









Electromagnetic emissivity of hot and dense matter

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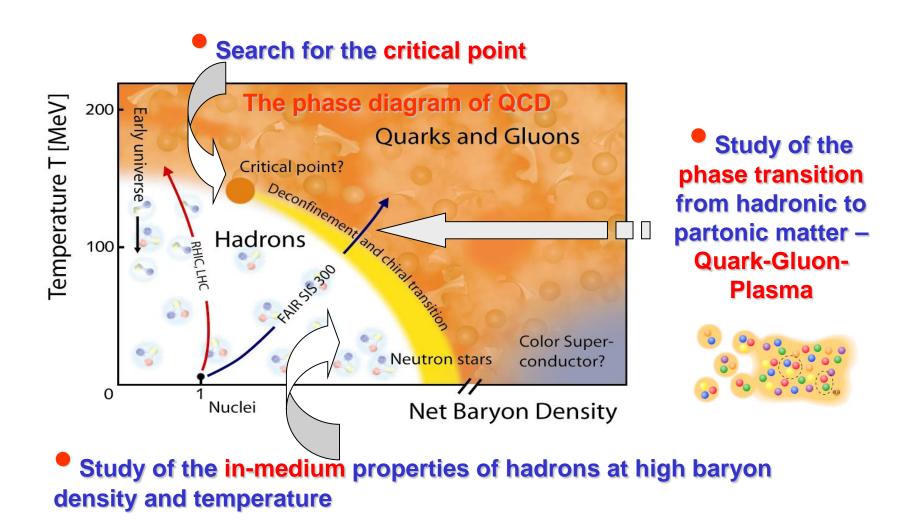
In collaboration with Taesoo Song, Wolfgang Cassing, Pierre Moreau



The 34th Winter Workshop on Nuclear Dynamics
The Langley Resort Fort Royal, Deshaies, Guadeloupe
March 25-31, 2018



The ,holy grail of heavy-ion physics:

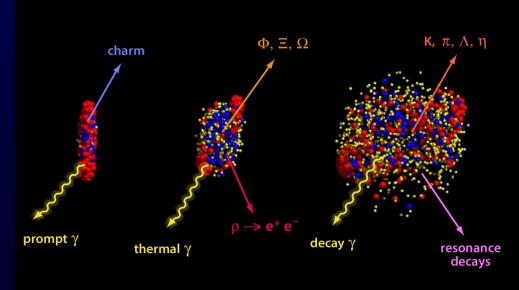


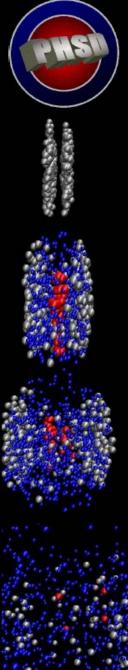
Electromagnetic probes: photons and dileptons

Feinberg (76), Shuryak (78)

- Advantages:
- ✓ dileptons and real photons are emitted from different stages of the reaction and not effected by final-state interactions
- ✓ provide undistorted information about their production channels
- ✓ promising signal of QGP ,thermal' photons and dileptons
- → Requires theoretical models which describe the dynamics of heavy-ion collisions during the whole time evolution!

- Disadvantages:
- low emission rate
- production from hadronic corona
- many production sources which cannot be individually disentangled by experimental data







PHSD is a non-equilibrium transport approach with

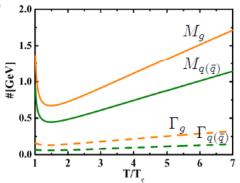
- explicit phase transition from hadronic to partonic degrees of freedom
- IQCD EoS for the partonic phase (,crossover at low μ_q)
- explicit parton-parton interactions between quarks and gluons
- dynamical hadronization
- ☐ QGP phase is described by the Dynamical QuasiParticle Model (DQPM)

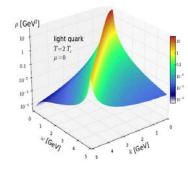
matched to reproduce lattice QCD

A. Peshier, W. Cassing, PRL 94 (2005) 172301; W. Cassing, NPA 791 (2007) 365: NPA 793 (2007)

- strongly interacting quasi-particles: massive quarks and gluons (g,q,q_{bar}) with sizeable collisional widths in a self-generated mean-field potential
- Spectral functions:

$$\rho_i(\omega,T) = \frac{4\omega\Gamma_i(T)}{\left(\omega^2 - \vec{p}^2 - M_i^2(T)\right)^2 + 4\omega^2\Gamma_i^2(T)}$$





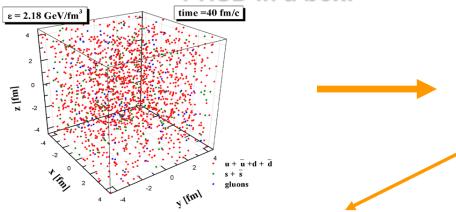
☐ Transport theory: generalized off-shell transport equations based on the 1st order gradient expansion of Kadanoff-Baym equations (applicable for strongly interacting systems!)



