

Cold nuclear matter effects on
quarkonium production @ RHIC and LHC:

Fractional energy loss on Υ and J/ψ production

Elena G. Ferreiro

Universidade de Santiago de Compostela, Spain

in collaboration with

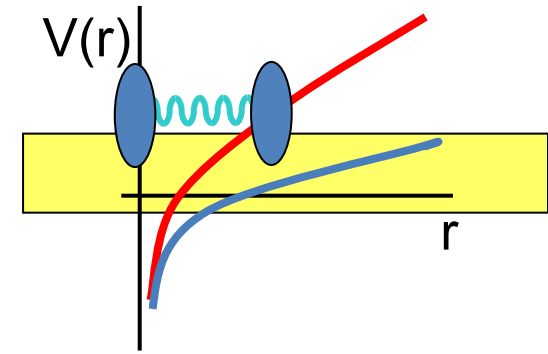
F. Fleuret, J-P. Lansberg, N. Matagne and A. Rakotozafindrabe

EPJC61(2009), PLB680(2009), PRC81(2010), NPA855(2011), arXiv:1110:5047

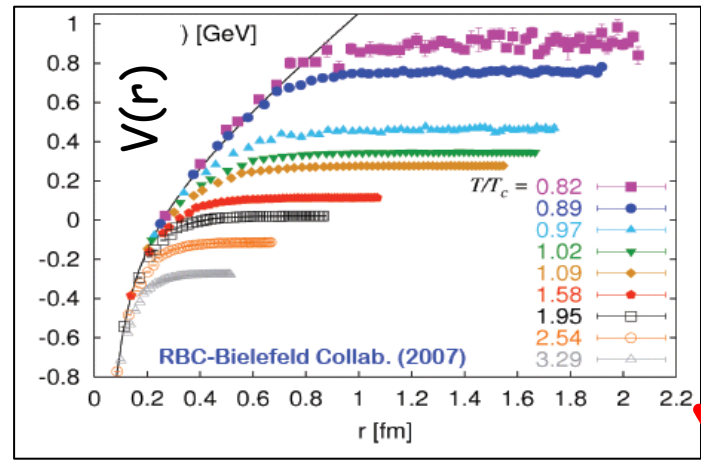
Introduction I: the intriguing story of J/ψ production

Potential between q-anti-q pair grows linearly at large distances

$$V(r) = -\frac{4}{3} \frac{\alpha_s}{r} + kr$$



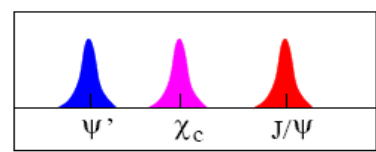
Screening of long range confining potential at high enough temperature or density.



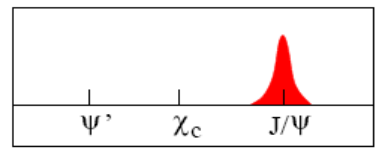
What happens when the range of the binding force becomes smaller than the radius of the state?
 different states “melting” at different temperatures due to different binding energies.

Matsui and Satz:

J/ψ destruction in a QGP by Debye screening

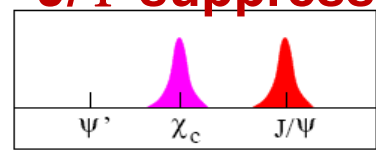


$T < T_c$

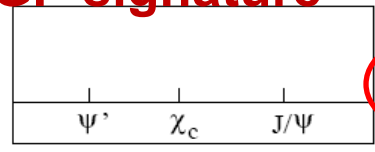


$T \sim 1.1 T_c$

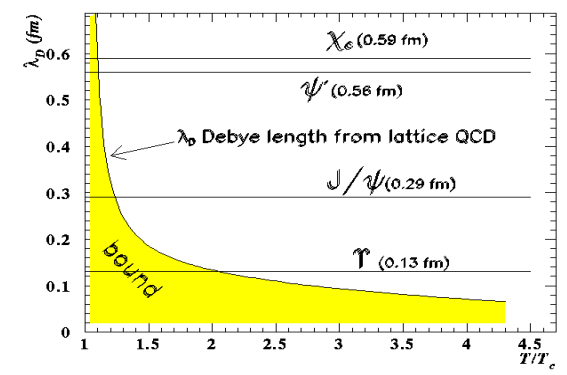
J/ψ suppression = QGP signature



$T \sim T_c$



$T \gg T_c$



Introduction II: motivation

- A lot of work trying to understand **A+A** data (since $J/\psi \equiv$ QGP signal)

Quarkonium as a hint of deconfinement

QGP

- If we focalise on **p+A** data (where no QGP is possible)
only cold nuclear matter (CNM) effects are in play here:

shadowing and nuclear absorption
EMC and energy loss

Quarkonium as a hint of coherence

nPDF

- In fact, the question is even more fundamental: **p+p** data
we do not know the specific production kinematics at a **partonic** level:

$(2 \rightarrow 2, 3, 4)$ vs $(2 \rightarrow 1)$

Quarkonium as a hint of QCD

QCD

Introduction III: contents

Our goal:

To investigate the **CNM effects** and the impact of the specific **partonic production** kinematics

3 ingredients:

- **Quarkonium partonic production mechanism**
- **Shadowing**
- **Nuclear absorption**

• Results on J/ψ production @ RHIC and LHC

To extend our study to Υ and other CNM effects :

• ***fractional energy loss***

• ***gluon EMC effect***

• Results on Υ and J/ψ production @ RHIC

• Results on Υ @ LHC

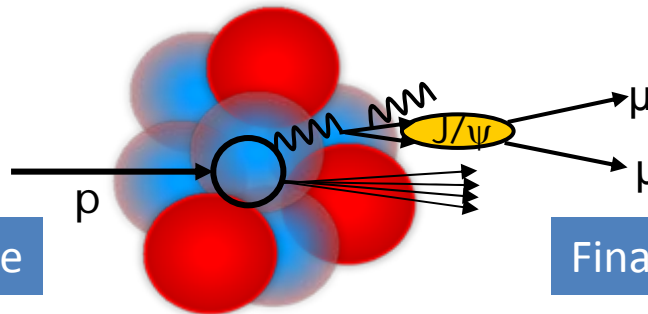
Quarkonium suppression in p+A collisions: CNM effects

Quarkonium production is suppressed in nuclear collisions ...but for a variety of reasons

- dissociation by screening (“melting”) and/or collisions in hot QGP

QGP effects
A+A collisions

- shadowing,
- saturation
- intrinsic charm



- nuclear absorption
- final energy loss
- comovers

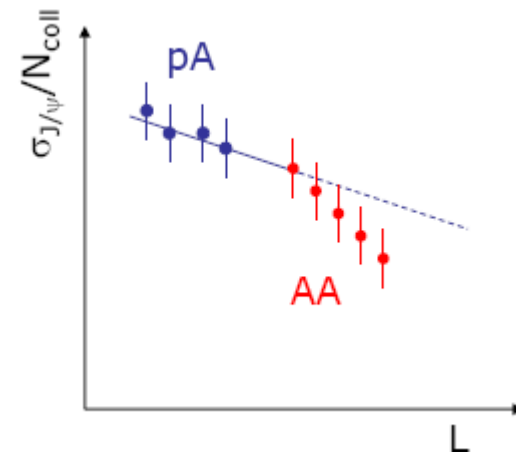
CNM effects
p+A and A+A collisions

To understand quarkonium behaviour in the hot medium, it's important to know its behaviour in the cold nuclear matter. This information can be achieved studying pA collisions

The cold nuclear matter effects present in pA collisions are of course present also in AA and can mask genuine QGP effects

It is very important to measure cold nuclear matter effects (CNM) before any claim of an “anomalous” (QGP) suppression in AA collisions

CNM, evaluated in pA, are extrapolated to AA, in order to build a reference for the J/ψ behaviour in hadronic matter



Quarkonium as a tool of COLD and HOT effects

● **cold effects:** wo thermalisation **NO QGP**

gluon shadowing
gribov shadowing
nuclear structure functions in nuclei \neq superposition of constituents nucleons
NI@SPS, IMP@RHIC

nuclear absorption
multiple scattering of a pre-resonance c-cbar pair within the nucleons of the nucleus
IMP@SPS, RHIC?

CGC
parton saturation
non-linear effects favoured by the high density of partons become important and lead to eventual saturation of the parton densities
non thermal colour connection

partonic comovers
hadronic comovers
dissociation of the c-cbar pair with the dense medium produced in the collision
partonic or hadronic
suppression by a dense medium, not thermalized

