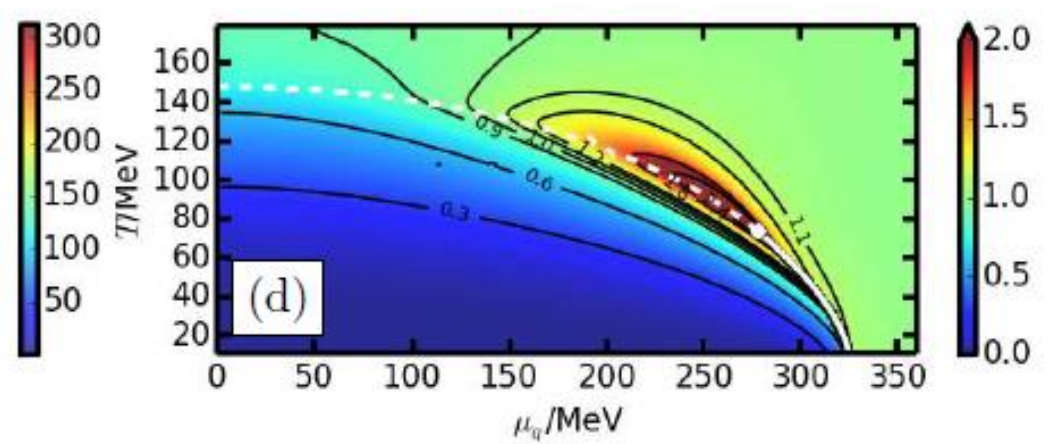
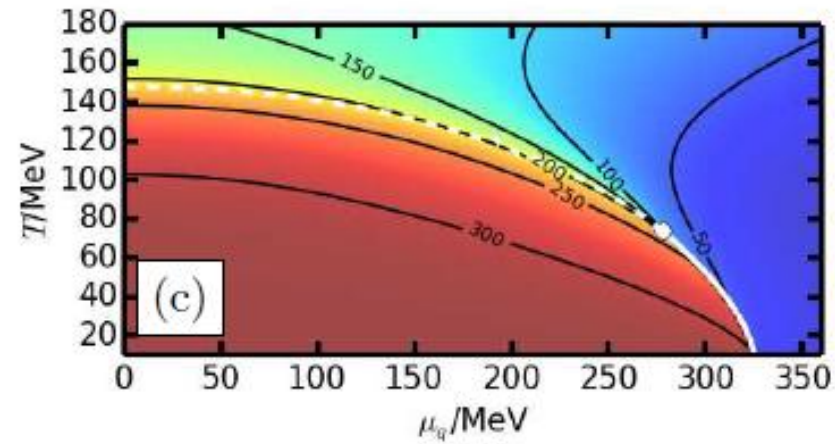
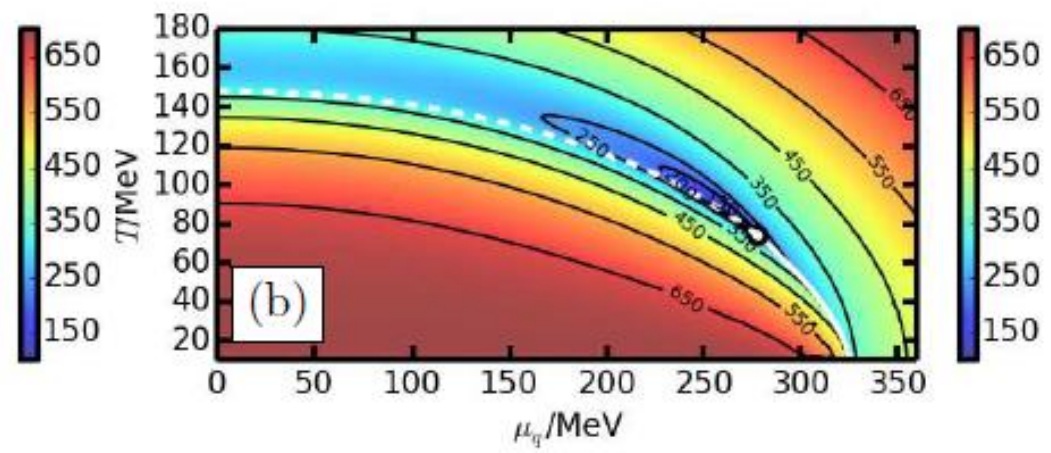
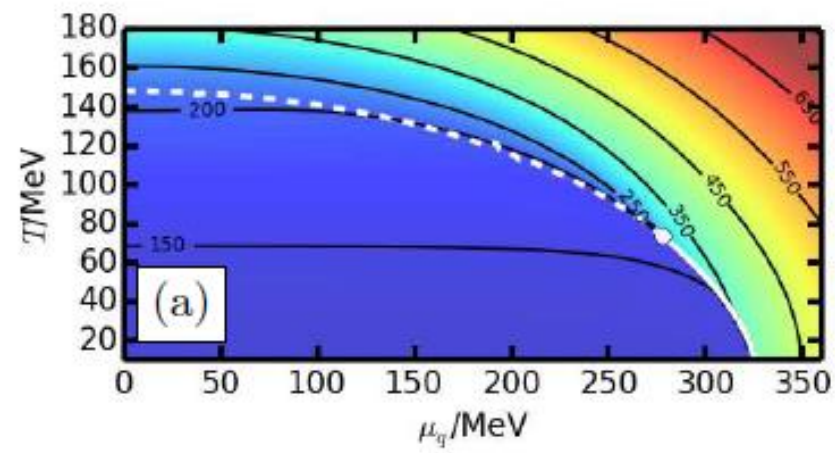
pattern of mass relations  
changes over the phase plane



disclaimer:

- at  $T = 0$  too small pressure (no proper baryons)
- LSM is of gas-liquid type





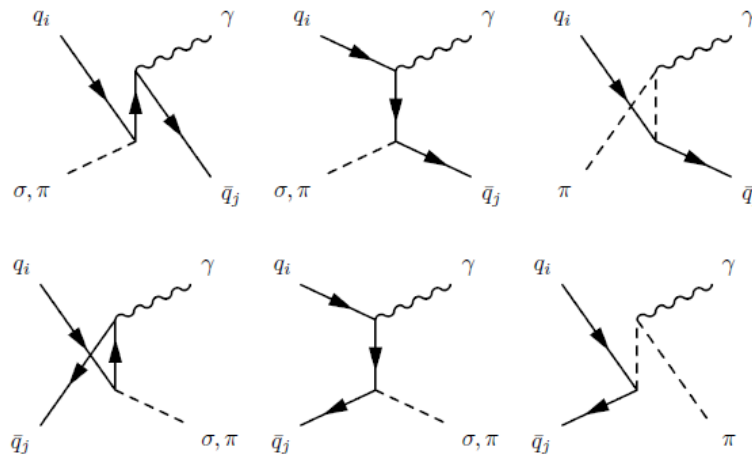
(a) pion mass    (b) sigma mass  
 (c) quark mass    (d) quark susceptibility (norm.)

wishful thinking:  $m(T, \mu)$  reflect phase diagram

# coupling in photons

$$\mathcal{L}_{\gamma\text{L}\sigma\text{M}} = \mathcal{L}_{\text{L}\sigma\text{M}} + \mathcal{L}_{\gamma} + \mathcal{L}_{\text{int}},$$

$$\mathcal{L}_{\text{int}} = -eQ_f\bar{\psi}A\psi + \frac{1}{2}e^2\pi^+\pi^-A^\nu A_\nu + \frac{1}{2}eA_\nu(\pi^-\partial^\nu\pi^+ + \pi^+\partial^\nu\pi^-),$$



NJL analog:

Fukushima, Ruggieri, Gatto, PRD (2010)

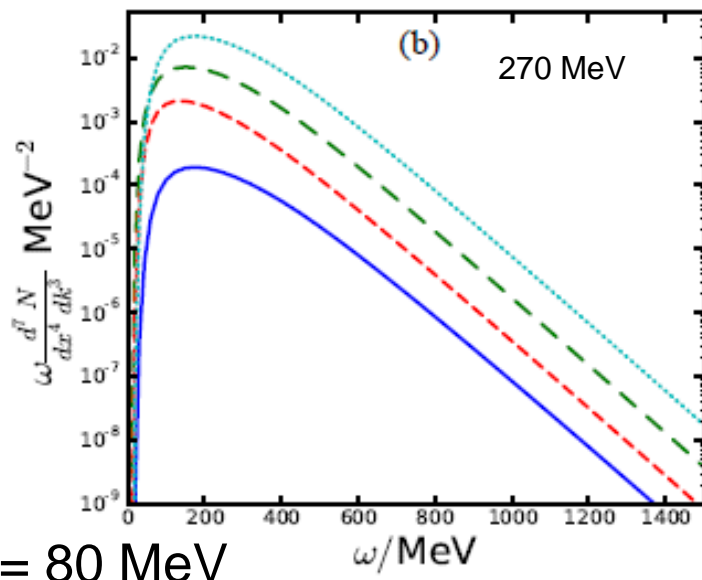
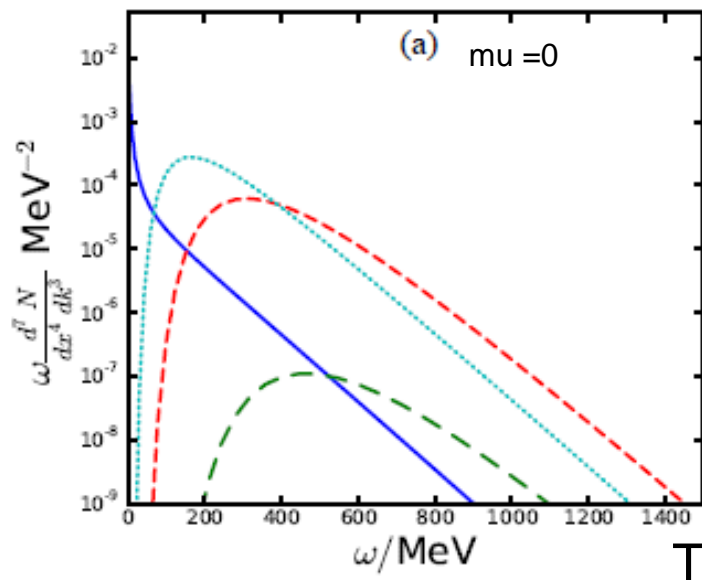
LsM:

Mizher, Chernodub, Fraga, PRD (2010)

$$\begin{aligned} q_i + \sigma, \pi &\rightarrow q_j + \gamma && \text{(Compton scatterings off quarks),} \\ \bar{q}_i + \sigma, \pi &\rightarrow \bar{q}_j + \gamma && \text{(Compton scat. off antiquarks),} \\ q_i + \bar{q}_j &\rightarrow \sigma, \pi + \gamma && \text{(annihilations)} \end{aligned}$$