QGP in equilibrium: Transport properties at finite (T, μ_q): η /s

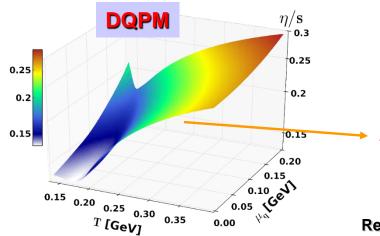
Infinite hot/dense matter =

PHSD in a box:



Shear viscosity η /s at finite (T, μ_q)

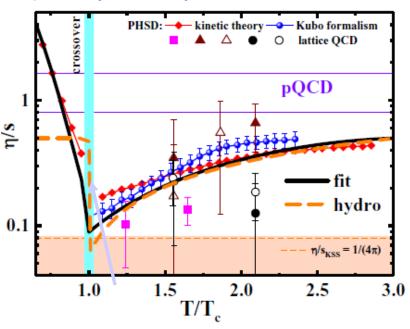
$$\frac{T_c(\mu_q)}{T_c(\mu_q = 0)} = \sqrt{1 - \alpha \, \mu_q^2} \approx 1 - \alpha/2 \, \mu_q^2 + \cdots$$



Shear viscosity η/s at finite T

PHSD: V. Ozvenchuk et al., PRC 87 (2013) 064903

Hydro: Bayesian analysis, S. Bass et al. ,1704.07671



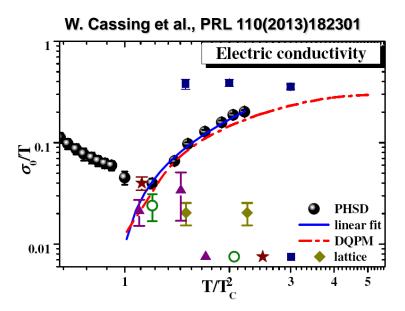
QGP in PHSD = stronglyinteracting liquid-like system

η/s: $\mu_q=0$ \rightarrow finite μ_q : smooth increase as a function of (T, μ_q)

Review: H. Berrehrah et al. Int.J.Mod.Phys. E25 (2016) 1642003

Transport properties at finite (T, μ_q): σ_e /T

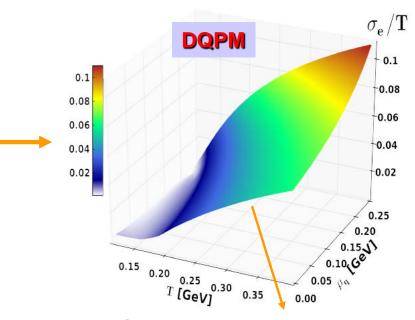
PHSD in a box: Electric conductivity σ_e/T at finite T



■ the QCD matter even at T~ T_c is a much better electric conductor than Cu or Ag (at room temperature) by a factor of 500 !

Electric conductivity σ_e/T at finite (T, μ_q)

H. Berrehrah et al. , PRC93 (2016) 044914



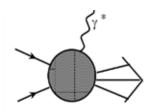
 σ_e/T : $\mu_q=0$ \rightarrow finite μ_q : smooth increase as a function of (T, μ_q)



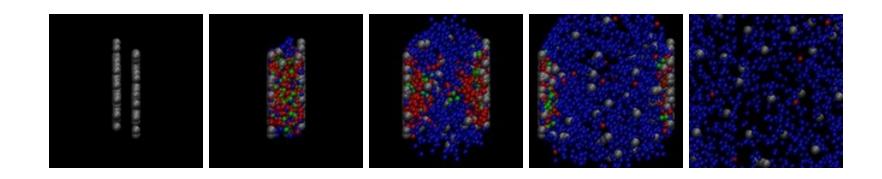
□ Photon emission: rates at $q_0 \rightarrow 0$ are related to electric conductivity σ_0

