

NLO updates of the EKRT model for central AA  
collisions at RHIC and LHC  
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# Outline

- Introduction: "Old" EKRT final-state saturation model
- NLO pQCD updated new EKRT setup
- Hydrodynamical evolution with EKRT initial state
- Results: multiplicities and identified hadron  $p_T$  spectra for RHIC & LHC
- Summary and Outlook

# The old EKRT model

EKRT = Eskola, Kajantie, Ruuskanen, Tuominen Nucl. Phys. B750 (2000) 379

- model combines pQCD minijet production with the saturation of gluons and (ideal) hydrodynamics
- one can Compute the initial conditions for hydrodynamical evolution from pQCD
- predictions for multiplicity scaling  $N_{ch} \propto A^{0.92}(\sqrt{s})^{0.40}$

# EKRT pQCD + Saturation + Hydrodynamics

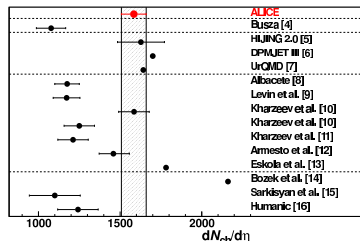
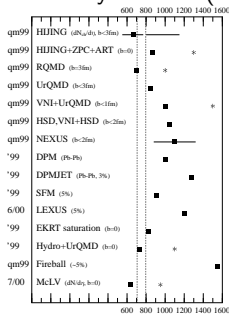
- Low- $p_T$  parton production is controlled by saturation among the produced gluons
- Saturation; Based on the geometric estimate, saturation sets in when produced gluons with  $p_T > p_0$  and transverse area  $\pi/p_0^2$  fill the whole transverse overlap area  $\pi R_A^2$  of the colliding nuclei

$$\Rightarrow K_{QCD} \underbrace{N_{AA}(p_0, \sqrt{s}, \Delta y, \mathbf{b} = \mathbf{0})}_{\text{\# of produced gluons with } p_T > p_0} \frac{\pi}{p_0^2} = K_{sat}(\pi R_A^2)$$

- Solution gives saturation scale  $p_0 = p_{sat}$  for any central ( $\mathbf{b} = \mathbf{0}$ ) AA collision. ( $K_{sat} \propto 1$  uncertainty constant)
- If  $p_{sat} \gg \Lambda_{QCD}$ , pQCD particles (q,g)'s with  $p_T > p_{sat}$  can give a good estimate of the # of partons and transverse energy  $E_T$  produced to midrapidity interval  $\Delta y$

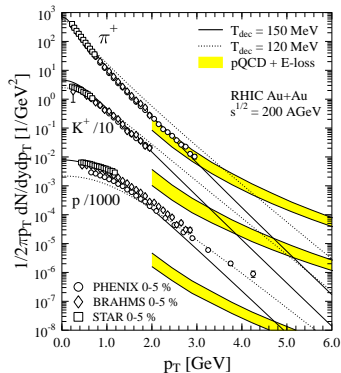
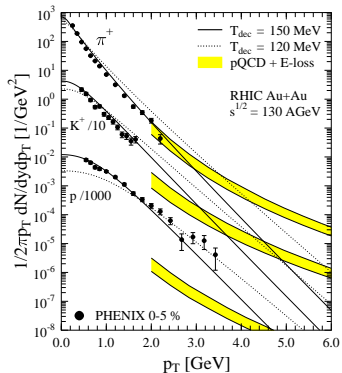
## Old EKRT model results: pQCD + Saturation + Hydro in good agreement with data

- Multiplicity predictions for RHIC  $\sqrt{s} = 200$  GeV and LHC  $\sqrt{s} = 2760, 5500$  GeV (10 years before LHC data! [ref. Eskola et al Nucl.Phys. A696 (2001) ] )



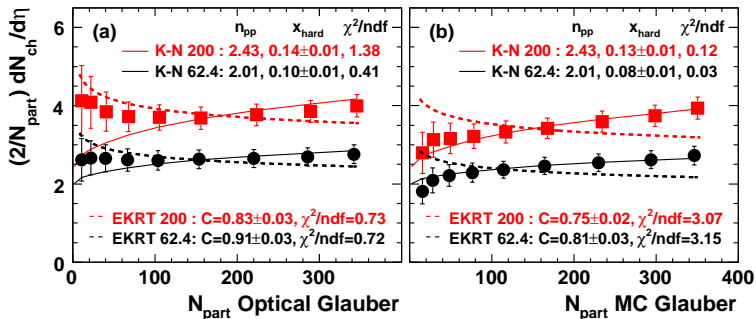
- ALICE LHC figure [ref. PRL 105 (2010) 252301]

- the  $p_T$  spectra of charged particles for 5% most central AuAu collisions at RHIC



- [ref. Eskola et al Phys.Rev. C72 (2005)]

■ Charged-particle multiplicity vs participants



- EKRT model (which uses optical Glauber) [ref. Eskola et al Phys.Lett. B497 (2001)] agrees with the data if  $N_{part}^{exp}$  is estimated with an optical calculation!! [ref. STAR figure Phys.Rev. C79 (2009) 034909]

## EKRT pQCD Modeling Problems

Saturation:

$$K_{QCD} \underbrace{N_{AA}(p_0, \sqrt{s}, \Delta y, \mathbf{b} = \mathbf{0})}_{\text{\#of produced gluons with } p_T > p_0} \frac{\pi}{p_0^2} = K_{sat}(\pi R_A^2)$$

$$dN_{AA} = \int d^2s \sum_{ij} T_A(s) f_i^A(Q^2) \otimes T_A(s) f_j^A(Q^2) \otimes d\hat{\sigma}_{ij}^{pQCD}$$

- Extension to NLO pQCD?  $K_{QCD} = NLO/LO$  rigorously defined only for  $E_T$ , which is infrared and collinear safe (ICS) quantity!!
- $\Delta y$  and  $\alpha_s$  dependence in the saturation criterion ?
- Before now, insufficient control over the uncertainties of the NLO nPDFs (NLO evolution of gluon shadowing, etc..)