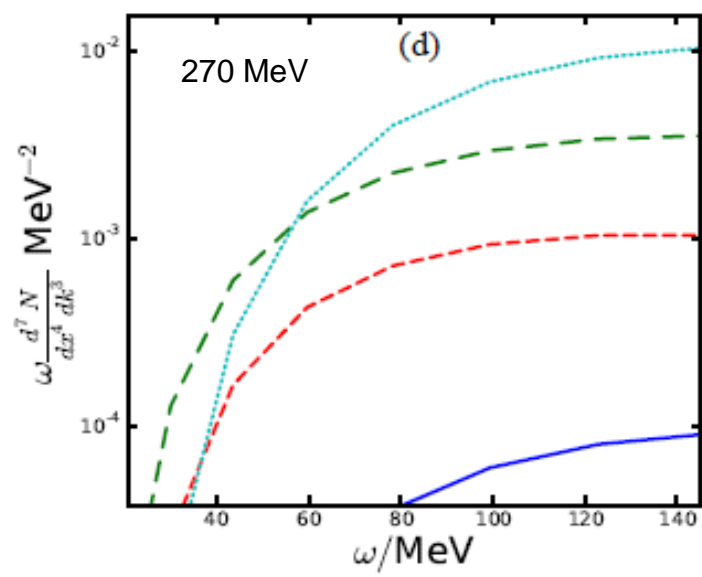
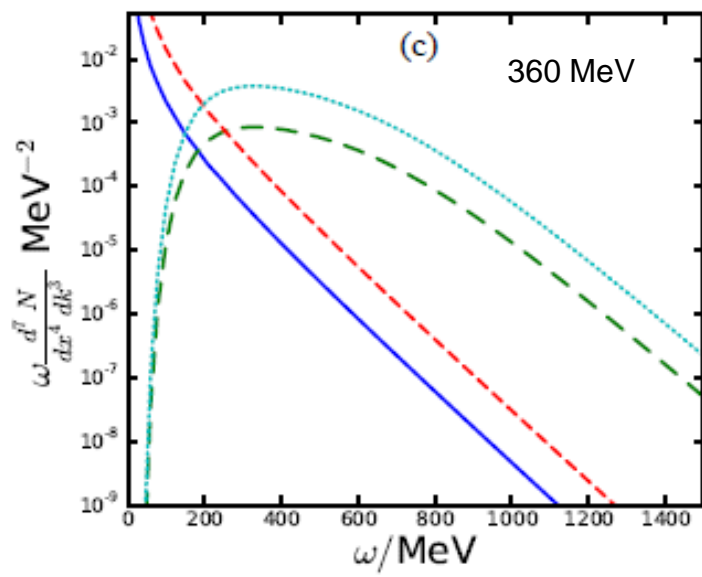
checked by  
FeynCal

$$\begin{aligned} \omega \frac{d^7 N_{12 \rightarrow 3\gamma}}{dx^4 dk^3} &= C \int \frac{d^3 p_1}{2p_1^0} \frac{d^3 p_2}{2p_2^0} \frac{d^3 p_3}{2p_3^0} \delta(p_1 + p_2 - p_3 - k) \\ &\times |\mathcal{M}_{12 \rightarrow 3\gamma}|^2 f_1(p_1) f_2(p_2) (1 \pm f_3(p_3)). \end{aligned}$$

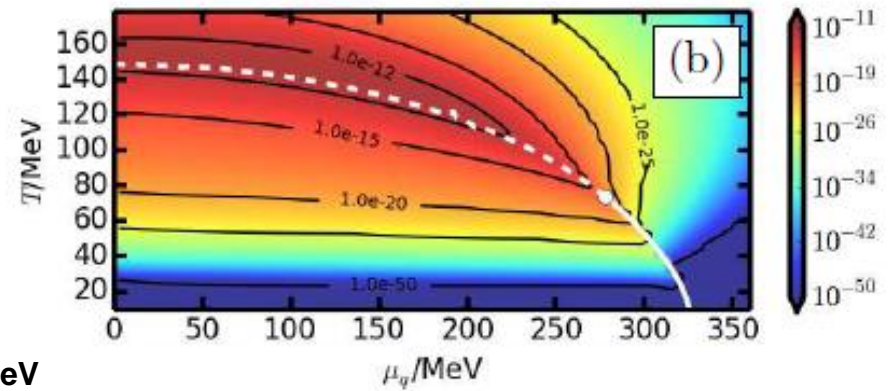
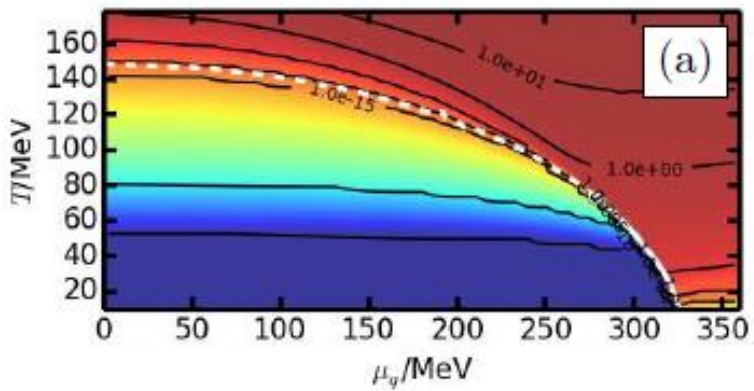


$q + \pi \rightarrow q + \text{gam}$   
 $q + \text{sigma} \rightarrow q + \text{gam}$   
 $q + q_{\text{bar}} \rightarrow \pi + \text{gam}$   
 $q + q_{\text{bar}} \rightarrow \text{sigma} + \text{gam}$

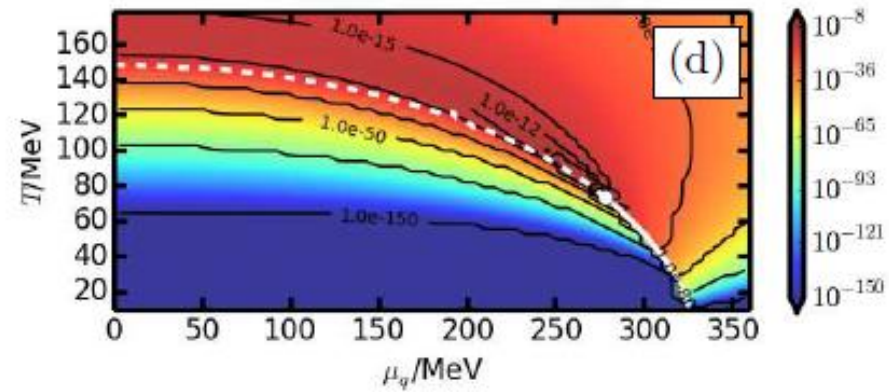
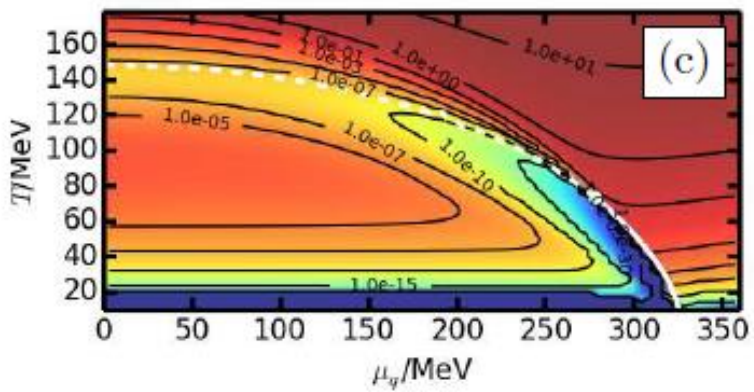
$T = 80 \text{ MeV}$



$$w d N / d^4 x d^3 k \text{ MeV}^2$$

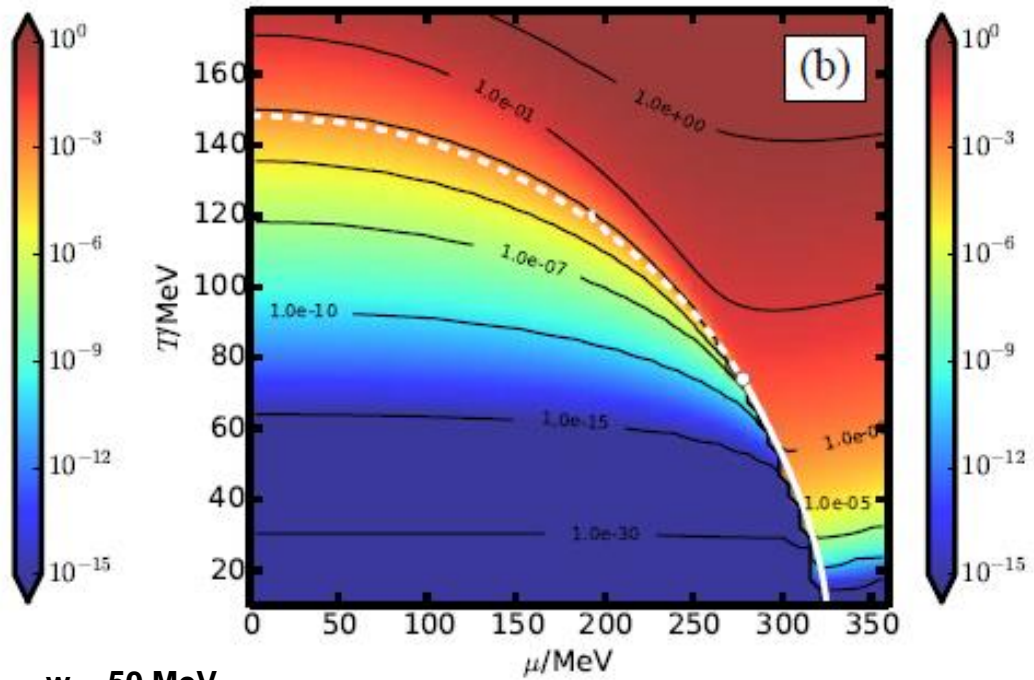
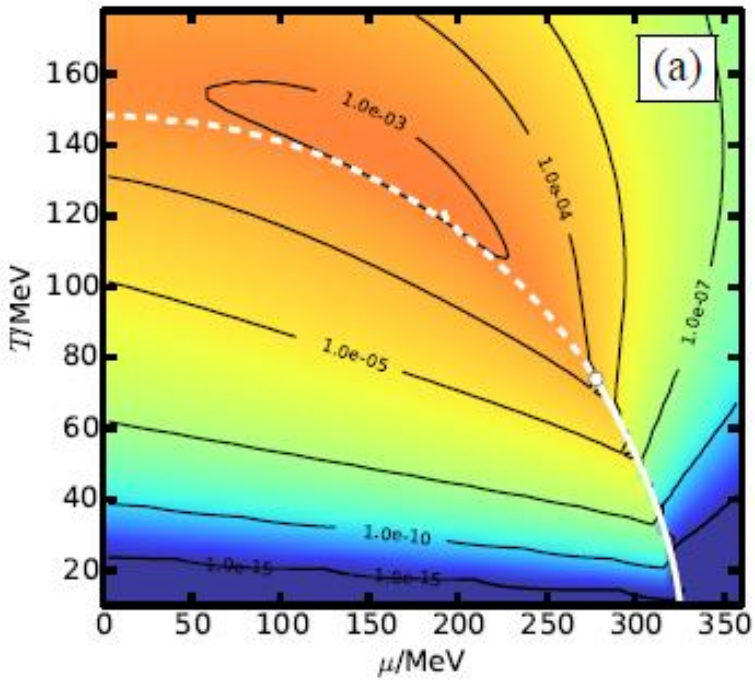