$$q_0 \frac{dR}{d^4 x d^3 q}\bigg|_{q_0 \to 0} = \frac{T}{4\pi^3} \sigma_0$$

 σ_0 \rightarrow Probe of electromagnetic properties of the QGP

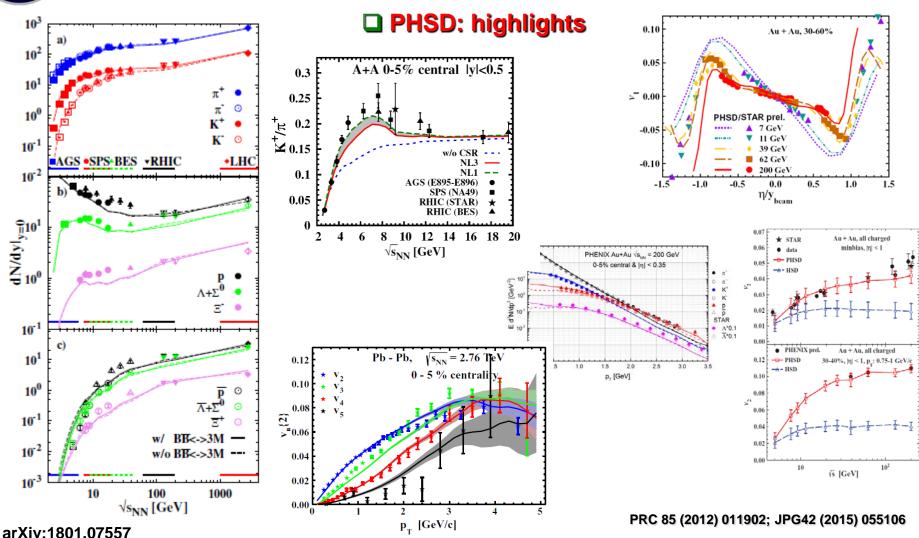


Bulk' properties in Au+Au collisions



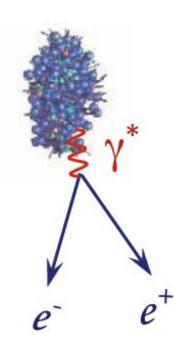


Non-equilibrium dynamics: description of A+A with PHSD

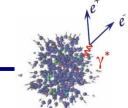


 \square PHSD provides a good description of ,bulk' observables (y-, p_T-distributions, flow coefficients v_n , ...) from SIS to LHC

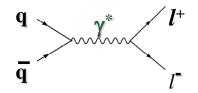
Dileptons as a probe of the QGP and in-medium effects

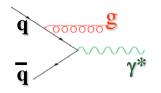


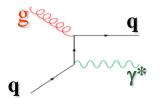
Dilepton sources

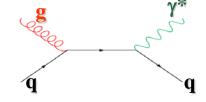












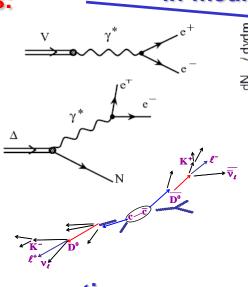
from hadronic sources:

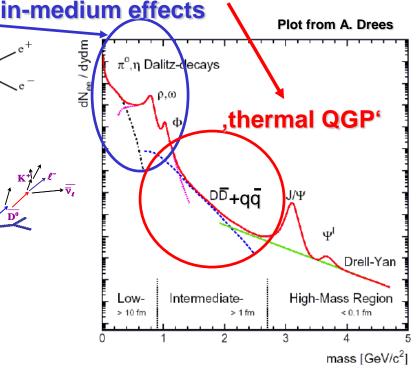
•direct decay of vector mesons $(\rho, \omega, \phi, J/\Psi, \Psi')$

Dalitz decay of mesons and baryons $(\pi^0, \eta, \Delta,...)$

•correlated D+Dbar pairs

•radiation from multi-meson reactions $(\pi+\pi, \pi+\rho, \pi+\omega, \rho+\rho, \pi+a_1)$ - 4π



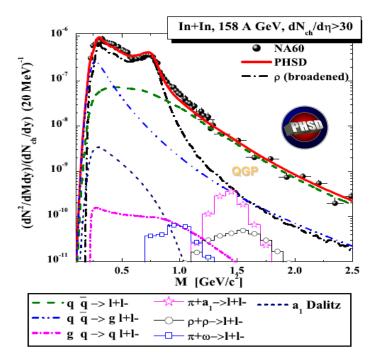


! Advantage of dileptons:

additional "degree of freedom" (M) allows to disentangle various sources

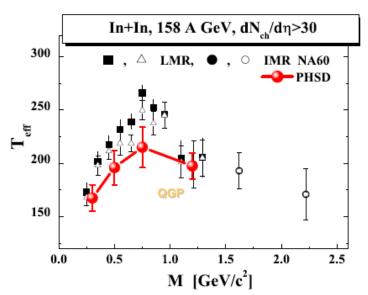
Lessons from SPS: NA60

■ Dilepton invariant mass spectra:



☐ Inverse slope parameter T_{eff}:

spectrum from QGP is softer than from hadronic phase since the QGP emission occurs dominantly before the collective radial flow has developed



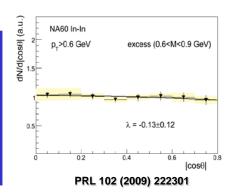
PHSD:

Linnyk et al, PRC 84 (2011) 054917

NA60: Eur. Phys. J. C 59 (2009) 607

Message from SPS: (based on NA60 and CERES data)

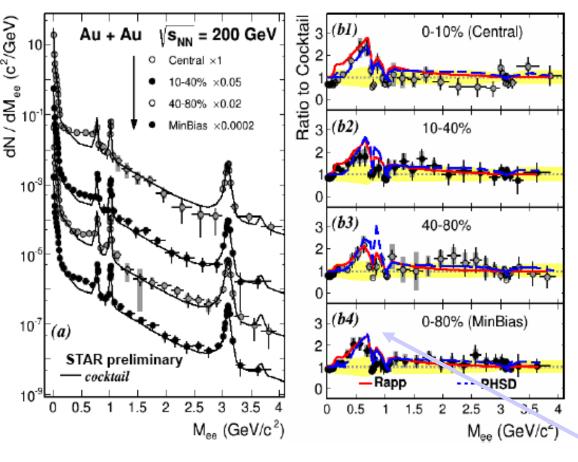
- 1) Low mass spectra evidence for the in-medium broadening of ρ-mesons
- 2) Intermediate mass spectra above 1 GeV dominated by partonic radiation
- 3) The rise and fall of Teff evidence for the thermal QGP radiation
- 4) Isotropic angular distribution indication for a thermal origin of dimuons