Others:
EMC effect, energy loss

● **hot effects:** w thermalisation **QGP**

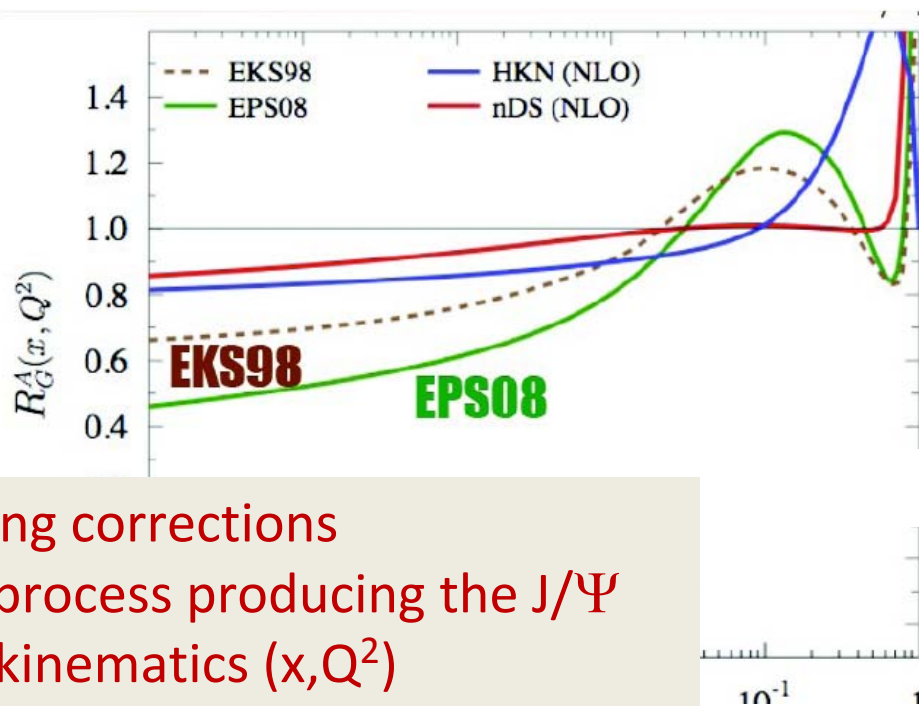
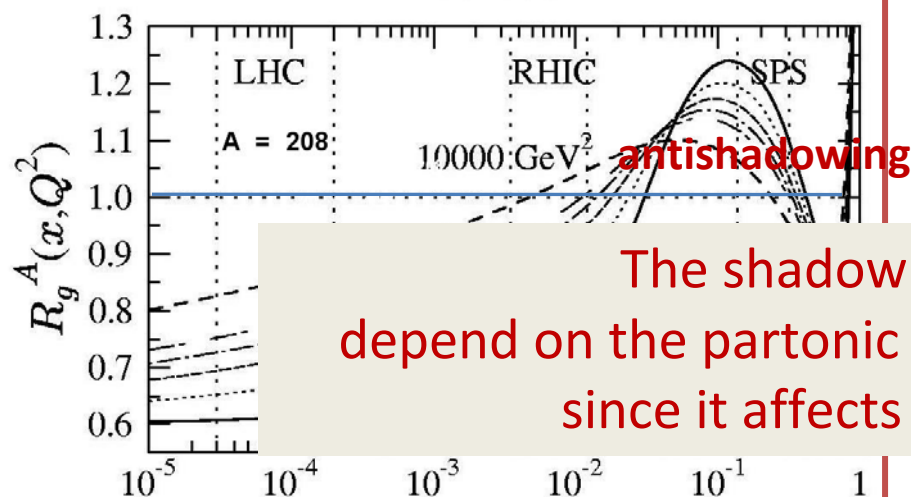
QGP sequential suppression recombination

Shadowing: an initial cold nuclear matter effect

- Nuclear shadowing is an initial-state effect on the partons distributions
- Gluon distribution functions are modified by the nuclear environment
- PDFs in nuclei different from the superposition of PDFs of their nucleons

Shadowing effects increases with energy ($1/x$) and decrease with Q^2 (m_T)

$$R_i^A(x, \mu_f) = \frac{f_i^A(x, \mu_f)}{A f_i^{\text{nucleon}}(x, \mu_f)}, \quad f_i = q, \bar{q}, g$$



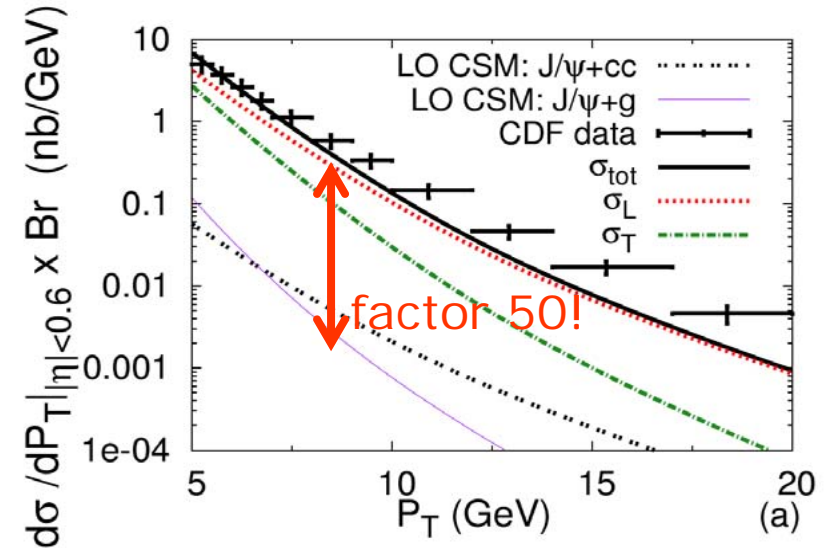
The shadowing corrections depend on the partonic process producing the J/Ψ since it affects kinematics (x, Q^2)

J/ψ production mechanisms at partonic level

- **First proposed: CSM @ LO**

- perturbative creation of the c \bar{c} pair in *color singlet state* with subsequent binding to J/ψ with same quantum numbers
- **hard gluon emission**

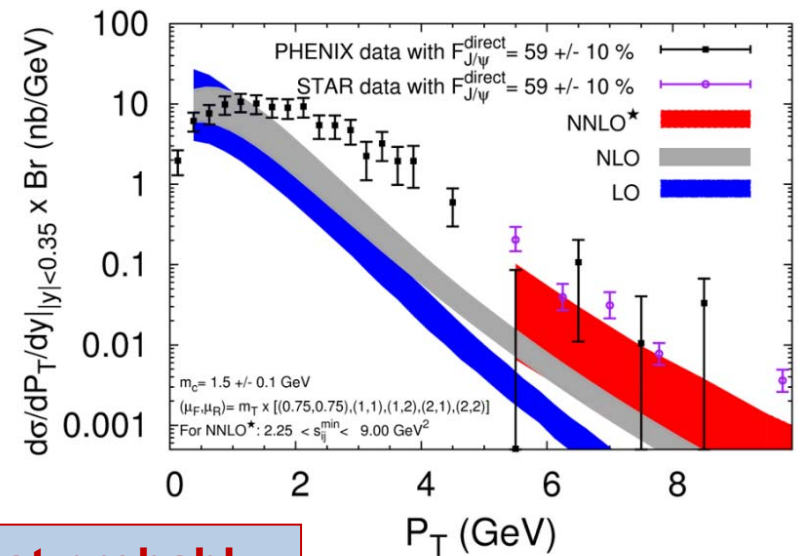
But CDF results on J/ψ direct production revealed a striking discrepancy wrt LO CSM



- **Second proposed: COM @ LO**

- uses NRQCD formalism to describe the non-perturbative hadronization of the c \bar{c} color octet to the color singlet state via soft gluon emission

The agreement improves in NRQCD approach but it does not describe polarization



Recently many step forwards (NLO and NNLO corrections...)

**Most probably:
CS+CO @ NLO**

On the kinematics of J/ψ production: two approaches

- CNM **-shadowing-** effects depends on J/ψ kinematics (x, Q²)
- J/ψ kinematics depends on the production mechanism =>

Investigating two production mechanisms (including p_T for the J/ψ):

$$g+g \rightarrow J/\psi \quad \mathbf{2 \rightarrow 1} \quad x_{1,2} = \frac{m_T}{\sqrt{s_{NN}}} \exp(\pm y)$$

- **intrinsic scheme:** the p_T of the J/ψ comes from initial partons
 - ❖ Not relevant for, say, p_T > 3 GeV
 - ❖ Only applies if COM(LO, α_s²) is the relevant production mechanism at low

p_T

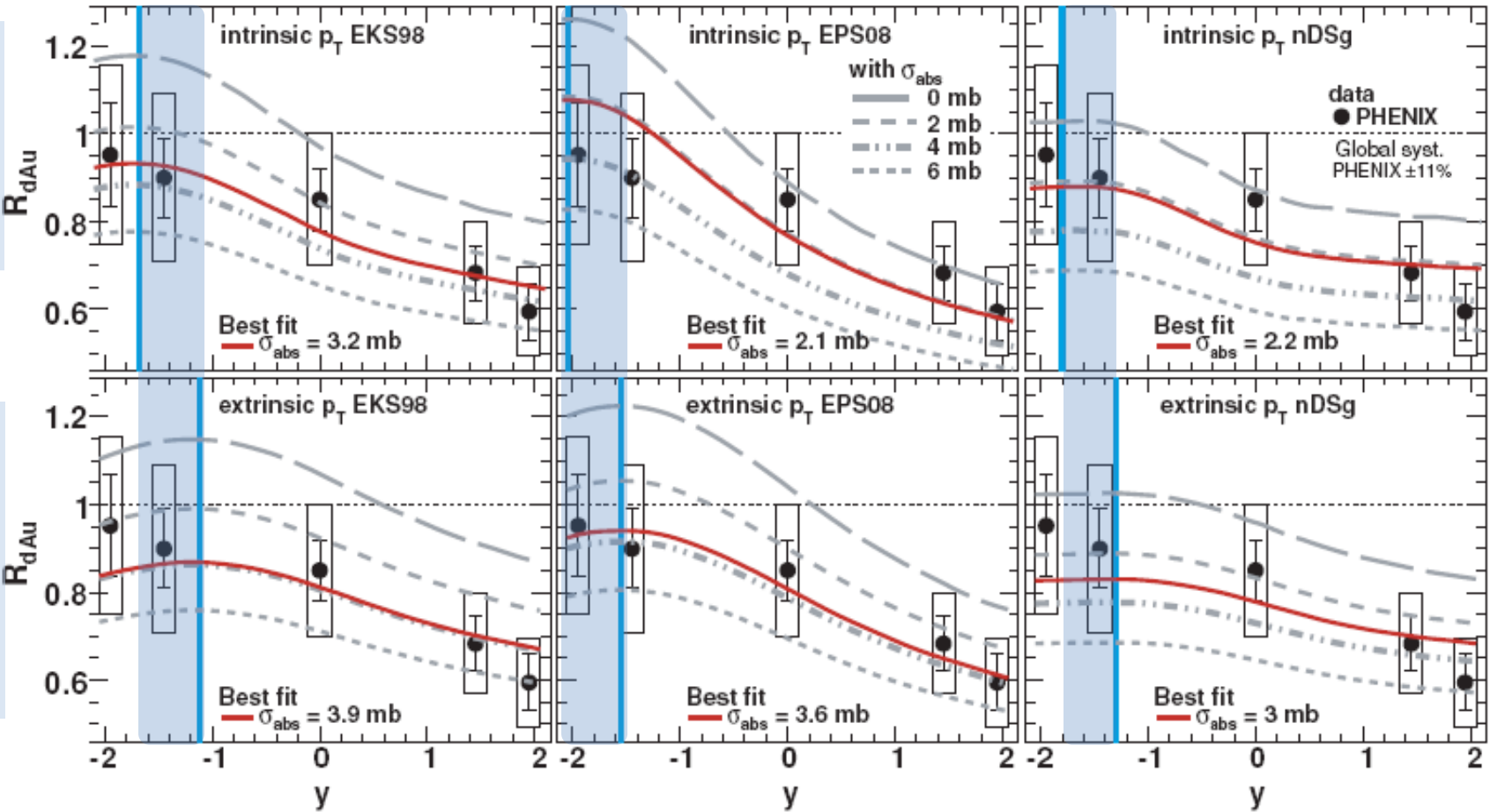
$$g+g \rightarrow J/\psi+g, gg, ggg, \dots \quad \mathbf{2 \rightarrow 2, 3, 4} \quad x_2 = \frac{x_1 m_T \sqrt{s_{NN}} e^{-y} - M^2}{\sqrt{s_{NN}} (\sqrt{s_{NN}} x_1 - m_T e^y)}$$