**w = 10 MeV**

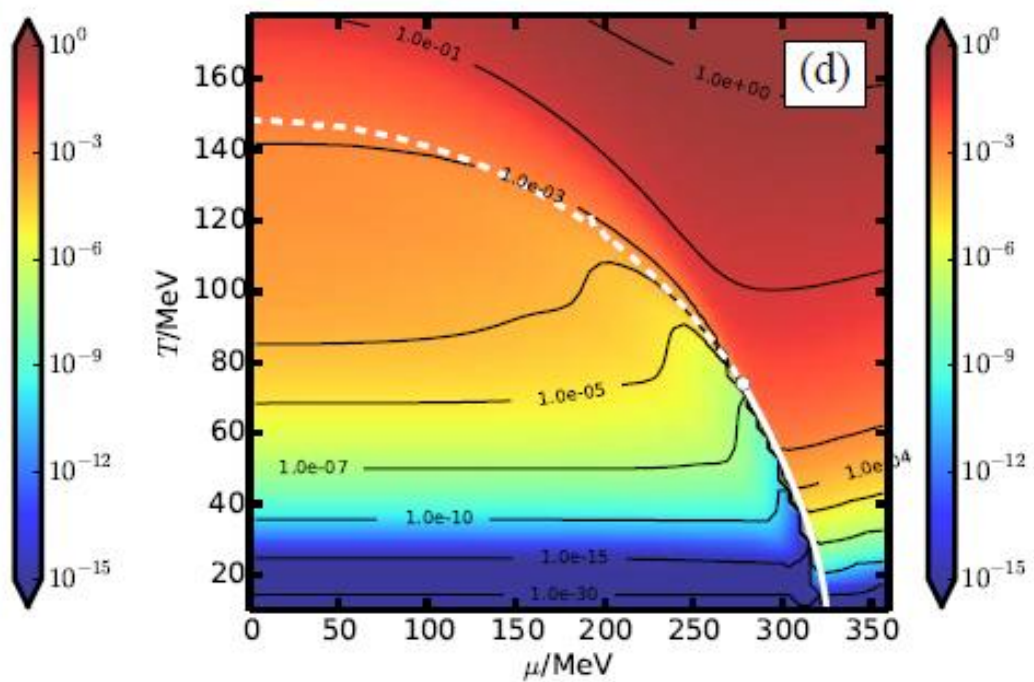
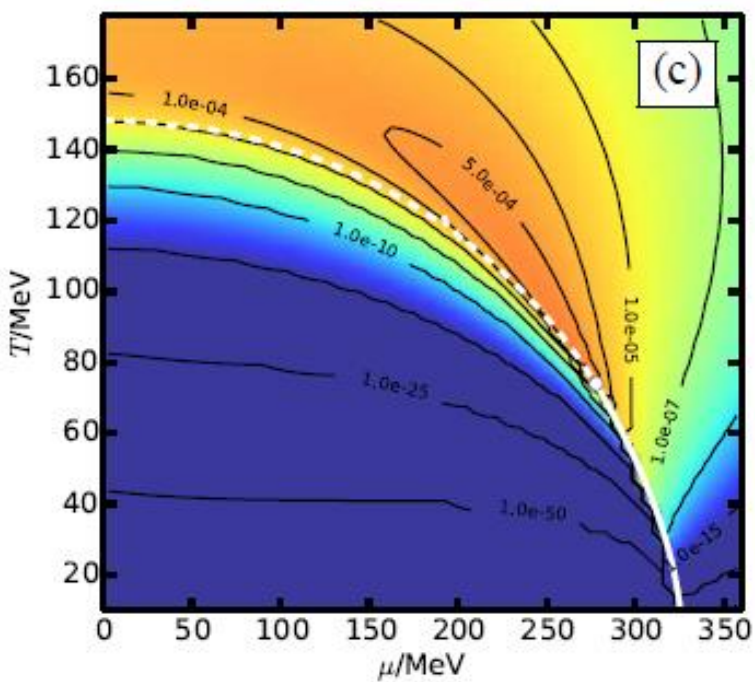


(a)  $q + q_{\text{bar}} \rightarrow \pi + \gamma$   
 (c)  $q + q_{\text{bar}} \rightarrow \sigma + \gamma$

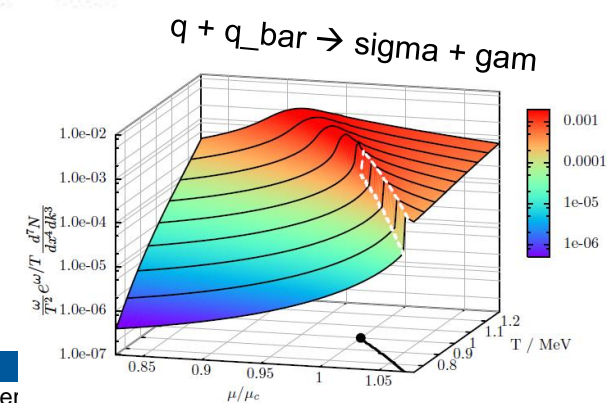
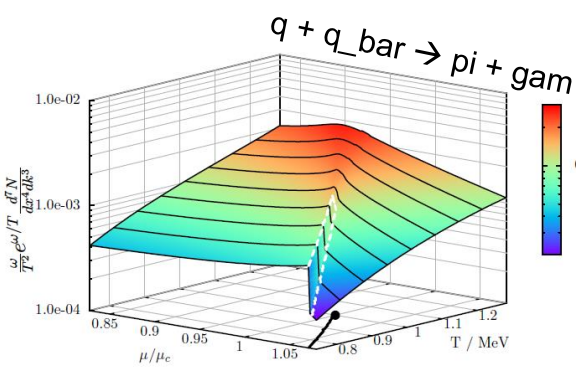
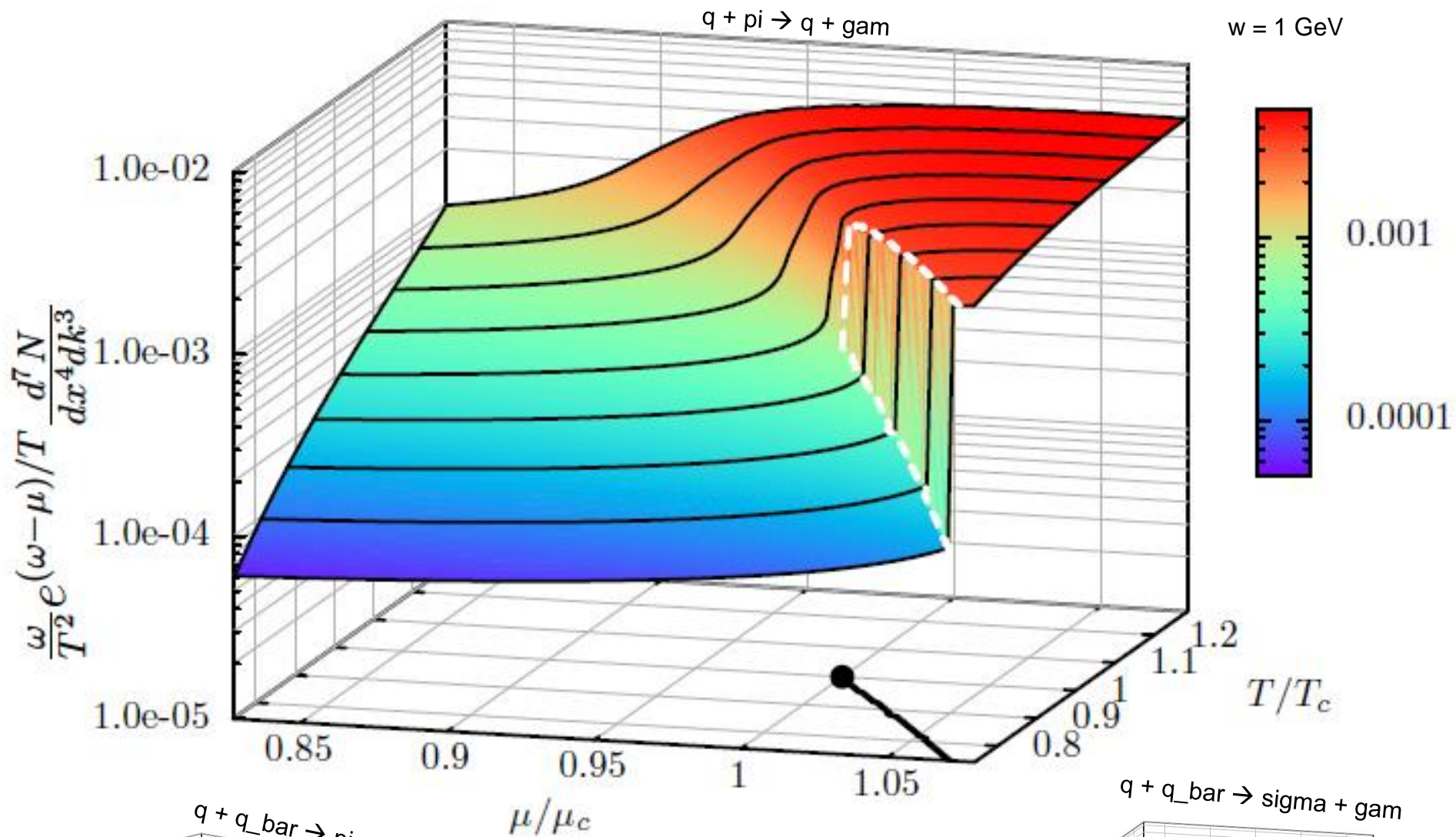
(b)  $q + \pi \rightarrow q + \gamma$   
 (d)  $q + \sigma \rightarrow q + \gamma$