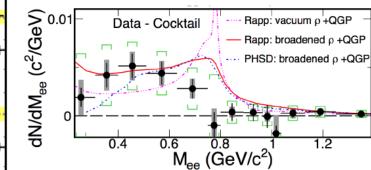
Dileptons at RHIC: STAR data vs model predictions

PRC 92 (2015) 024912

Centrality dependence of dilepton yield



Excess in low mass region, min. bias



Models:

- Fireball model R. Rapp
- PHSD

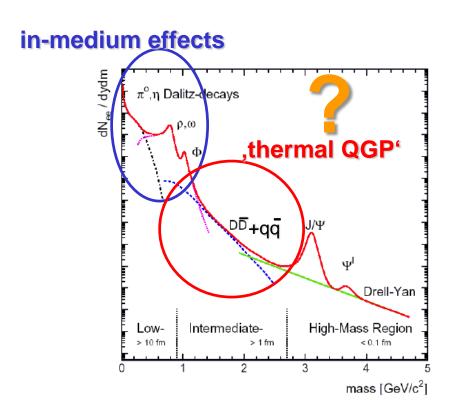
Low masses:

collisional broadening of ρ Intermediate masses:

QGP dominant

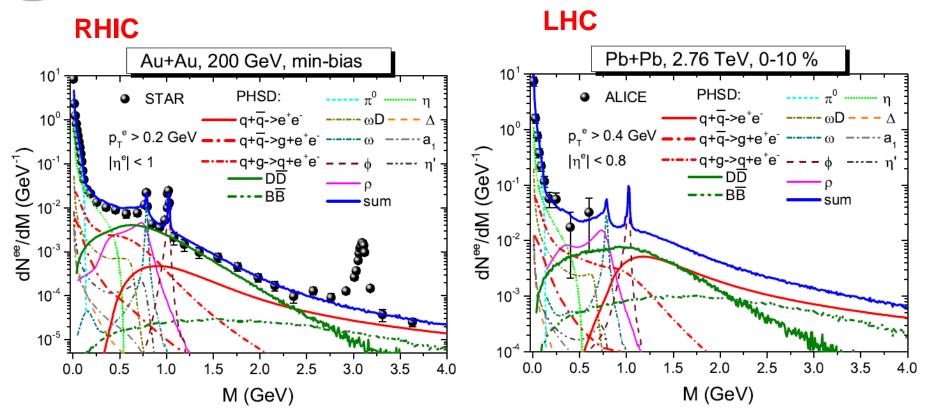
Message: STAR data are described by models within a collisional broadening scenario for the vector meson spectral function + QGP

What is the best energy range to observe thermal dileptons from QGP?





Dileptons at RHIC and LHC



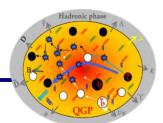
Message:

STAR data at 200 GeV and the ALICE data at 2.76 TeV are described by PHSD within

- 1) a collisional broadening scenario for the vector meson spectral functions
 - + QGP + correlated charm
- 2) Charm contribution is dominant for 1.2 < M < 2.5 GeV



Charm dynamics in PHSD

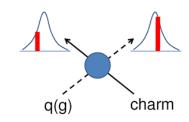




In order to get information about the QGP in HIC via dileptons, the charm dynamics must be under control

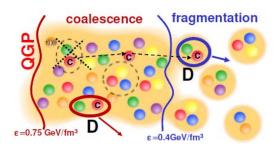
Dynamics of heavy quarks in A+A:

- 1. Production of heavy (charm and bottom) quarks in initial binary collisions + shadowing and Cronin effects
- Interactions in the QGP according to the DQPM:
 elastic scattering with off-shell massive partons Q+q→Q+q
 → collisional energy loss



3. Hadronization: c/cbar quarks \rightarrow D(D*)-mesons:

Dynamical hadronization scenario for heavy quarks : coalescence with <r>>=0.9 fm & fragmentation $0.4 < \epsilon < 0.75 \text{ GeV/fm}^3$ $\epsilon < 0.4 \text{ GeV/fm}^3$



4. Hadronic interactions:

D+baryons; D+mesons with G-matrix and effective chiral Lagrangian approach with heavy-quark spin symmetry

- T. Song et al., PRC 92 (2015) 014910, PRC 93 (2016) 034906, PRC 96 (2017) 014905
- T. Song, W.Cassing, P.Moreau and E.Bratkovskaya, arXiv:1803.02698