- **extrinsic scheme:** the p_T of the J/ψ is balanced by the outgoing parton(s)
 - ❖ COM (NLO, NNLO),
 - ❖ CSM (LO, NLO, NNLO)
- for a given y, larger x in extrinsic scheme => modification of shadowing effects

How partonic kinematics affects CNM: J/ψ in d+Au @ RHIC

2 → 1
g+g → J/ψ

2 → 2
g+g → J/ψ+g



E. G. Ferreiro, F. Fleuret, J. P. Lansberg and A. Rakotozafindrabe PRC 81 (2010) 064911

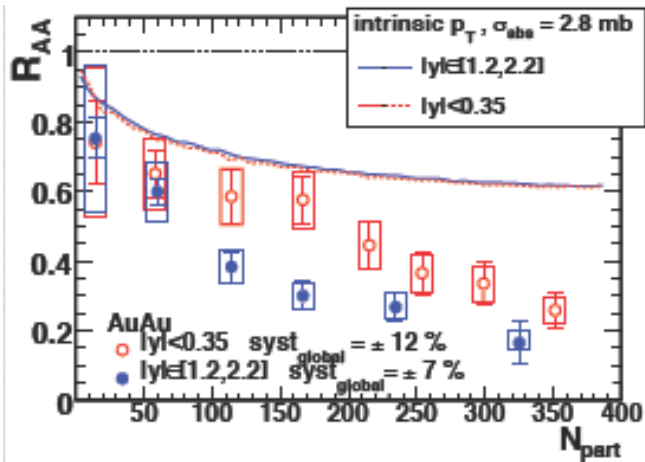
- shadowing depends on the partonic process: 2 → 1 or 2 → 2
- in order to reproduce data @ RHIC: nuclear absorption

σ_{abs} extrinsic > σ_{abs} intrinsic the kinematics matter for the extraction of σ_{abs}

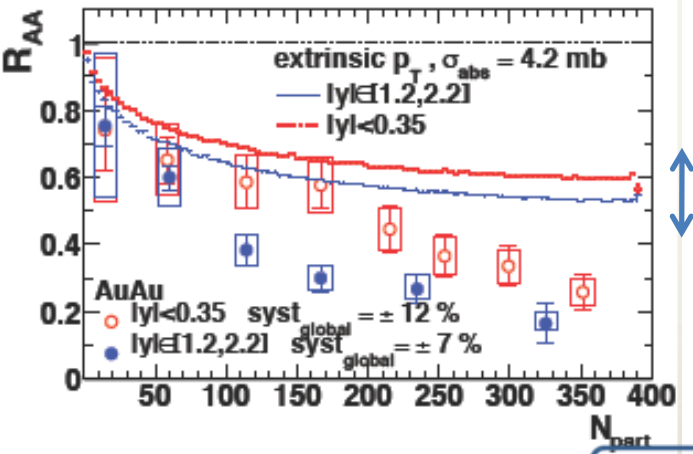
How partonic kinematics affects CNM: J/ψ in Au+Au @ RHIC

mid-y & forward-y

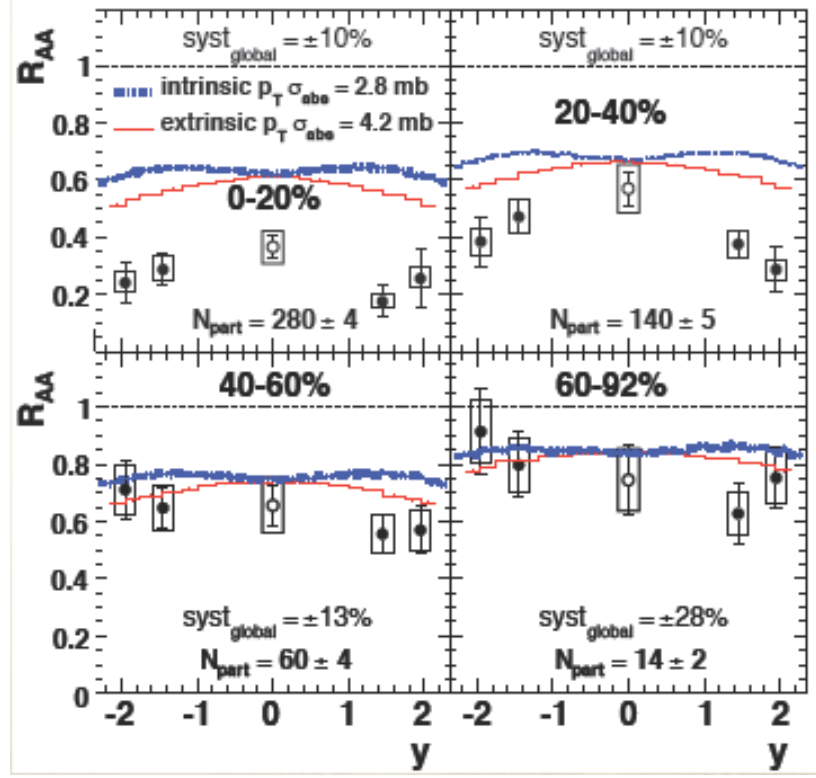
Intrinsic scheme: 2→1



Extrinsic scheme: 2→2



2→1 & 2→2 process

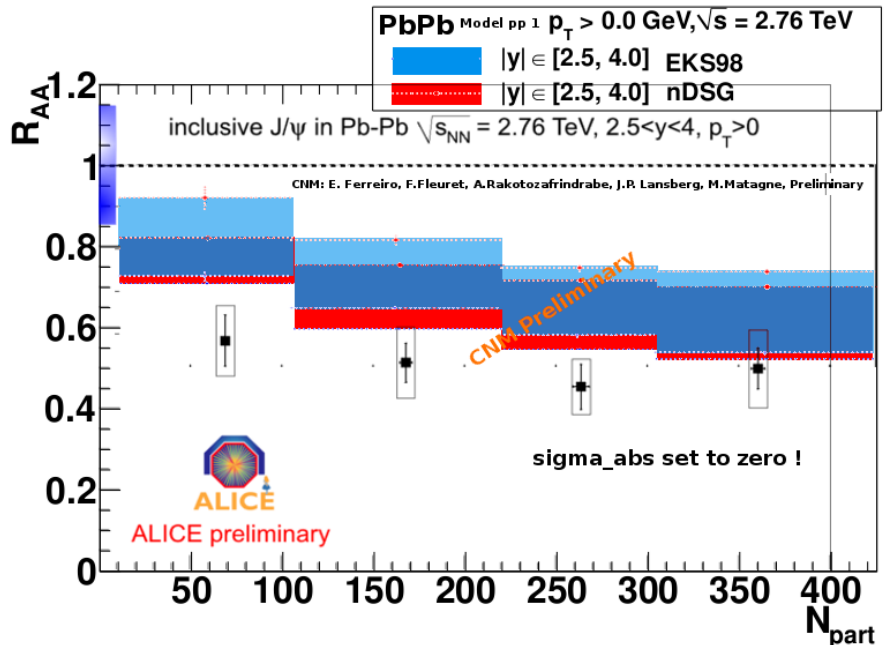


Hot Nuclear matter effects of course needed, but...
by considering the correct partonic kinematics