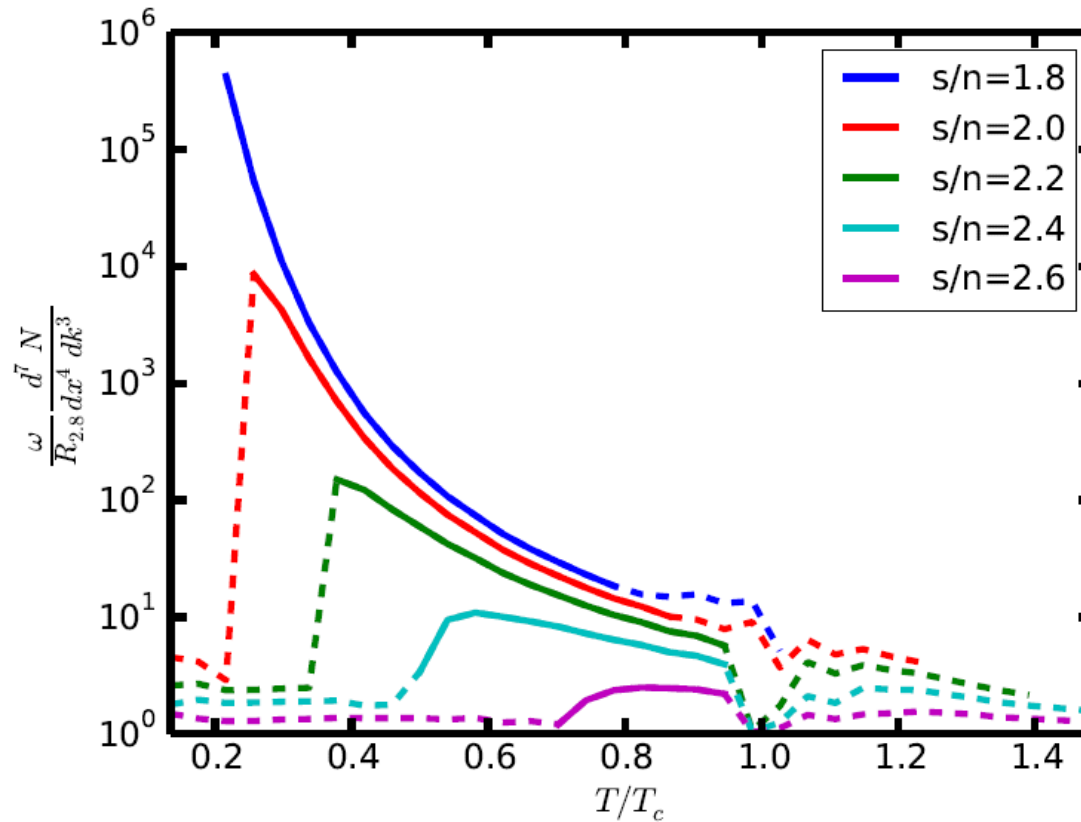
$w = 50 \text{ MeV}$



$\cdot \sigma + \lambda \leftrightarrow \psi + \bar{\psi} : (p) ; \psi + \lambda \leftrightarrow \sigma + \psi : (q) ; \psi + \lambda \leftrightarrow \psi + \psi : (r) ; \psi + \lambda \leftrightarrow \psi + \psi : (s)$

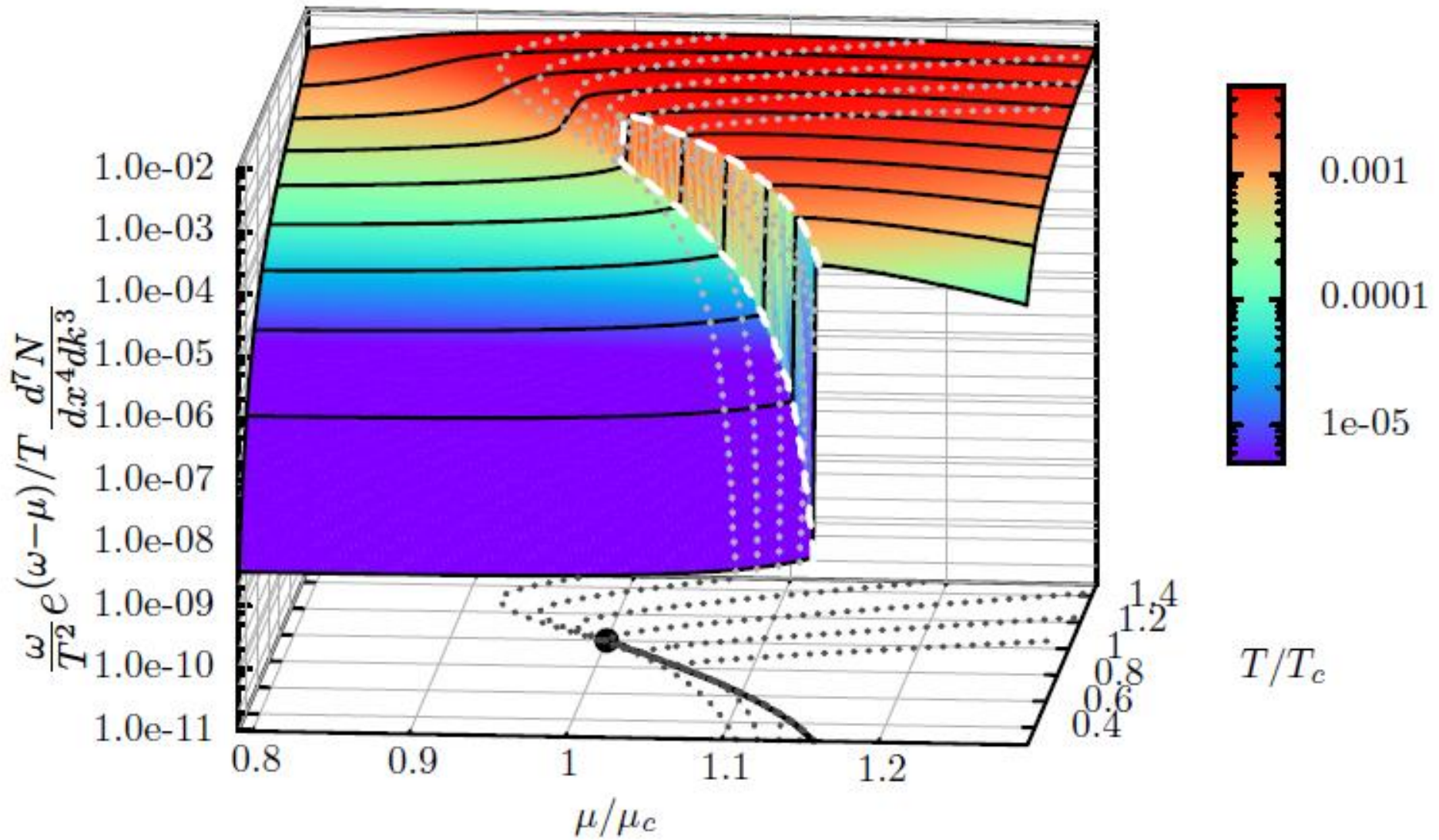


adiabatic expansion:  $s/n = \text{const}$



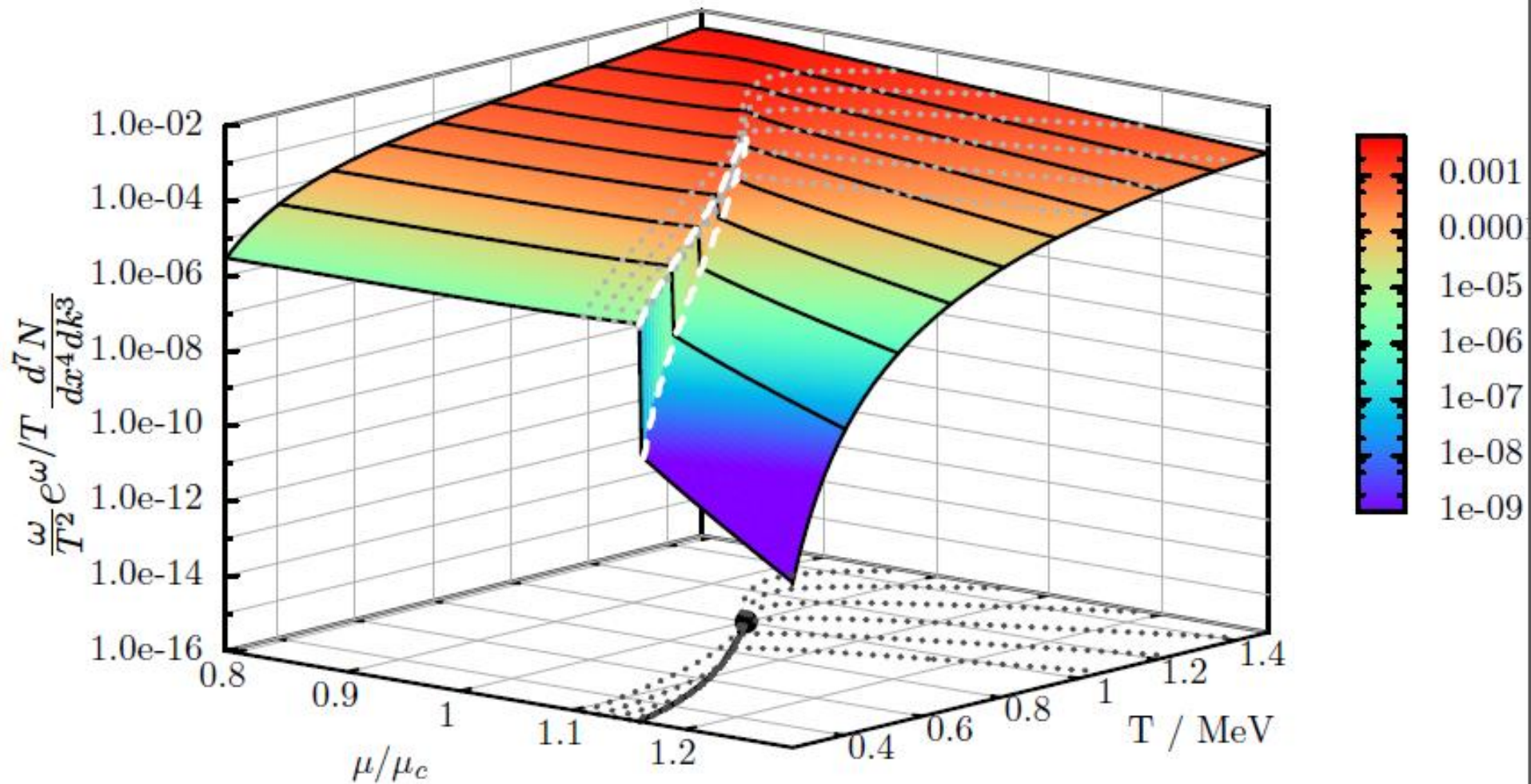
reference:  $s/n = 2.8$  (does not touch CEP and phase border curve)