Charm production in NN collisions

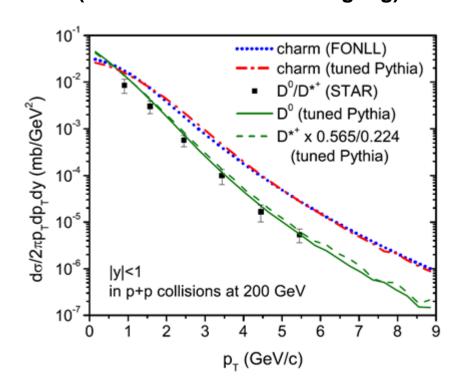
☐ A+A: charm production in initial NN binary collisions: probability

$$P = \frac{\sigma(c\overline{c})}{\sigma_{NN}^{inel}}$$

The total cross section for charm production in p+p collisions $\sigma(cc)$

Alice_ATLAS (a) 10⁴ LHCb STAR Alice 10³ PHENIX 2011 $\sigma_{pp}(c\overline{c})$ (μb) PHENIX 2017 10² E653 (pA) NA16 (pA) 7 E743 (pA) E769 (pA) 10° Experimental data 10⁻¹ **PHSD** 10⁻² -10¹ 10² 10^{3} 10⁴ s^{1/2} (GeV)

Momentum distribution of heavy quarks: use ,tuned' PYTHIA event generator to reproduce FONLL (fixed-order next-to-leading log) results



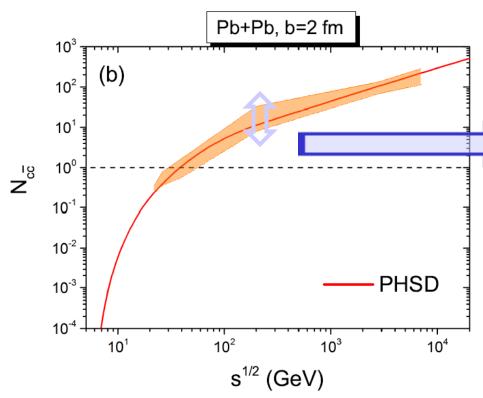
T. Song, W.Cassing, P.Moreau and E.Bratkovskaya, arXiv:1803.02698

T. Song et al., PRC 92 (2015) 014910, PRC 93 (2016) 034906, PRC 96 (2017) 014905



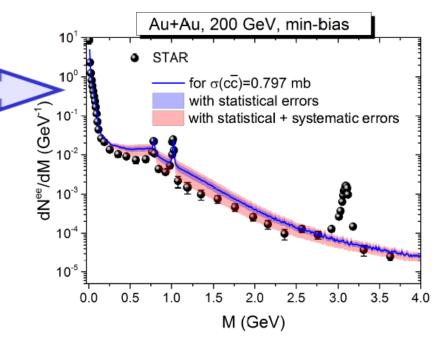
Charm at RHIC and LHC

The number of primary cc pairs in Pb+Pb collisions at b=2 fm as a function of s^{1/2}



* The shaded area shows the uncertainty in the number of cc pairs due to the uncertainty in the charm production cross section in p+p collisions

The invariant mass spectra of dielectrons for min-bias Au+Au at 200 GeV with the $\sigma(cc)$ from the STAR with statistical and systematic errors

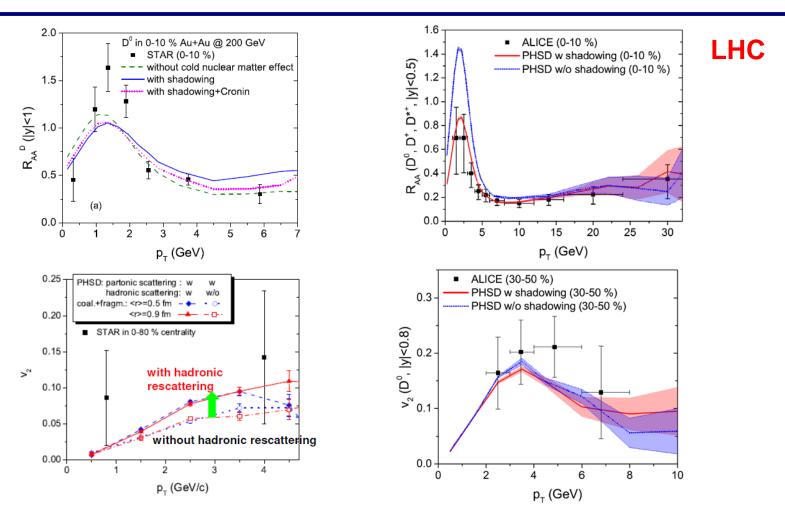


- Uncertainty in σ(cc) from pp leads to the uncertainty in the charm production in AA and in the dilepton spectra!
- \rightarrow Reliable data for $\sigma(cc)$ from pp are needed!