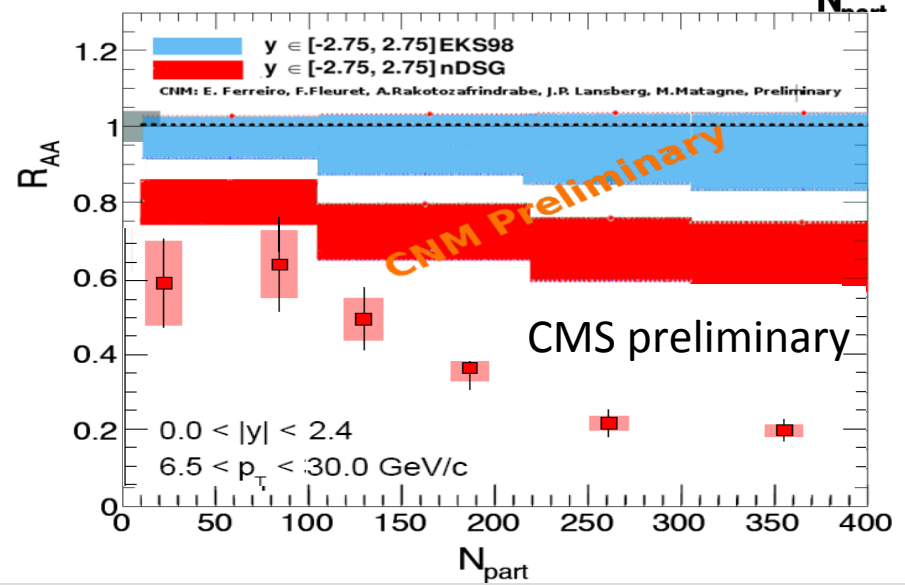
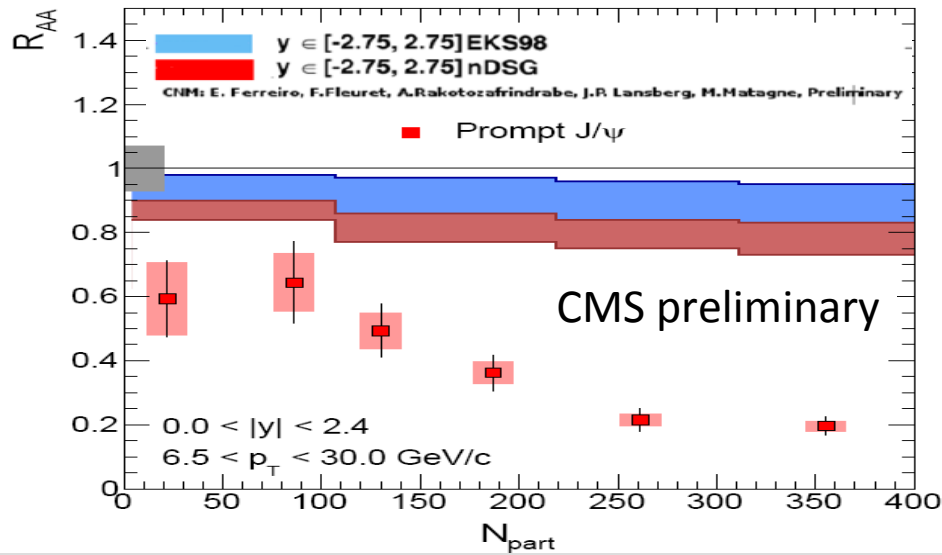
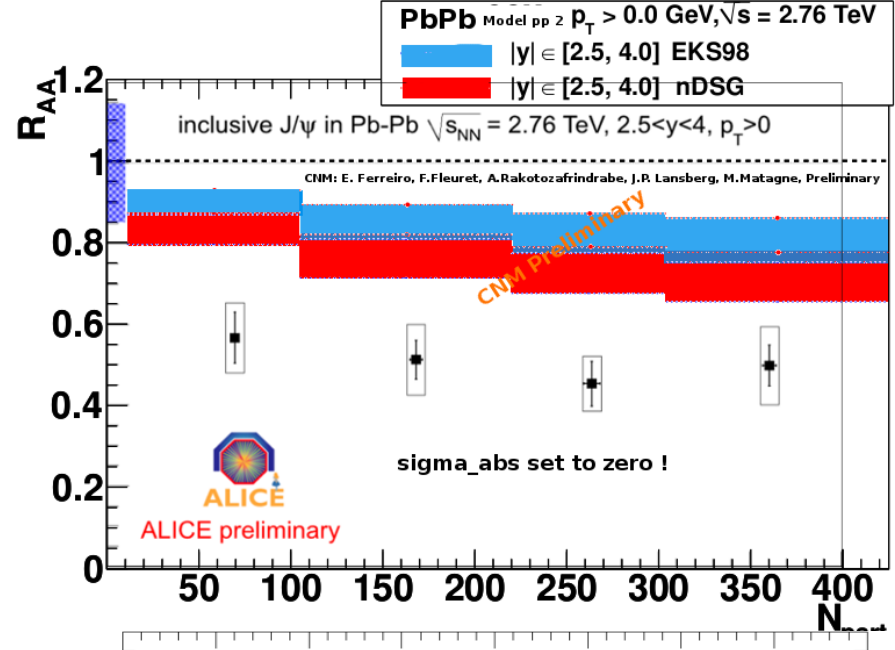
CNM effects could explain the mid/forw @ RHIC