$$qp \rightarrow gq, \quad \omega = 1000\text{MeV}$$





$qq \rightarrow gp, \quad \omega = 1000\text{MeV}$

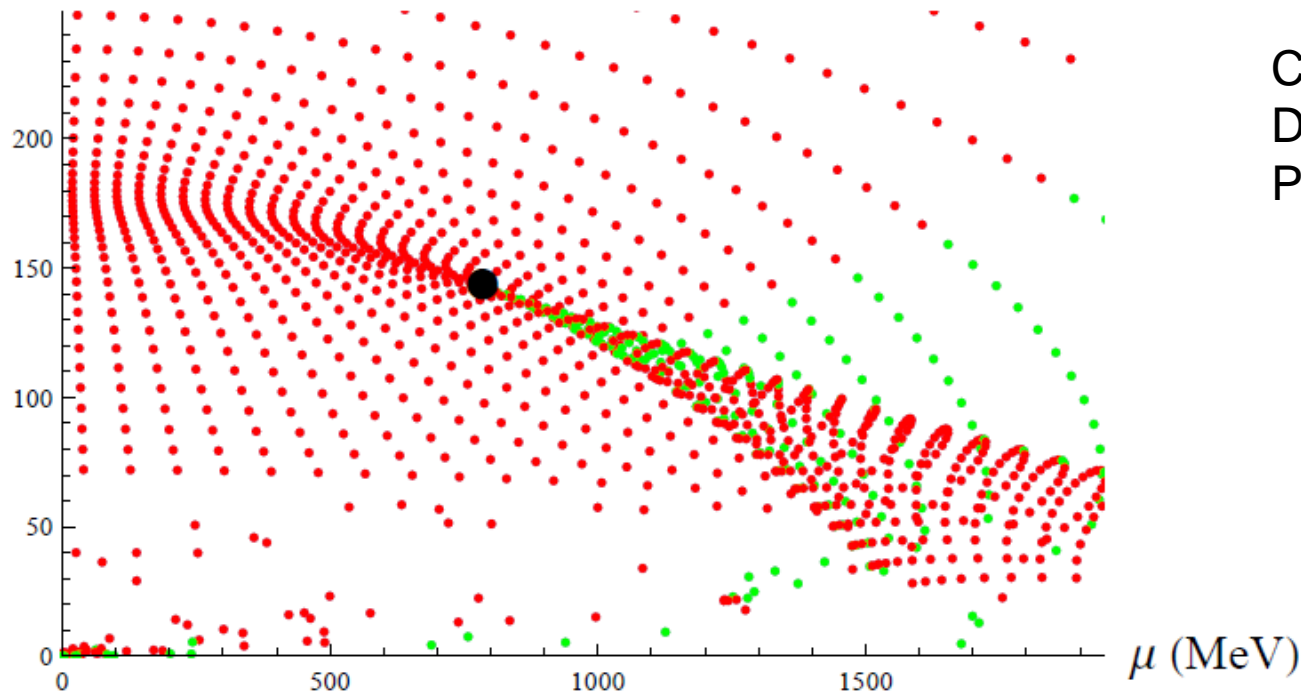


two effects: - rate(s) over  $T - \mu$  plane reflect  $m(T, \mu)$   
 - adiabatic paths affected by phase border curve

we leave out the improper CEP (reminder: CEP  $\rightarrow$  critical opalescence)

# holographic avenues

T (MeV)



Criticality at CEP:  
DeWolfe, Gubser, Rosen  
PRD (2011)

Production of Prompt Photons: Holographic Duality and Thermalization

Baier, Stricker, Taanila, Vuorinen, Phys.Rev. D86 (2012) 081901

Holographic Dilepton Production in a Thermalizing Plasma

Baier, Stricker, Taanila, Vuorinen, JHEP 1207 (2012) 094

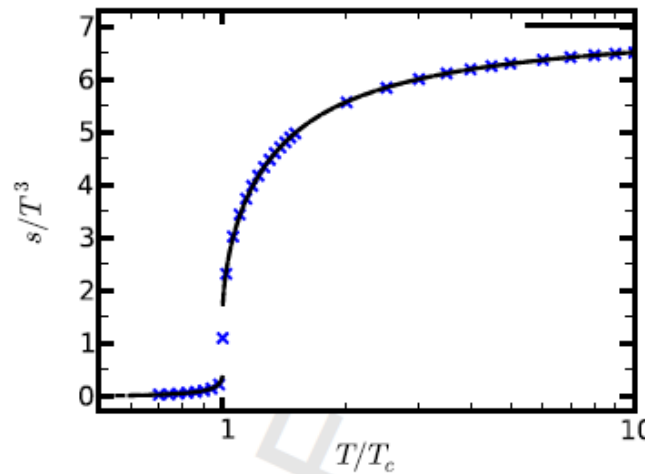
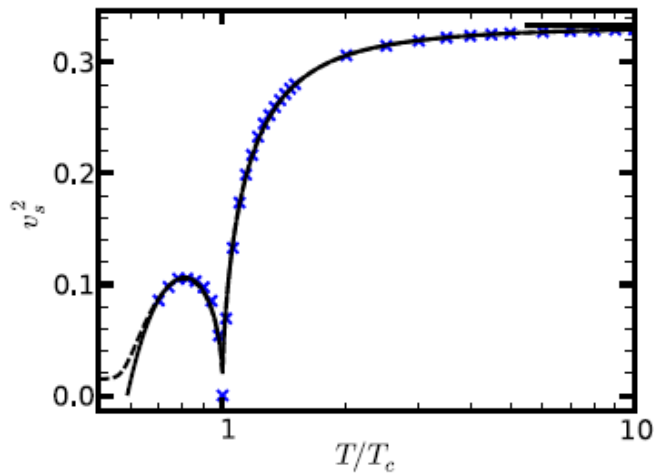
# 2. Viscosities: Holographic Input for midly NeD in HICs

## 1. Pure gluon medium: SU(3) – 1st order p.t.

exercise and model for early gluon-rich stage

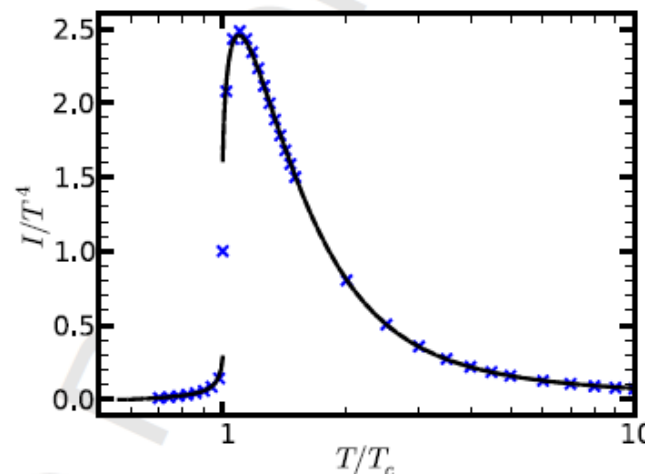
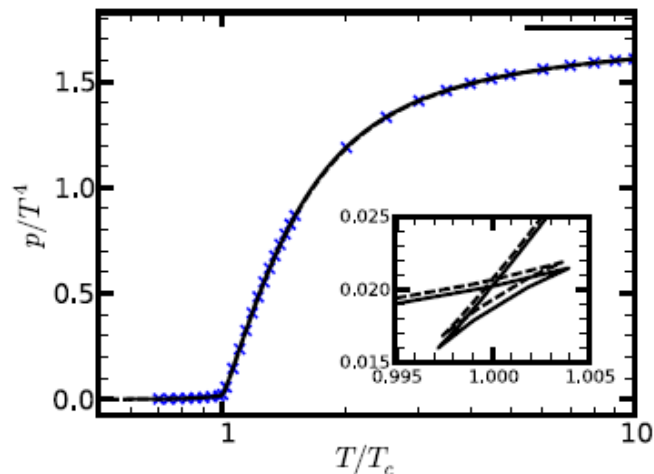
QPM: Chabrobay, Kapusta PRC (2012)  
Bluhm, BK, Redlich, PRC (2012)

holography: Kiritsis et al.  
Gubser et al.

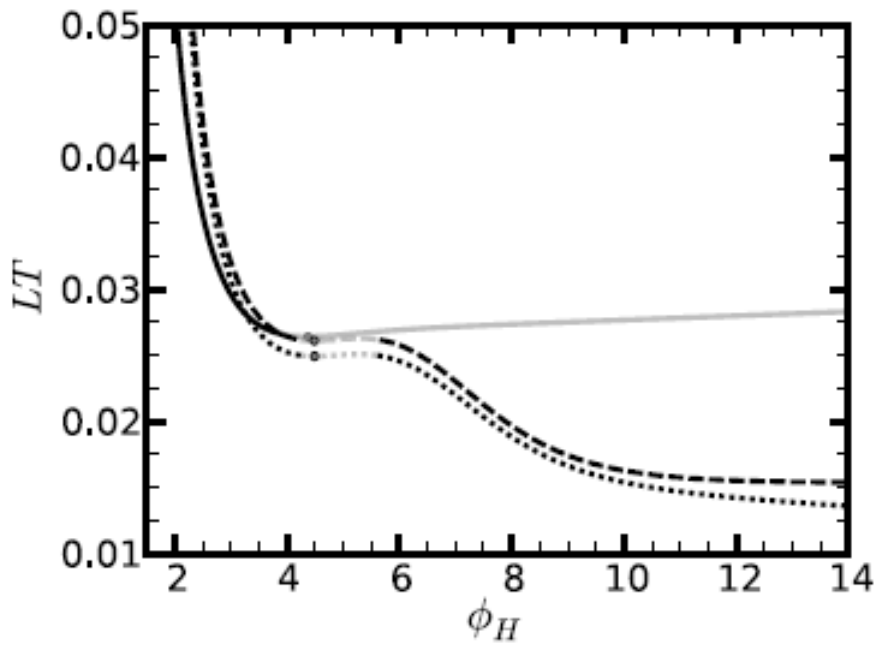
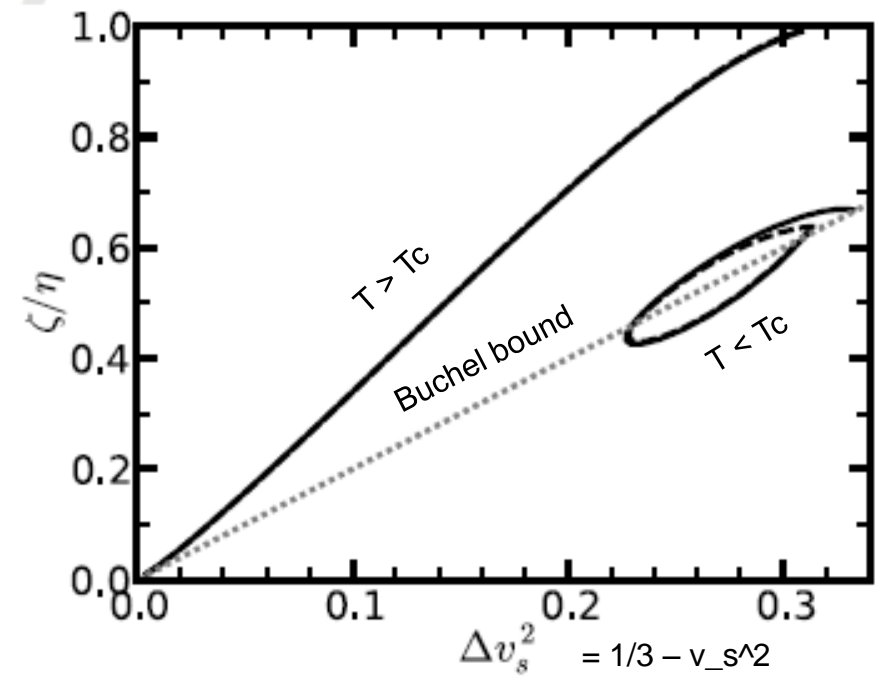
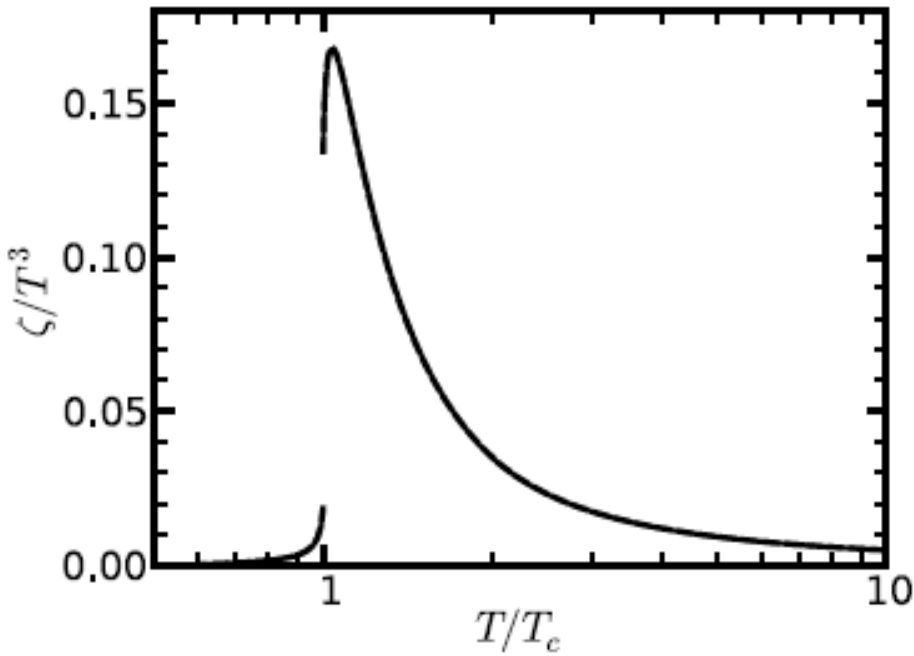