PHSD vs charm observables at RHIC and LHC

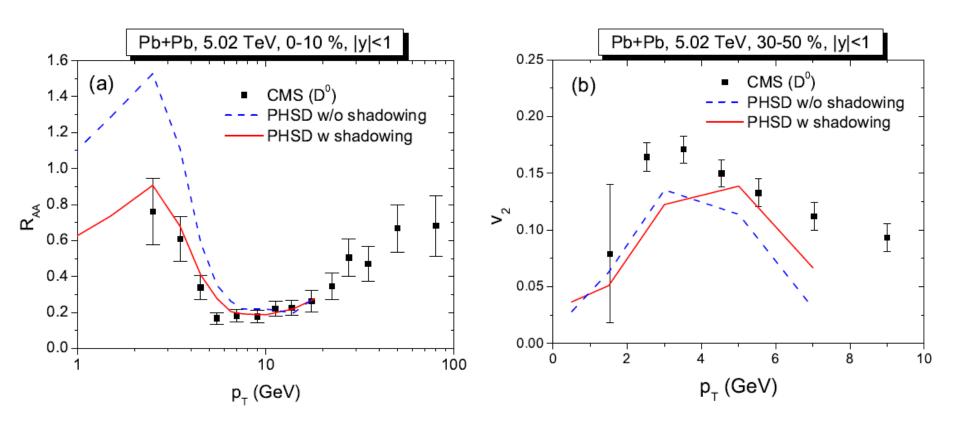
RHIC



- ☐ The exp. data for the R_{AA} and v₂ at RHIC and LHC are described in the PHSD by QGP collisional energy loss due to elastic scattering of charm quarks with massive quarks and gluons in the QGP
- + by the dynamical hadronization scenario "coalescence & fragmentation"
- + by strong hadronic interactions due to resonant elastic scattering of D,D* with mesons and baryons



Charm at LHC: central Pb+Pb at 5.02 TeV



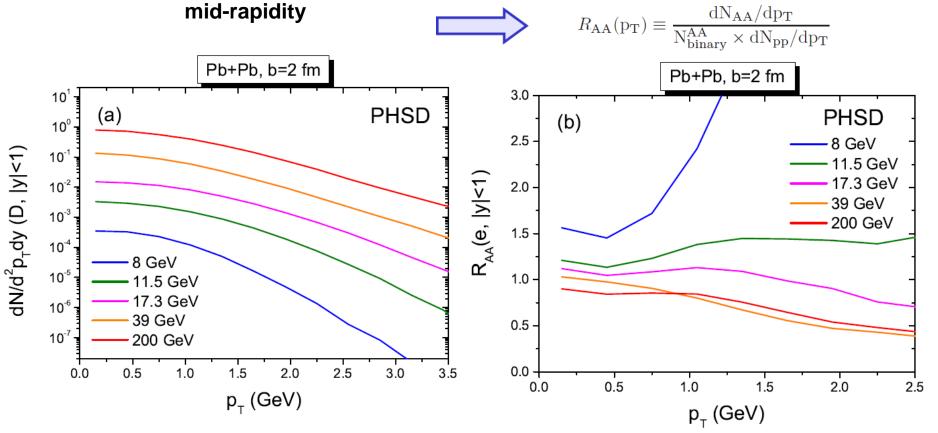
→ PHSD shows a good agreement with CMC data



Nuclear modification of dielectrons from heavy flavor

The transverse momentum spectra of D-mesons at s^{1/2} from 8 to 200 GeV at mid-rapidity

R_{AA}(p_T) of single electrons from semi-leptonic decay of D-mesons

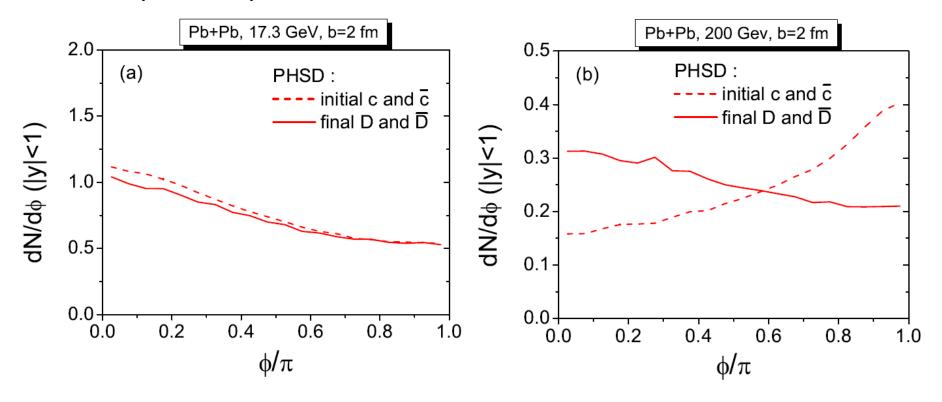


- □ Hardening of the p_T spectra of D-mesons with increasing incoming energy
- □ R_{AA}(p_T) of single electrons from suppression at high energy to enhancement at low energy



Angular correlation between D-Dbar

Azimuthal angular distribution between the transverse momentum of D-Dbar at midrapidity (|y| < 1) before (dashed lines) and after the interactions with the medium (solid lines) in central Pb+Pb collisions at $s^{1/2} = 17.3$ and 200 GeV