Going to higher energy: J/ψ in Pb+Pb @ LHC

“CEM NLO” before k_T smearing

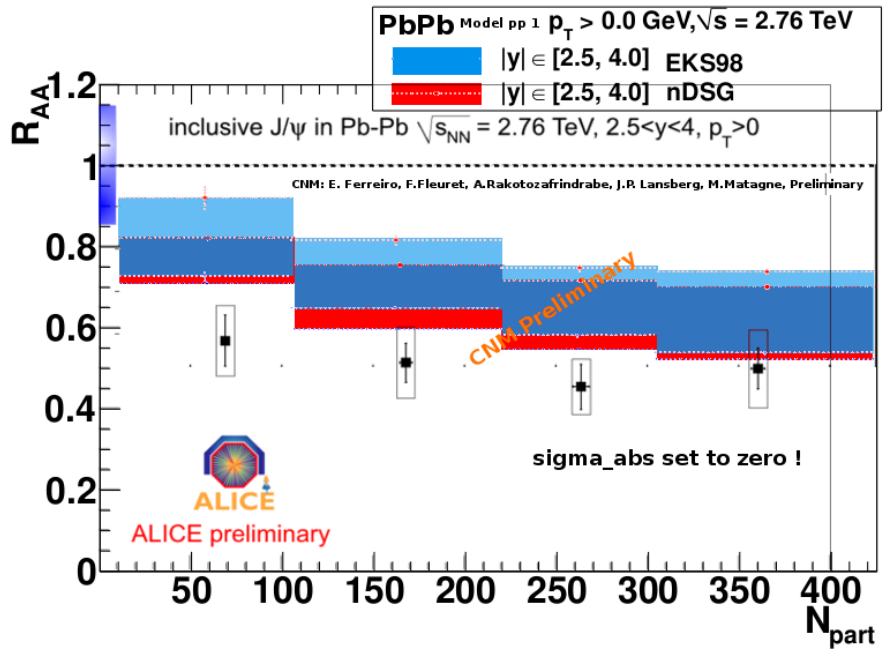


“Traditional” 2 → 2

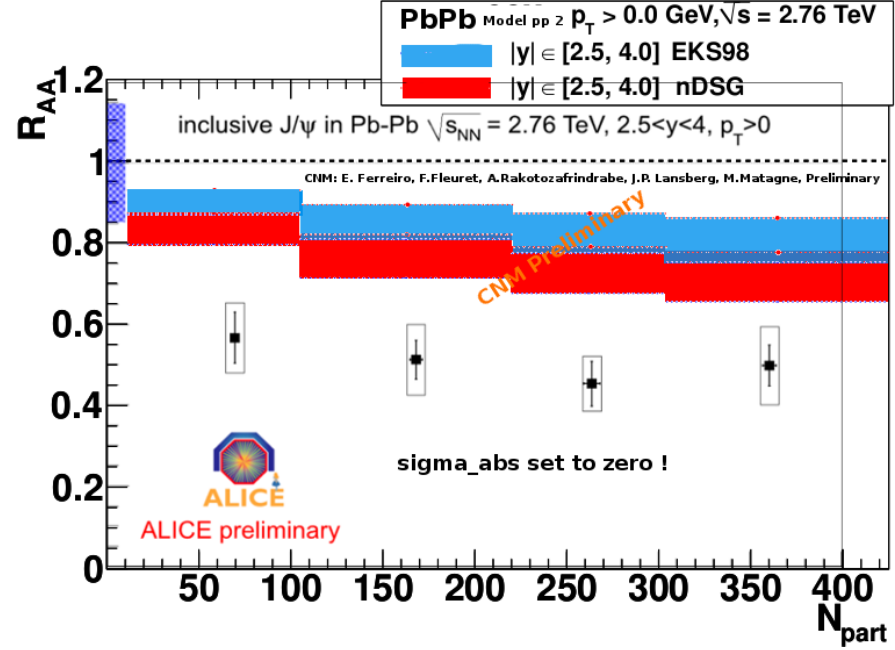


Going to higher energy: J/ψ in Pb+Pb @ LHC

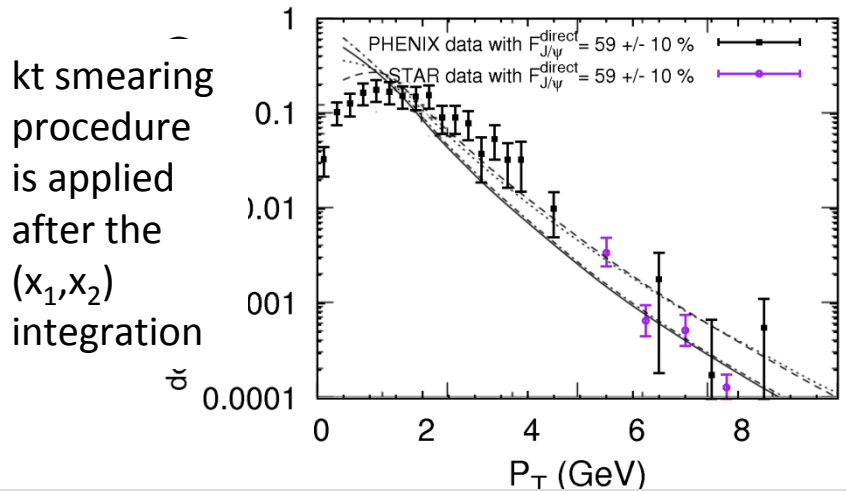
“CEM NLO” before k_T smearing



“Traditional” 2 → 2

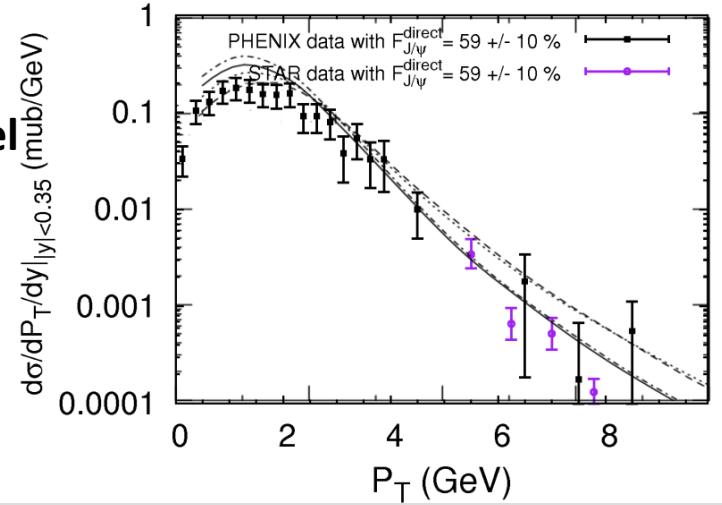


without k_T smearing



underlying partonic model

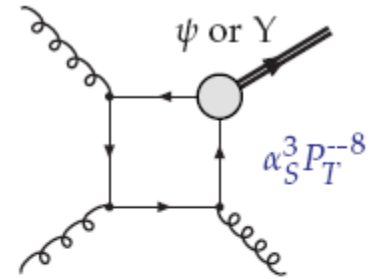
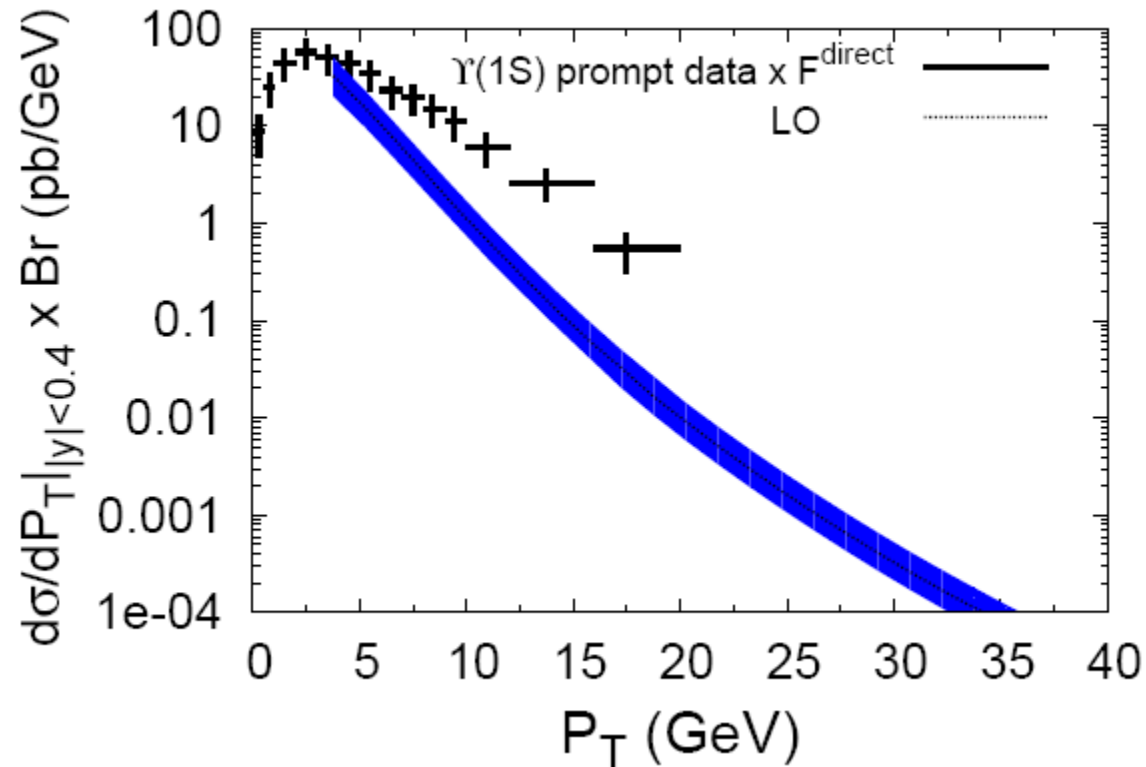
with k_T smearing



- **Gluc** EMC effect
- **Fractional** energy loss

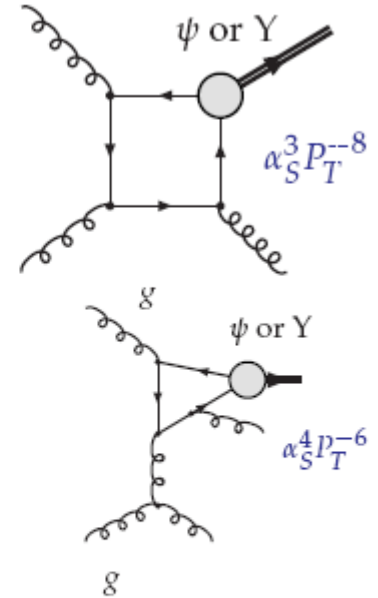
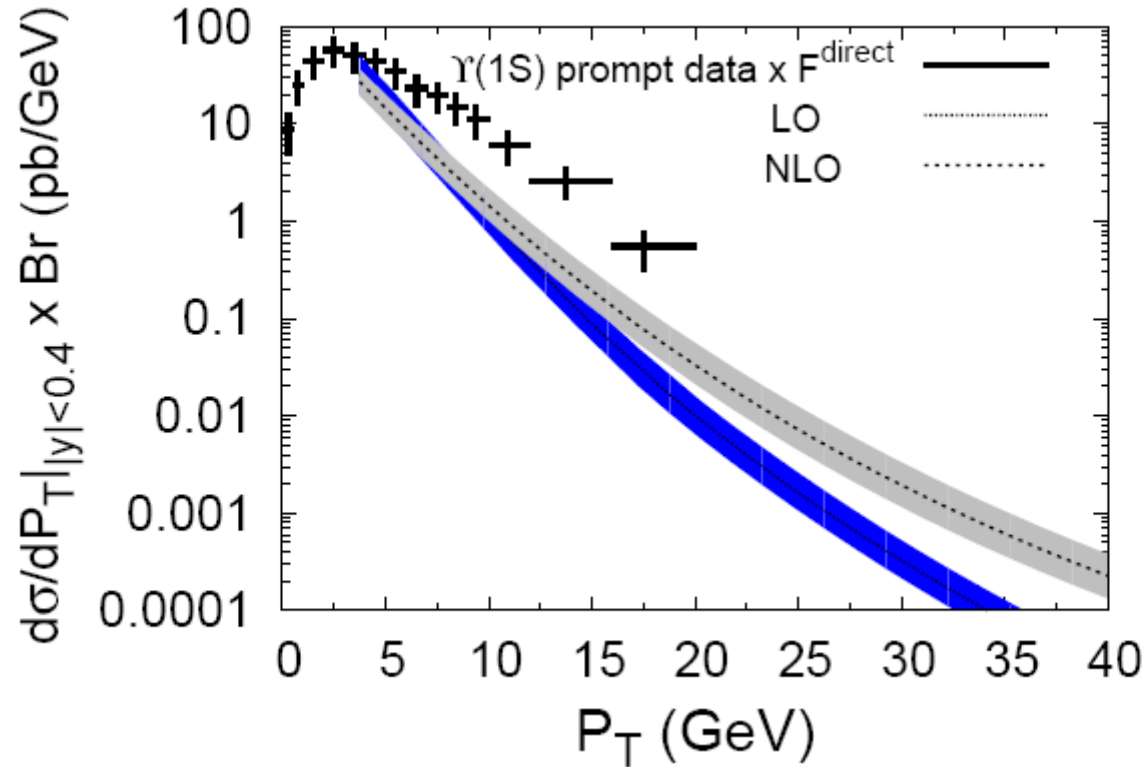
Remainder: QCD corrections for Υ in p+p @ Tevatron