lattice data:  
Borsanyi et al., JHEP (2012)

consistent with  
Boyd et al., NPB (1996)

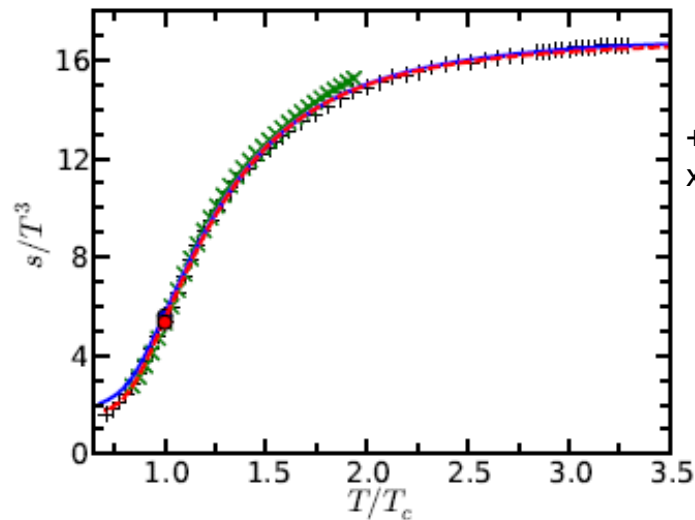
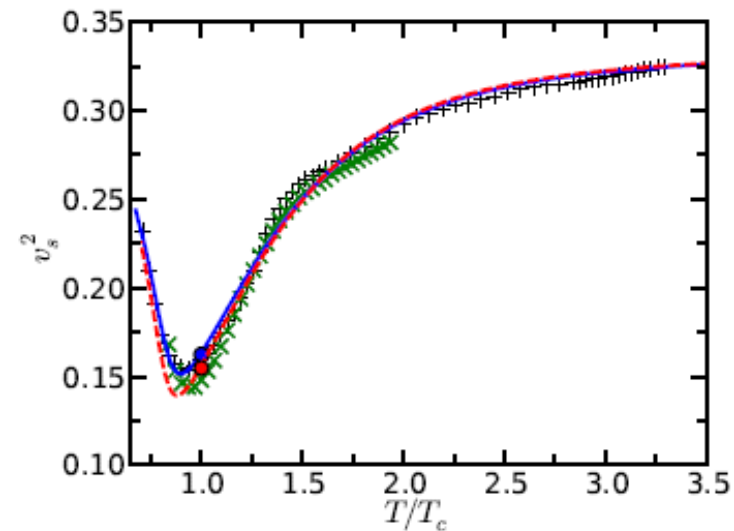


Yaresko, BK, PLB (2015)

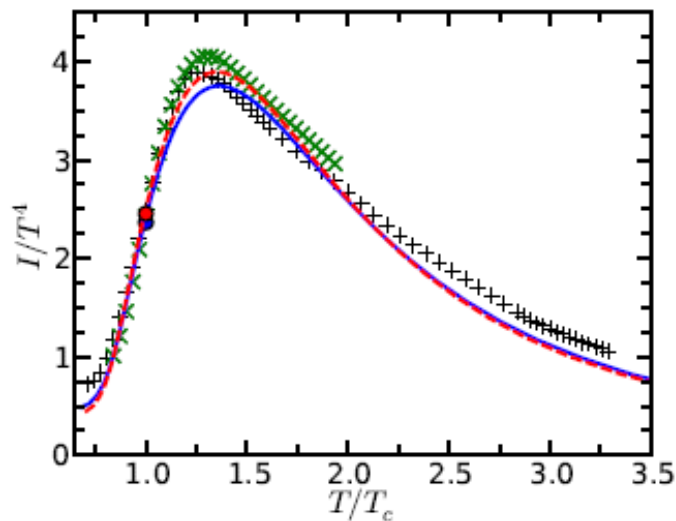
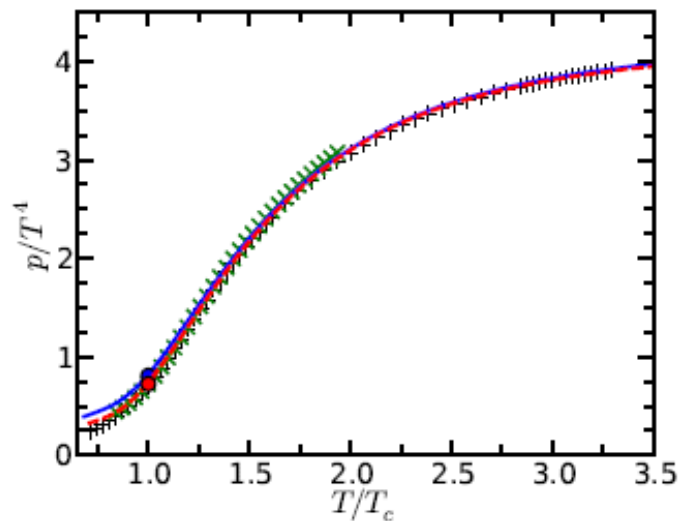


disclaimer: holographic model has  $\eta/s = 1/4 \pi$

## 2. QCD: cross over – $T_c \rightarrow T_{pc}$



lattice data:  
 +++: Borsanyi et al. PLB (2014)  
 xxx: Bazavov et al., PRD (2014)



Yaresko, Knaute, BK  
 EPJC (2015)

# Holography

based on AdS/CFT correspondence



5D Riemann  
class. gravity + fields  
asyp. AdS + black hole  
metric  
dilaton

Hawking T  
Bekenstein-Hawking s

4D Minkowski  
QFT (operators)  
thermo-field theory  
energy-momentum tensor  
 $\langle (\text{gluon field})^2 \rangle$

tachyon

quark condensate

At  $\mu = 0$ , the equation of state, in parametric form, follows from [13]

$$LT(\phi_H) = \frac{V(\phi_H)}{\pi V(\phi_0)} \exp \left( A(\phi_0) + \int_{\phi_0}^{\phi_H} d\phi \left[ \frac{1}{4X} + \frac{2}{3}X \right] \right), \quad (1)$$

$$G_5 s(\phi_H) = \frac{1}{4} \exp \left( 3A(\phi_0) + \frac{3}{4} \int_{\phi_0}^{\phi_H} d\phi \frac{1}{X} \right), \quad (2)$$

for entropy density  $s$  and temperature  $T$ , where the scalar function  $X(\phi; \phi_H)$  [14] is determined by the system (a prime means a derivative w.r.t.  $\phi$ )

$$X' = - \left( 1 + Y - \frac{2}{3}X^2 \right) \left( 1 + \frac{3}{4X} \frac{V'}{V} \right), \quad (3)$$