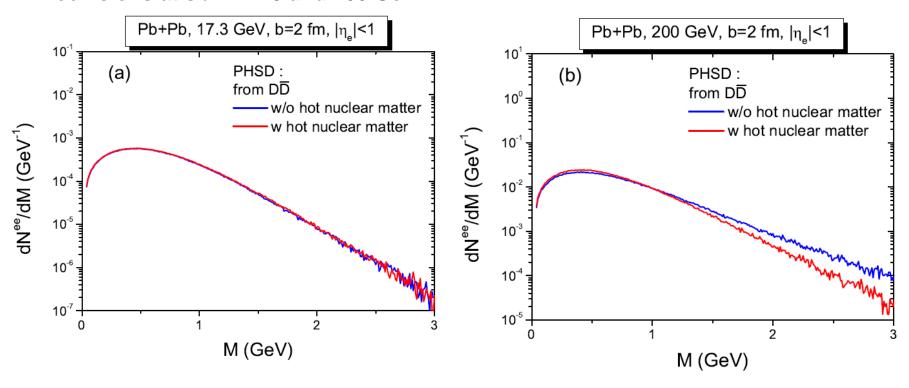


- □ Initial correlations from PYTHIA : peaks around ϕ = 0 for \sqrt{s} = 17.3 GeV, while around ϕ = π for \sqrt{s} = 200 GeV
- □ Final correlations: smeared at √s= 200 GeV due to the interaction of charm quarks in QGP



Modification of dielectron spectra due to the in-medium interaction of D-Dbar

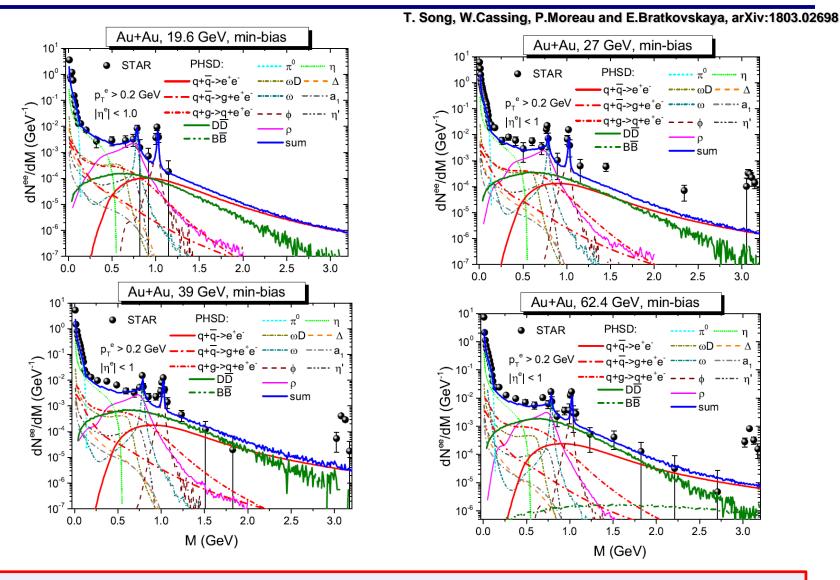
The invariant mass spectra of dielectrons from charm pairs with (red lines) and without the interactions with the hot medium (blue lines) in central Pb+Pb collisions at $s_{1/2} = 17.3$ and 200 GeV



- \square Softening of dN/dM at \sqrt{s} = 200 GeV due to the interaction of charm quarks in QGP
- Note: the invariant mass of the dielectrons depends on the momenta of e⁺, e⁻ and also on the angle between them $\Rightarrow R_{AA}(p_T)$ shows that the momenta of e⁺, e⁻ are suppressed and $dN/d\phi$ shows that the azimuthal angle between them decreases at \sqrt{s} = 200 GeV



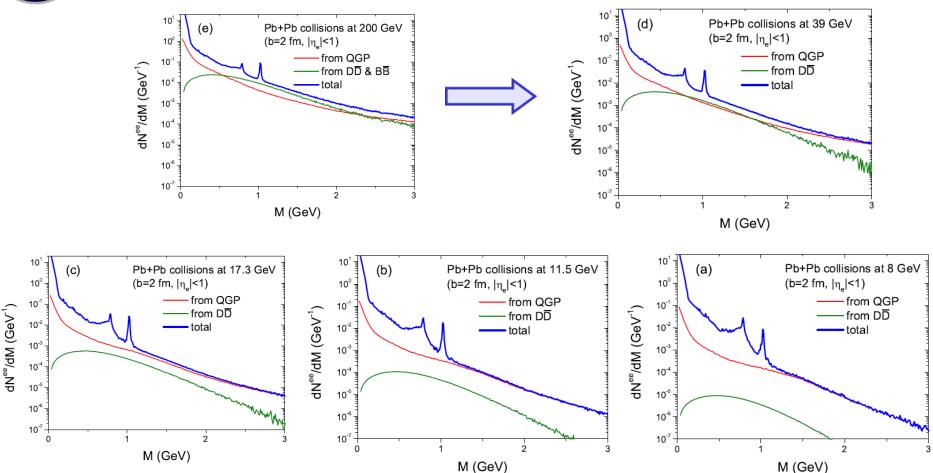
Dileptons from RHIC BES: STAR



QGP and charm are dominant contributions for intermediate masses at BES RHIC → measurements of charm at BES RHIC are needed to control charm production!