P. Artoisenet, J. Campbell, J.P. Lansberg, F. Maltoni, PRL 101 (2008) 152001.



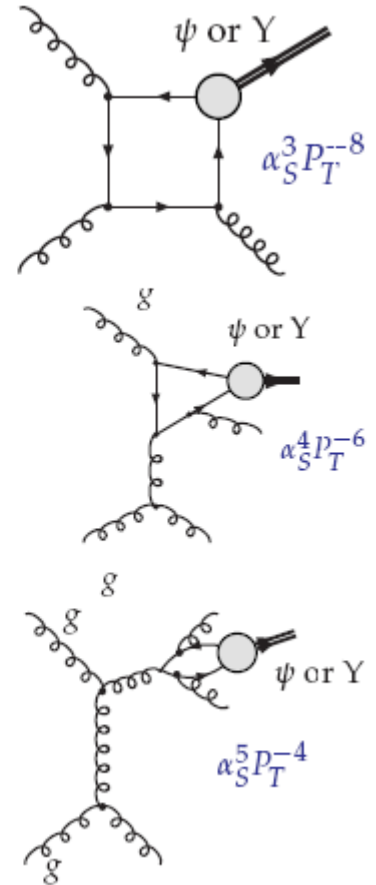
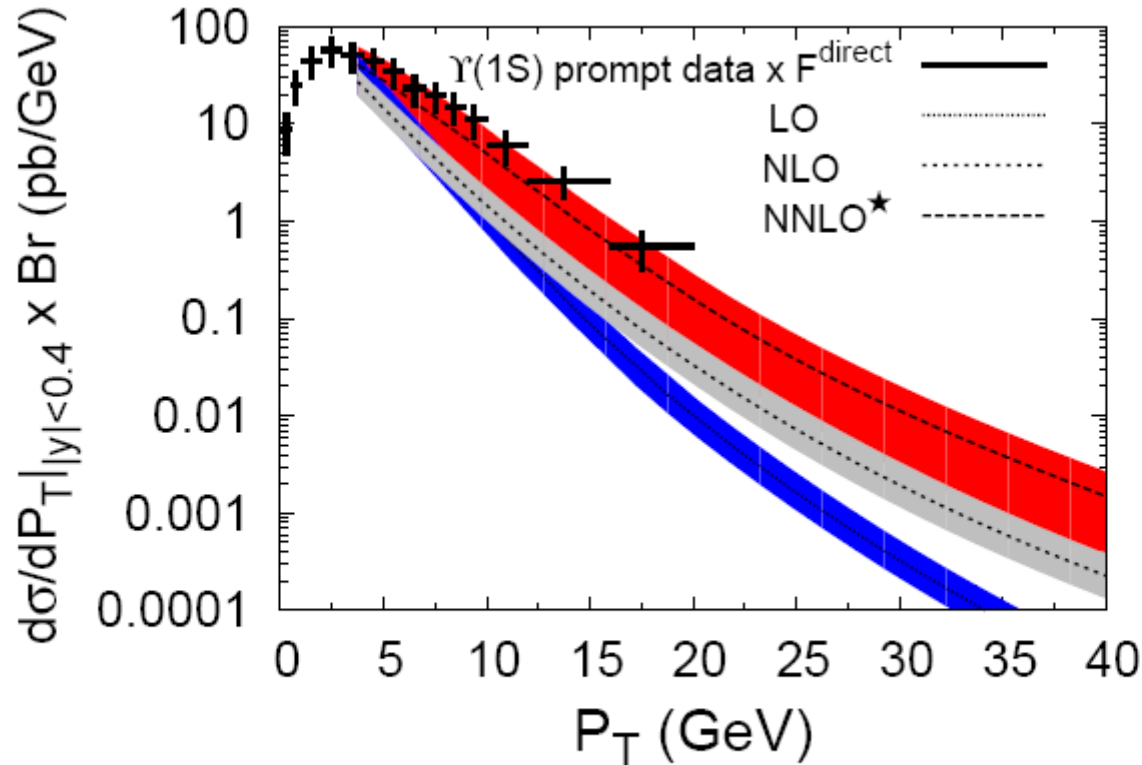
Remainder: QCD corrections for Υ in p+p @ Tevatron

P. Artoisenet, J. Campbell, J.P. Lansberg, F. Maltoni, PRL 101 (2008) 152001.



Remainder: QCD corrections for Υ in p+p @ Tevatron

P. Artoisenet, J. Campbell, J.P. Lansberg, F. Maltoni, PRL 101 (2008) 152001.



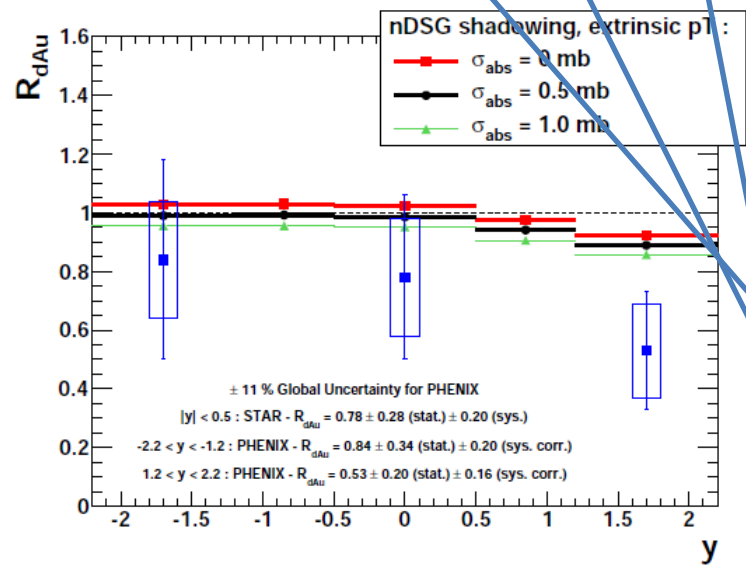
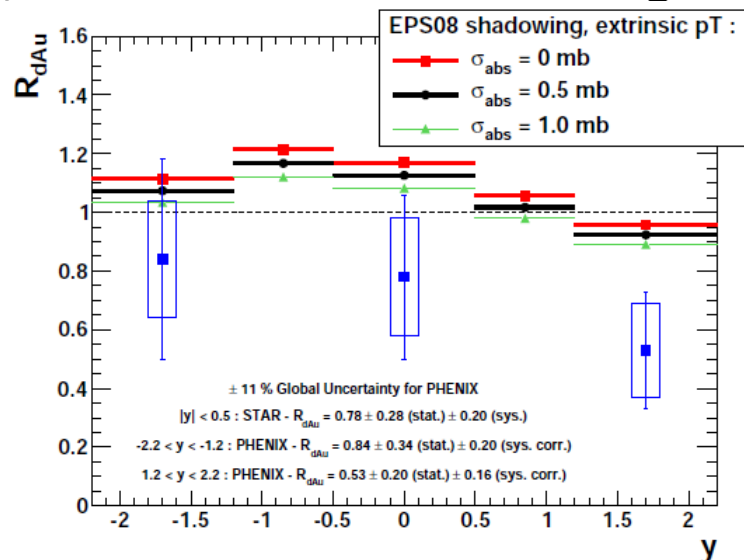
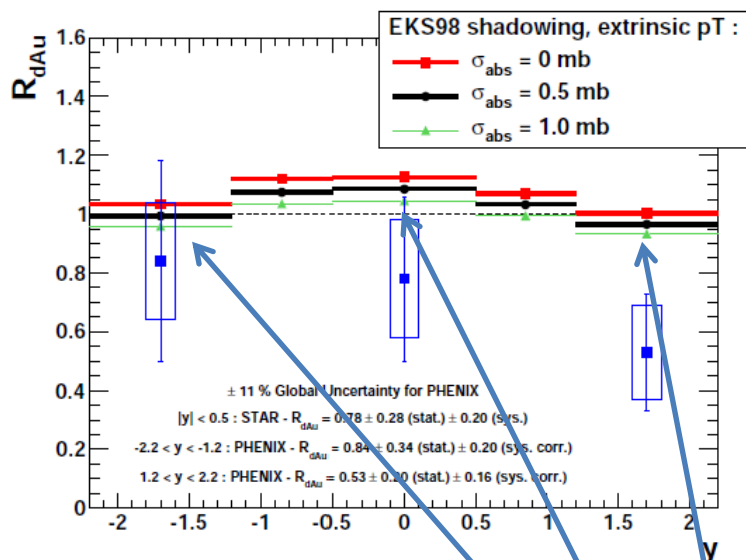
CSM LO sufficient to describe low pT data

Υ vs J/ψ :

- better th. control on the prod. process (larger m_Q)
- smaller uncertainty on σ_{abs} (since small)

Searching new CNM effects: Υ in d+Au @ RHIC

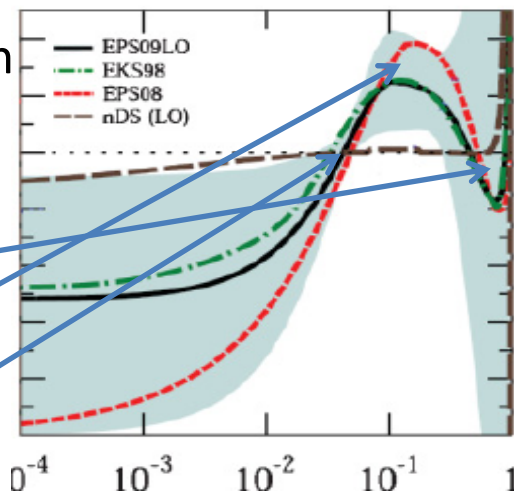
Extrinsic scheme: $\sigma_{abs}=0$ mb, $\sigma_{abs}=0.5$ mb, $\sigma_{abs}=1$ mb in 3 shadowing models



- backward: ok within uncertainties
- central: reasonable job
- forward : clearly too high