$$Y' = - \left( 1 + Y - \frac{2}{3}X^2 \right) \frac{Y}{X}, \quad (4)$$

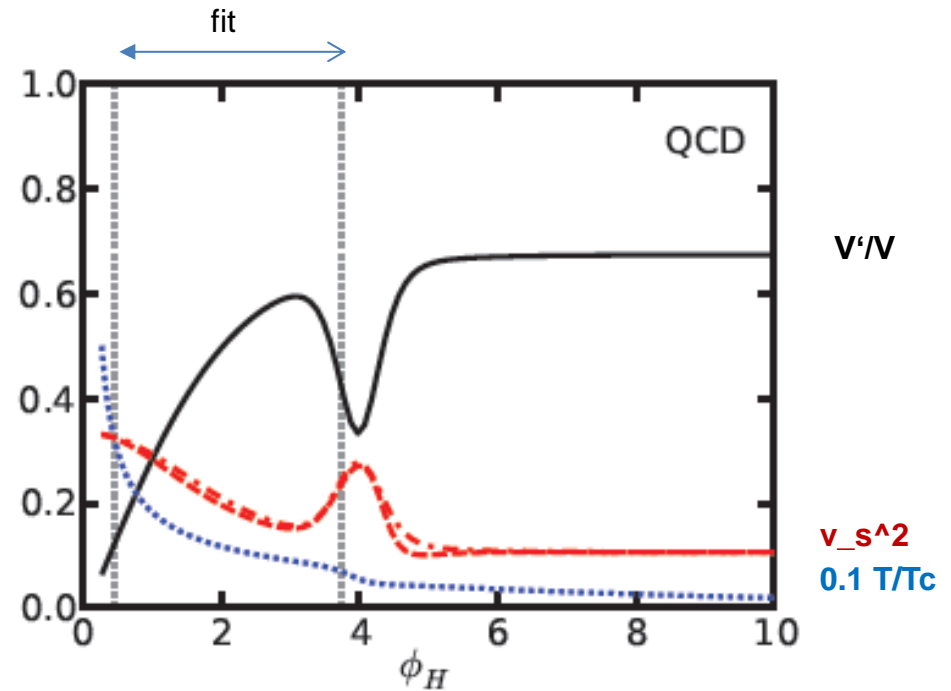
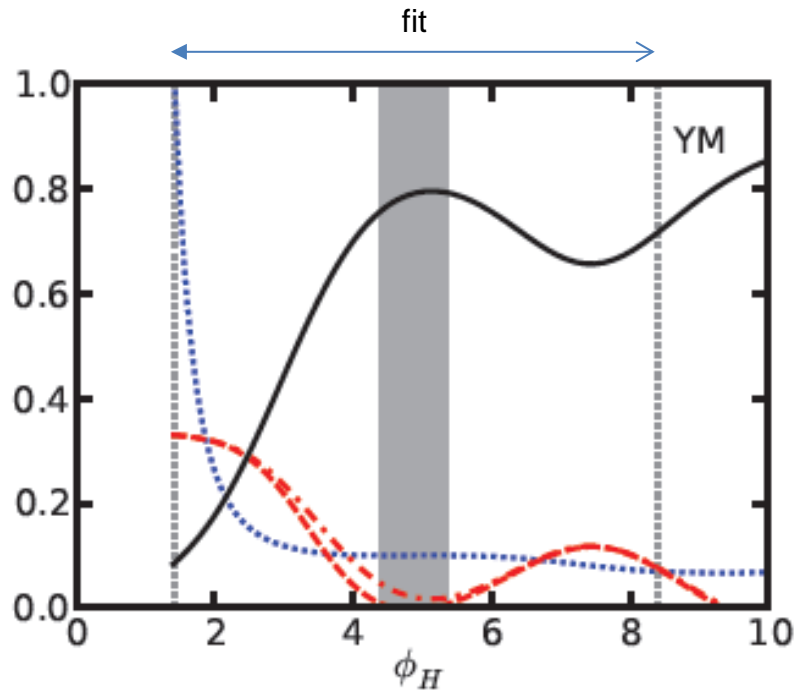
which is integrated from the horizon  $\phi_H - \epsilon$  to the boundary  $\phi_0$  with initial conditions

$$X(\phi_H - \epsilon) = - \frac{3}{4} \frac{V'(\phi_H)}{V(\phi_H)} + \mathcal{O}(\epsilon^1), \quad (5)$$

$$Y(\phi_H - \epsilon) = - \frac{X(\phi_H - \epsilon)}{\epsilon} + \mathcal{O}(\epsilon^0), \quad (6)$$

and  $\epsilon \rightarrow 0$ . The quantity  $A(\phi_0)$  encodes the near-boundary behavior of the model. We assume  $L^2 V(\phi) \approx -12 + \frac{L^2 M^2}{2} \phi^2$  for  $\phi \rightarrow \phi_0 = 0$  which results in  $A(\phi_0) = \frac{\log \phi_0}{\Delta - 4}$ , whereby we have set  $L\Lambda = 1$  [13] and, as usual,  $L^2 M^2 = \Delta(\Delta - 4)$ . We consider  $2 < \Delta < 4$ .

# adjusting potentials (phi self-interaction)



dream: lattice QCD thermodynamics  $\rightarrow V(\phi)$

outlook:  $\mu > 0$  & phase diagram a la deWolfe, Gubser, Rosen



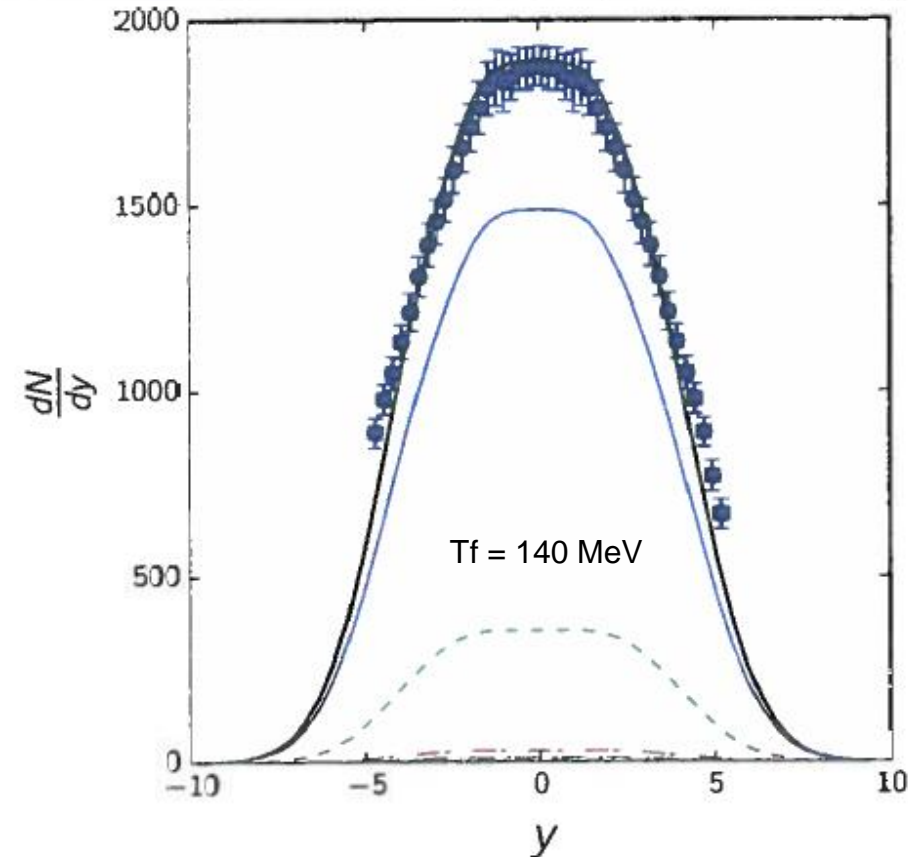
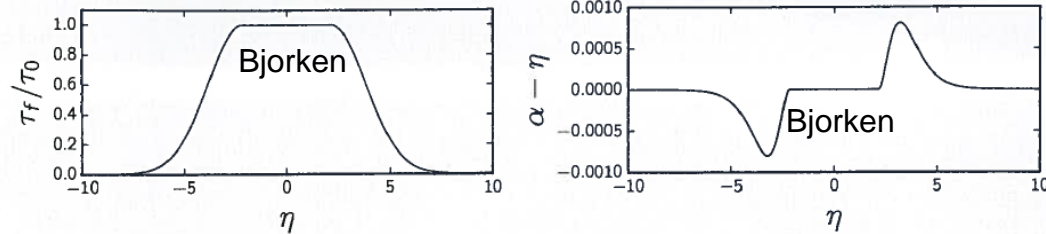
# 3. Longitudinal Dynamics

## (i) CF freeze-out

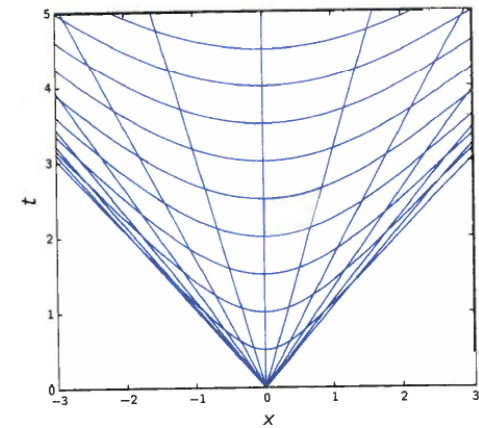
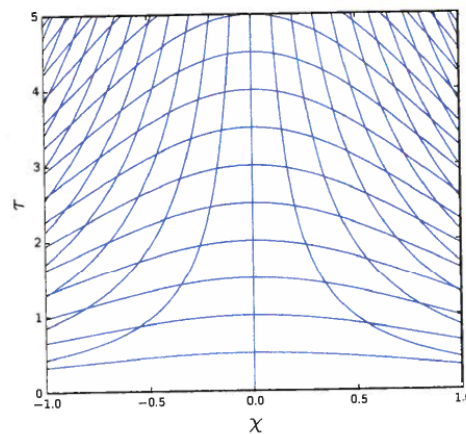
$$\frac{dN}{dy} = \frac{A_{\perp}}{(2\pi)^2} \int d\chi \frac{g(\chi, y)}{f(\chi, y)} e^{-mf(\chi, y)} \left[ m^2 + \frac{2m}{f(\chi, y)} + \frac{2}{f^2(\chi, y)} \right]$$

$$f(\chi, y) = \frac{1}{T_f} \cosh(y - \alpha(\chi)),$$

$$g(\chi, y) = \sinh(\chi - y) \partial_{\chi} \tau_f(\chi) + \tau_f(\chi) \cosh(\chi - y).$$



## Milne coordinates



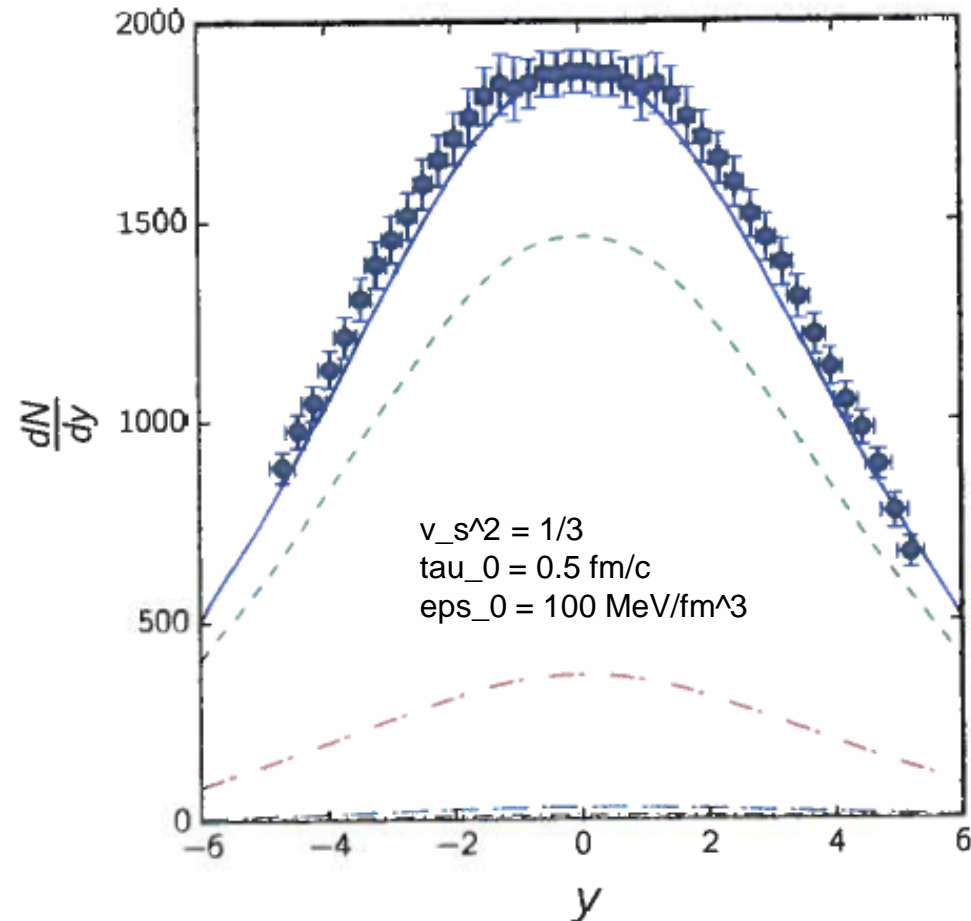
data: ALICE, PLB (2013) Pb + Pb, sqrt(s\_NN) = 2.75 TeV