Dileptons at FAIR/NICA energies: predictions

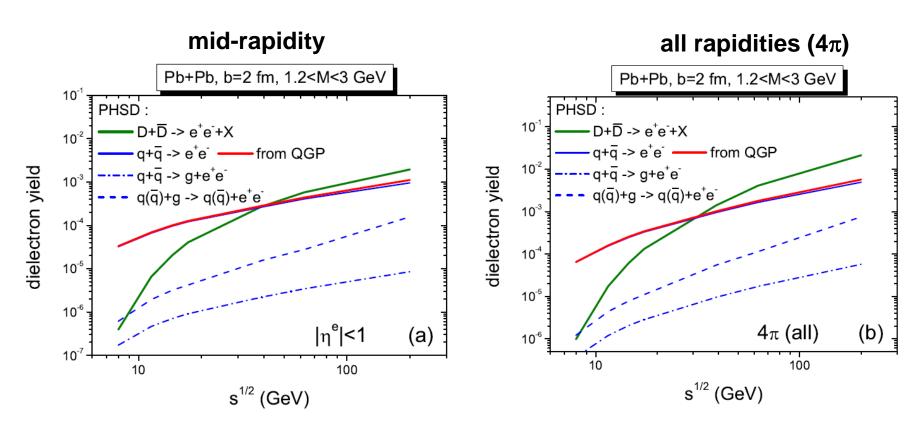


Relative contribution of QGP versus charm increases with decreasing energy!



Dileptons: QGP vs charm

Excitation function of dilepton multiplicity integrated for 1.2<M<3GeV

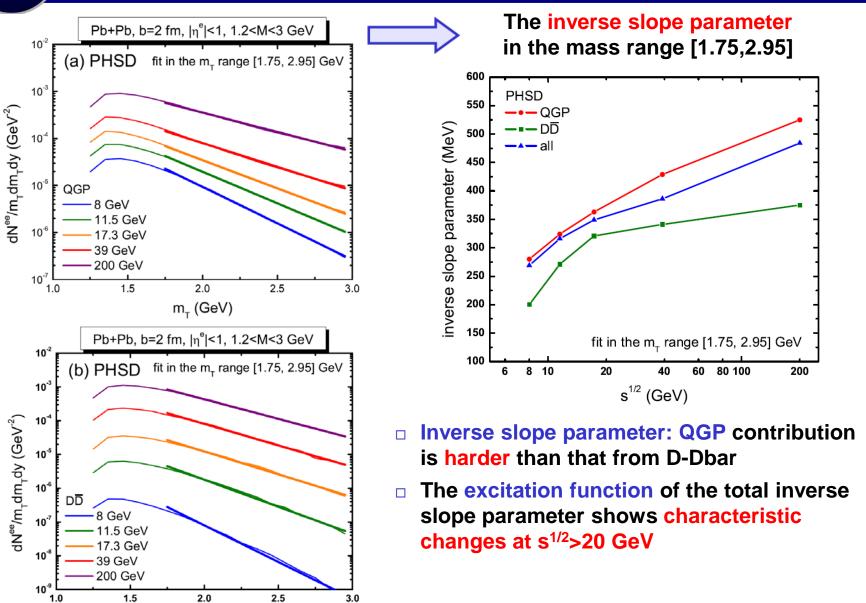


QGP contribution overshines charm with decreasing energy!

→ Good perspectives for FAIR/NICA and BES RHIC!



Dilepton transverse mass spectra



m₋ (GeV)

Messages from the dilepton study

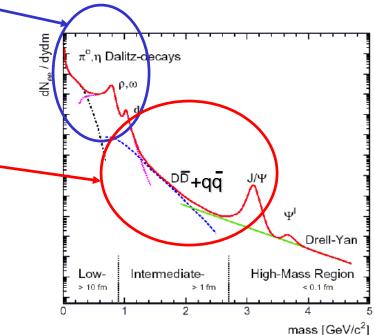


■ Low dilepton masses:

Dilepton spectra show sizeable changes due to the in-medium effects – modification of the properties of vector mesons (as collisional broadening) – which are observed experimentally

In-medium effects can be observed at all energies from SIS to LHC; excess increasing with decreasing energy due to a longer ρ-propagation in the high baryon-density phase

- Intermediate dilepton masses M>1.2 GeV :
- Dominant sources : QGP (qbar-q) , correlated charm D/Dbar
- Fraction of QGP grows with increasing energy; however, the relative contribution of QGP to dileptons from charm pairs increases with decreasing energy
- → Good perspectives for FAIR/NICA



Review: O. Linnyk et al., Prog. Part. Nucl. Phys. 89 (2016) 50

Thank you for your attention!



WWND-2018