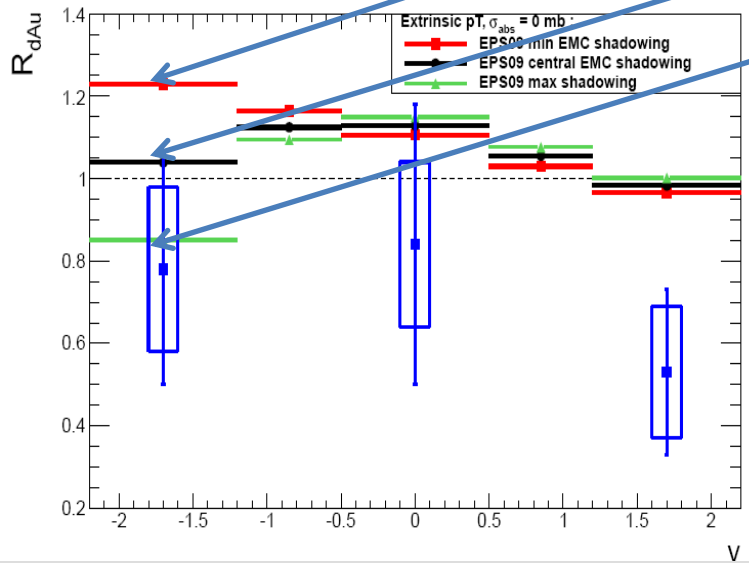
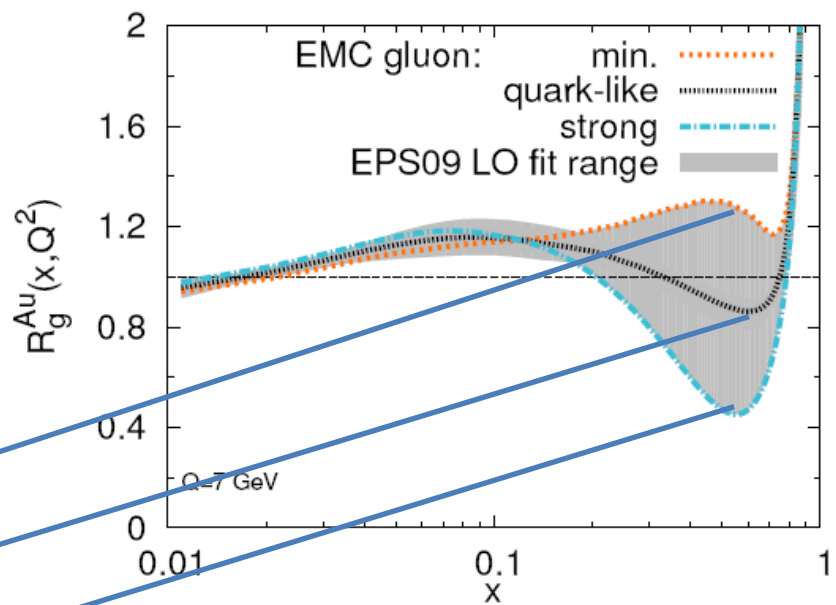
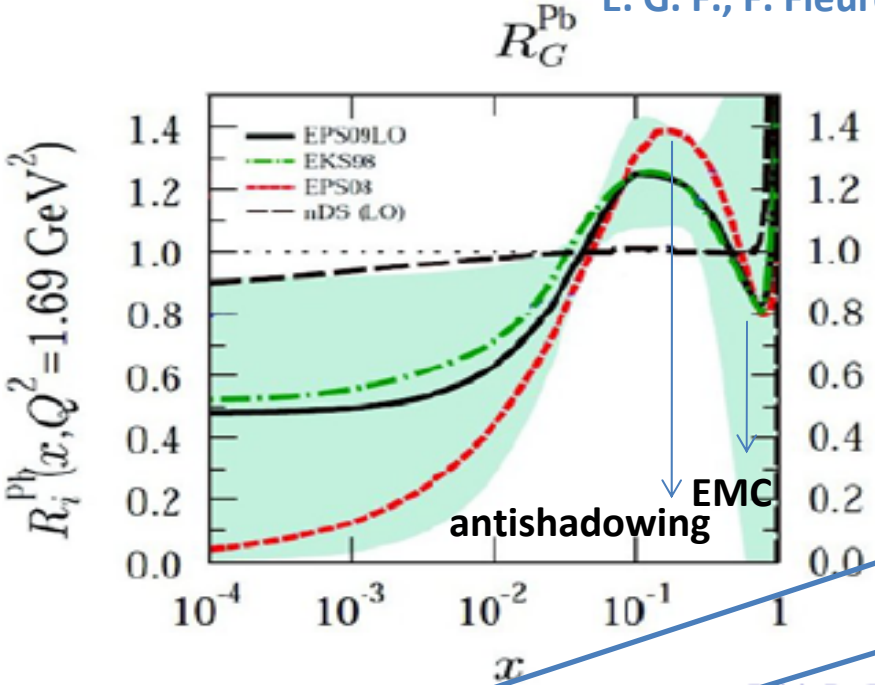
Physical interpretation

- backward: EMC effect
- central: antishadowing
- forward : shadowing ≈ 1
energy loss is needed



Gluon EMC effect

E. G. F., F. Fleuret, J. P. Lansberg and A. Rakotozafindrabe arXiv:1110:504



Let us try to increase the suppression of $g(x)$ in the EMC region, thr shadow incert.:

we have used three of the EPS09 LO sets: one with a quark-like EMC gluon suppression, and the two limiting curves

Works better for backward region
What about the forward region?



Energy loss effect: Radiative energy loss

- **Basic idea:** An energetic parton traveling in a large nuclear medium undergoes multiple elastic scatterings, which induce gluon radiation
=> radiative energy loss (BDMPS)
- **Intuitively:** due to parton energy loss, a hard QCD process probes the incoming PDFs at higher x , where they are suppressed, leading to nuclear suppression
- **The problem:** This energy loss is subject to the LPM bound
=> ΔE is limited and does not scale with E (Brodsky-Hoyer)
- **At RHIC and LHC** (contrary to SPS), typical partons (for $x_1 \sim 10^{-2}$) have energies of the order of hundreds of GeV in the nucleus rest frame
=> radiative energy loss has a negligible effect on the parton x_1



Energy loss effect: Fractional energy loss

- Still, in order to explain large x_F data at RHIC, it would be useful to have
=> a fractional energy loss: $\Delta E \propto E$
(Old idea by Gavin Milana, thought to be ruled out by LPM bound)
- Recently (Arleo, Peigner, Sami arxiv:10006.0818) it has been probed that the **notion of radiated energy associated to a hard process is more general than the notion of parton energy loss.**

The medium-induced gluon radiation associated to large- x_F quarkonium hadroproduction:

- ❖ arises from large gluon formation times $t_f \gg L$
- ❖ scales as the incoming parton energy E
- ❖ cannot be identified with the usual energy loss
- ❖ qualitatively similar to Bethe-Heitler energy loss
- ❖ the Brodsky-Hoyer bound does not apply for large formation times

Thus, the Gavin-Milana assumption of an “energy loss” scaling as E turns out to be qualitatively valid for quarkonium production provided this “energy loss” is correctly interpreted as the radiated energy associated to the hard process, and not as the energy loss of independent incoming and outgoing color charges.

Coherent energy loss effect

(i) **Incoherent** radiation in the initial/final state

Arleo, Peigner, Sami arxiv:1006.0818
Arleo HP2012

Radiation of gluons with large formation times **cancels out** in the **induced** gluon spectrum, leading to $t_f \sim L$

$$\Delta E \propto \hat{q}L^2$$

- Hadron production in nuclear DIS and Drell-Yan in p A collisions
- Jets and hadrons produced in hadronic collisions at large angle

(ii) **Coherent** radiation (interference) in the initial/final state

Induced gluon spectrum dominated by large formation times

$$\Delta E \propto \frac{\sqrt{\hat{q}L}}{M} E$$

- Production of light and open heavy-flavour hadrons at forward rapidities in the medium rest frame (nuclear matter or QGP)
- Production of heavy-quarkonium if color neutralisation occurs on long time-scales $t_{\text{octet}} \gg t_{\text{hard}}$

Energy loss effect @ RHIC

$$\Delta E/E = \Delta x_1/x_1 \simeq N_c \alpha_s \sqrt{\Delta \langle p_T^2 \rangle} / M_T$$

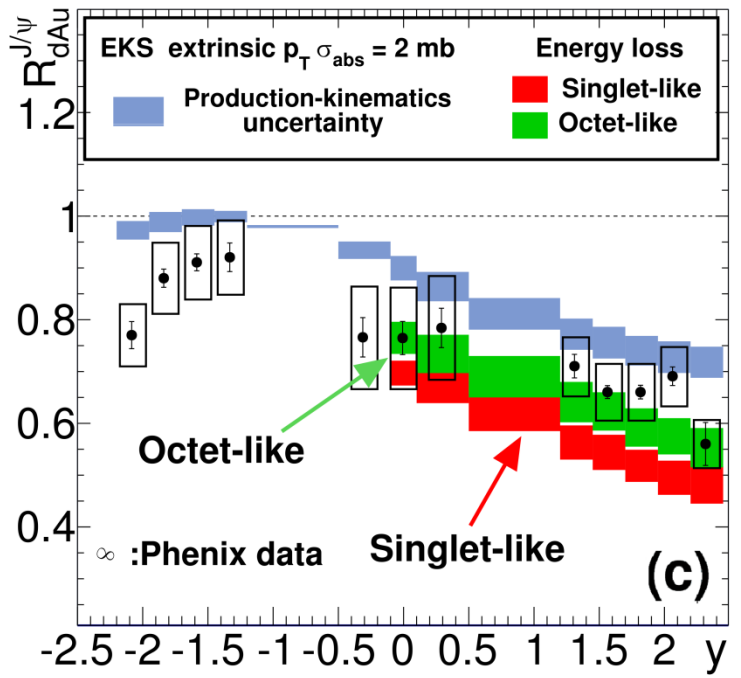
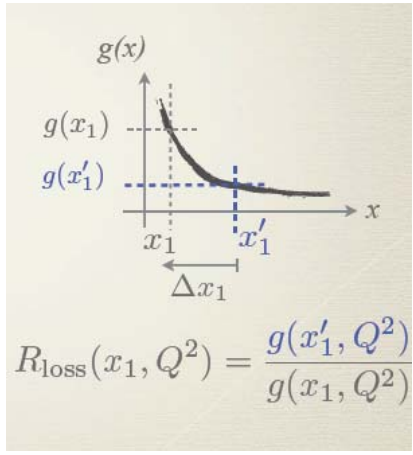
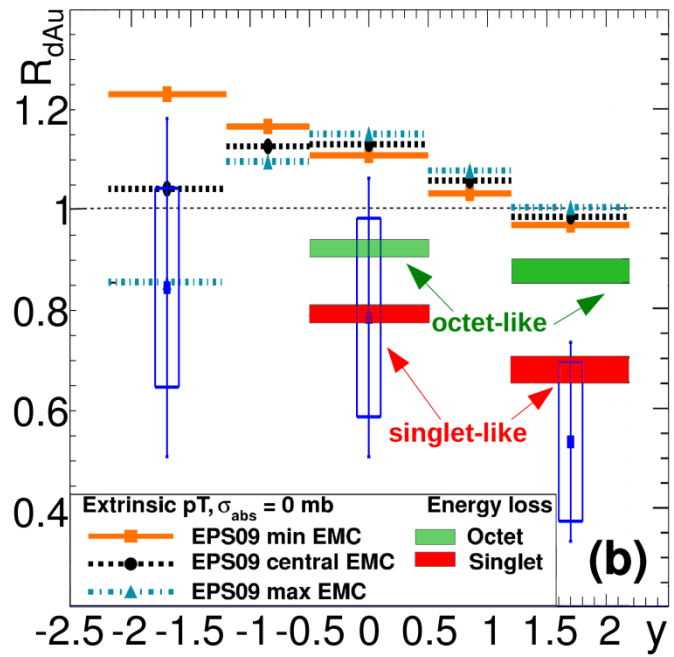
E. G. F., F. Fleuret, J. P. Lansberg and A. Rakotozafindrabe arXiv:1110:5047