## (ii) dynamics

initial conds.

$$\epsilon(\tau_0, \chi) = \begin{cases} \epsilon_0 & \text{für } |\chi| < a, \\ \epsilon_0 e^{-b(\chi - a \operatorname{sign}(\chi))^2} & \text{für } |\chi| \geq a, \end{cases}$$

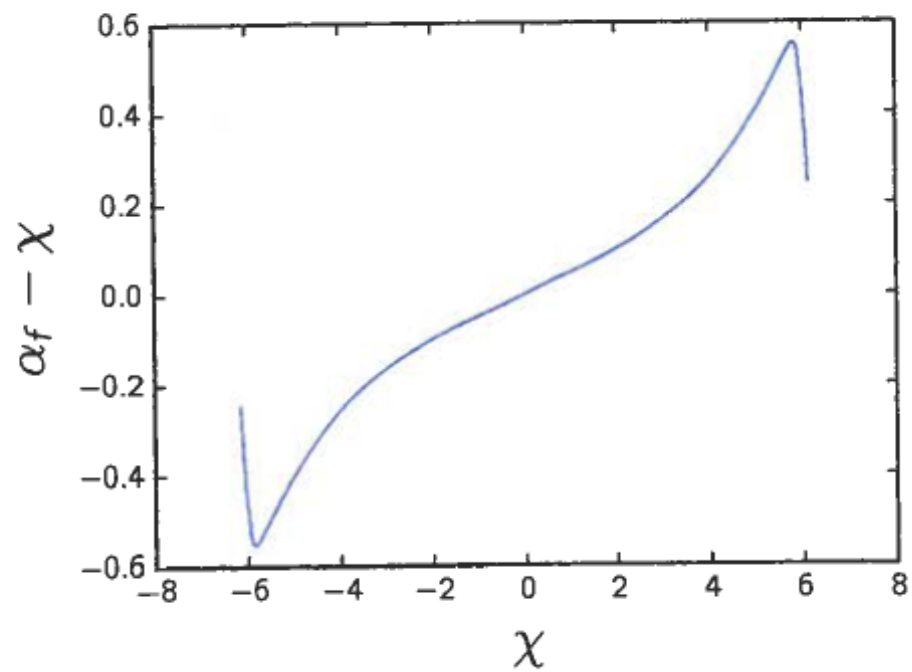
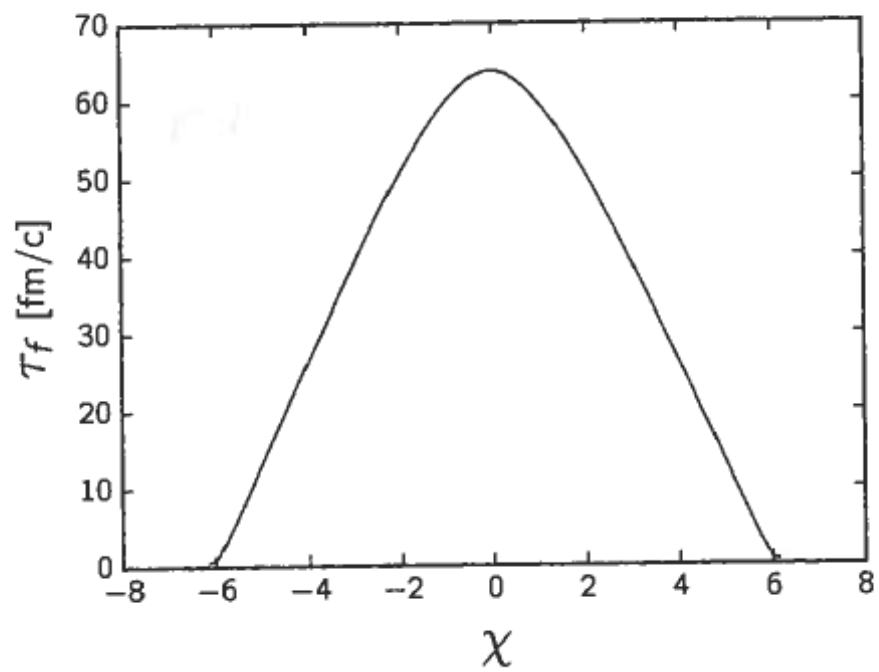
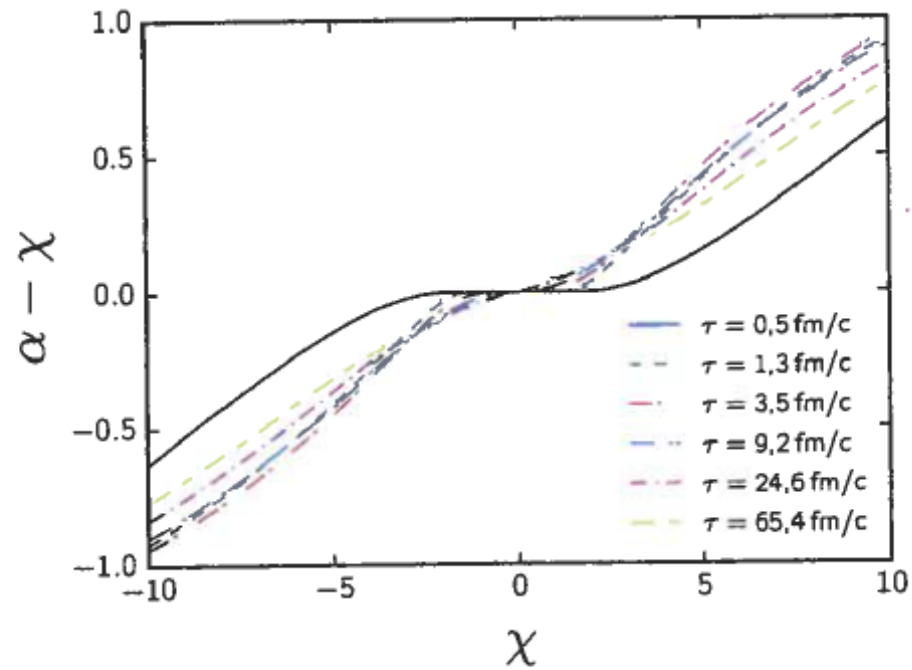
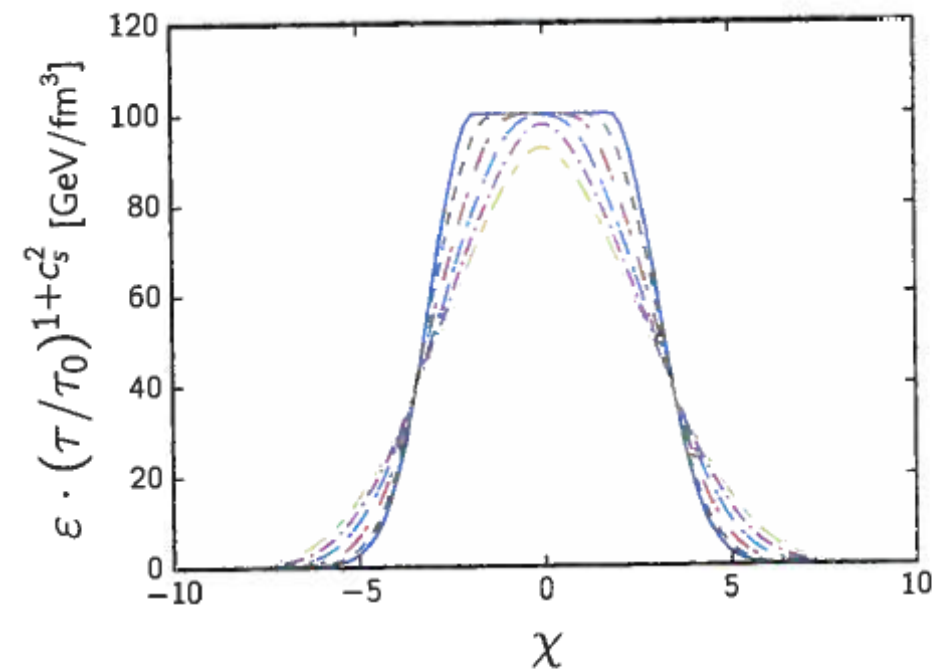
$$\alpha(\tau_0, \chi) = \begin{cases} \chi & \text{für } |\chi| < a, \\ \chi + b \operatorname{sign}(\chi) \left( \sqrt{0,1 \cdot (|\chi| - a)^2 + 1} - 1 \right) & \text{für } |\chi| \geq a, \end{cases}$$



reminder:

Bjorken flow:  $\epsilon(\tau)$

$$\alpha = \chi$$



next step: inverting a heavy-ion collision

a la Stephanov, Yin (2014)

reformulate hydro in  $T, \alpha$  coordinates  
→ Cauchy problem from freeze-out into past  
supposed (i) viscosities are small  
(ii) 1+1 dynamics is applicable

1 + 1 hydro (two 1st order pDEs) → Chalatnikov eq. (one 2nd order pDEs)

# Summary/Outlook

- photon emissivities: phase structure  $\rightarrow$  rates,  
adiabatic expansion trajectories  
+ phase mixture
- 1-dilaton holography: lattice QCD  $\rightarrow$   $V(\phi)$   $\rightarrow$  viscosities,  
improvements: 2-field model with  
chiral condensate, non-zero  $\mu$
- $dN/dy$  (ALICE)  $\rightarrow$  f.o. hypersurface  
improvements: more dynamics + EoS