Independently of the gluon PDF parameterization, this energy loss will induce

For Υ
 a suppression of 80% - 90% at mid y
 a suppression of 65% -85%at forward y

For J/ψ
 a suppression of 80% - 90% at mid y
 a suppression of 70% -80%at forward y

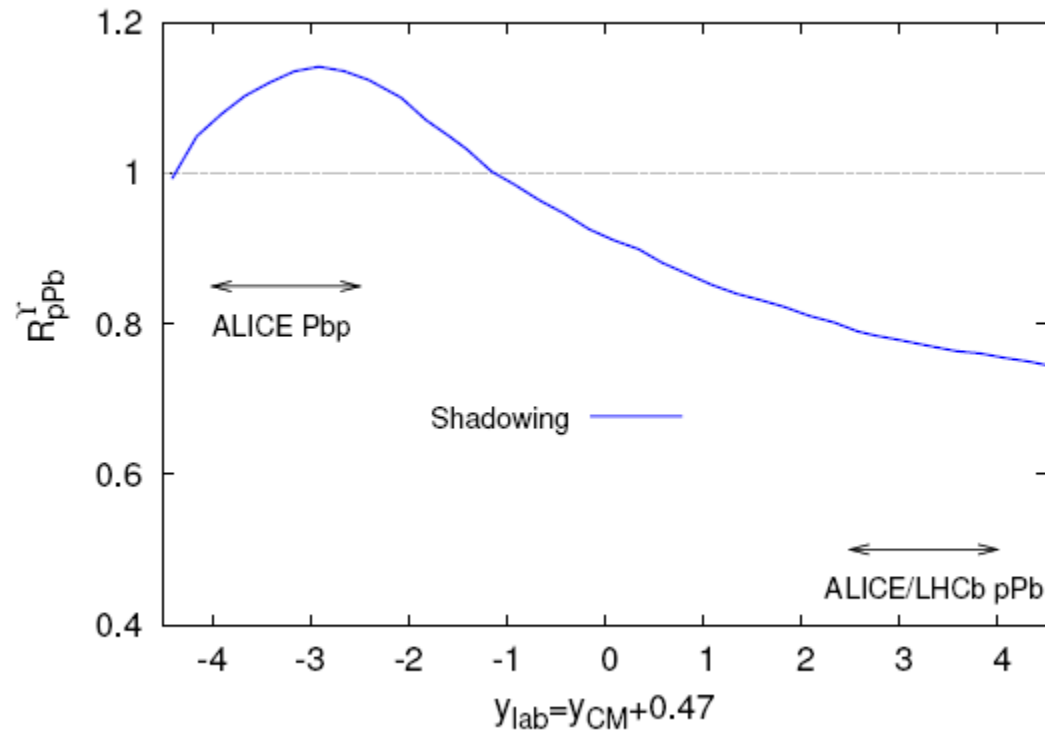
Due to t_f of the order of nuclear size, this energy loss is not applicable in the backward y



Work in progress: Υ production @ LHC

- **Shadowing** effect is not small

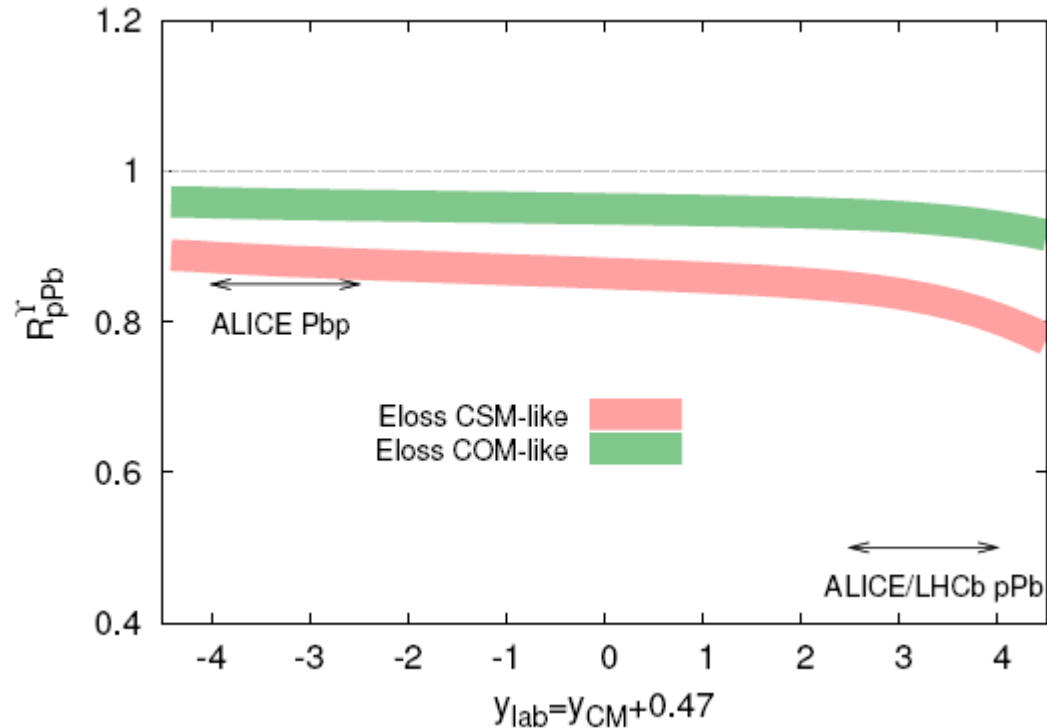
Remember that in PbPb collisions at $y = 0$ shadowing effect is squared compared to pPb



Work in progress: Υ production @ LHC

- **Shadowing** effect is not small

Remember that in PbPb collisions at $y = 0$ shadowing effect is squared compared to pPb

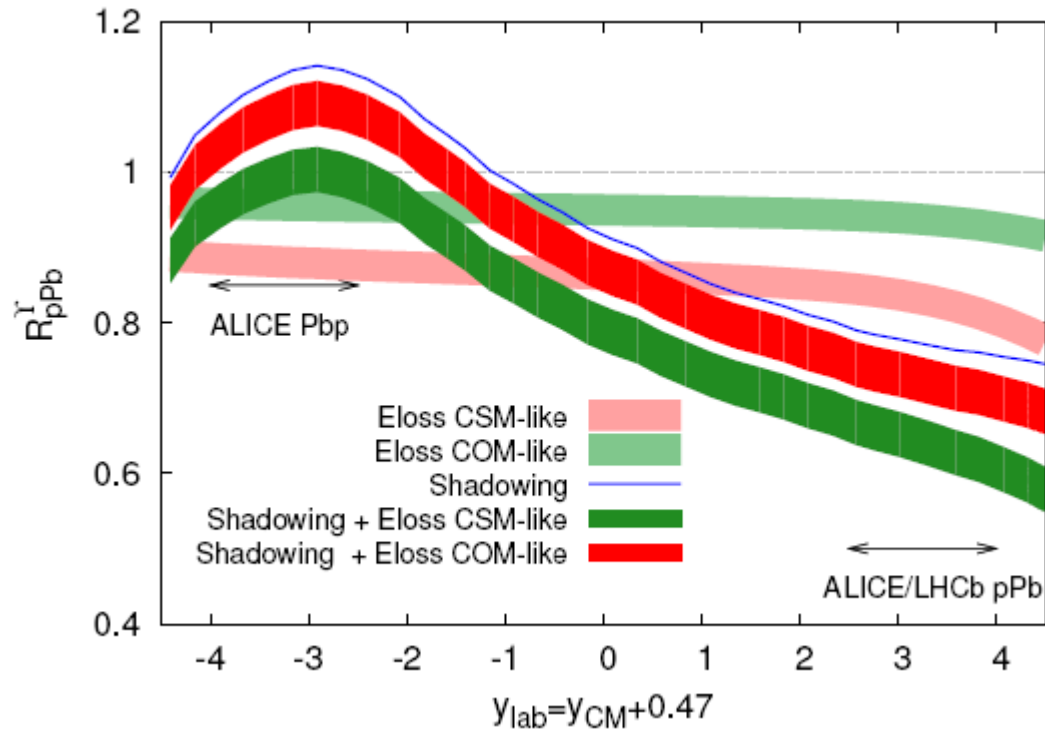


- **Energy loss** also likely matters

Work in progress: Υ production @ LHC

- **Shadowing** effect is not small

Remember that in PbPb collisions at $y = 0$ shadowing effect is squared compared to pPb



- **Energy loss** also likely matters
- Overall, **nuclear matter effects** are not small

they should be accounted when analysing PbPb data

Conclusions

- We have studied the influence of specific partonic kinematics within **2 schemes**: intrinsic ($2 \rightarrow 1$) and extrinsic ($2 \rightarrow 2$) p_T for **different shadowings** including **nuclear absorption** and different partonic models

- for J/ψ
d+Au collisions @ RHIC
A+A collisions @ RHIC and LHC
CNM strong effects at both energy

- for Υ in d+Au collisions @ RHIC:
EMC effect in the backward region
fractional energy loss in the central & forward region
- **energy loss** can also affect J/ψ @ RHIC
- work in progress: Υ production @ LHC