Can we fix these problems ?? YES we can!!



# NLO pQCD updated new EKRT setup

New EKRT saturation

$$\frac{dE_T}{d^2sdy}(2 \rightarrow 2) \sim \frac{dE_T}{d^2sdy}(3 \rightarrow 2)$$

$$\Rightarrow (T_{AGA})^2 \frac{\alpha_s^2}{p_0} \sim (T_{AGA})^3 \left(\frac{\alpha_s}{p_0}\right)^3 \Rightarrow T_{AGA} \sim \frac{p_{sat}^2}{\alpha_s} \Rightarrow \frac{dE_T}{dy} \sim R_A^2 p_{sat}^3$$

Thus a new saturation criterion for  $E_T$  in a region  $\Delta y = 1$

$$E_T(p_0, \sqrt{s}, \Delta y = 1, \mathbf{b} = \mathbf{0}) = T_{AA}(\mathbf{0}) \sigma_{QCD} \langle E_T \rangle_{\Delta y, p_0} = K_{sat} (R_A^2 p_0^3)$$

- No explicit  $\alpha_s$  appears!
- Standard nuclear overlap function  $T_{AA}(\mathbf{0})$  accounts for the nuclear collision geometry (Woods-Saxon profile).
- $\sigma_{QCD} \langle E_T \rangle_{\Delta y, p_0}$  is the first moment of the minijet  $E_T$  distribution in  $NN$ .
- We perform a rigorous NLO pQCD computation of  $\sigma_{QCD} \langle E_T \rangle$ : no  $K_{QCD}$  factors anymore!

$$\sigma_{QCD}\langle E_T \rangle_{p_0, \Delta y} = \sigma_{QCD}\langle E_T \rangle_{p_0, \Delta y}^{2 \rightarrow 2} + \sigma_{QCD}\langle E_T \rangle_{p_0, \Delta y}^{2 \rightarrow 3}$$

where

$$\sigma_{QCD}\langle E_T \rangle_{p_0, \Delta y}^{2 \rightarrow 2} = \frac{1}{2!} \int [DPS]_2 \frac{d\sigma^{2 \rightarrow 2}}{[DPS]_2} \tilde{S}_2(p_1, p_2)$$

$$\sigma_{QCD}\langle E_T \rangle_{p_0, \Delta y}^{2 \rightarrow 3} = \frac{1}{3!} \int [DPS]_3 \frac{d\sigma^{2 \rightarrow 3}}{[DPS]_3} \tilde{S}_3(p_1, p_2, p_3)$$

- Partonic  $2 \rightarrow 2$  ( $gg \rightarrow gg$ , etc...) processes for LO & NLO corrections (1-loop level)
- $2 \rightarrow 3$  ( $gg \rightarrow ggg$ , etc...) processes - only NLO corrections
- UV renormalized  $|M|^2$  in  $4 - 2\epsilon$  dimensions (R.K Ellis at all)
- IR/CL divergencies handled with NLO def. of PDFs & EKS subtraction method

The measurement functions  $\tilde{S}_2$  and  $\tilde{S}_3$  fulfil the IR/CL criteria, which ensure that  $\sigma_{QCD}\langle E_T \rangle$  is a well defined IR/CL safe quantity

$$\tilde{S}_2 = \left[ \epsilon(y_1) + \epsilon(y_2) \right] p_{T2} \Theta(p_{T2} \geq p_0)$$

$$\tilde{S}_3 = E_T \Theta(p_{T1} + p_{T2} + p_{T3} \geq 2p_0) \Theta(E_T \geq C \times p_0)$$

where  $\epsilon(y_i) = 1$  if  $y_i \in \Delta y$  otherwise  $\epsilon(y_i) = 0$

- $\tilde{S}_3 \rightarrow \tilde{S}_2$  at IR/CL limits
- We introduce a new set of measurement functions,  $0 \leq C \leq 1$  which control the amount of  $E_T$  in  $\Delta y$  carried by the partons
- New: any  $C$  between 0&1 is equally good and IR/CL safe!

Bound proton PDFs  $f_{i/A}(x, Q^2)$  for each parton flavor  $i$

$$f_{i/A}(x, Q^2) \equiv R_i^A(x, Q^2) f_i^P(x, Q^2)$$

$R_i^A$  denotes the nuclear modification to the free proton PDF  $f_i^P$

- Old EKRT setup
  - GRV94 LO parton densities with nuclear effects from the EKS98 [ref. Eskola et al Eur.Phys.J. C9 (1999)] LO parametrization
  - $K_{QCD}$  NLO factor - GRV94 and CTEQ5 (LO& NLO) and LO EKS98: NO full NLO evolution
- New EKRT setup
  - CTEQ6 NLO parton densities with EPS09 [ref. JHEP 0904 (2009) 065] NLO parametrization
  - Study also the propagation of nPDF uncertainties with the 30 error sets in EPS09

## Hydrodynamical evolution; Based on the 1+1 ideal hydrodynamics (H. Holopainen)

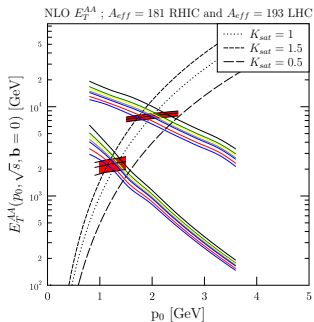
- Impact parameter  $\mathbf{b} = \mathbf{0}$  fm
- Boost invariance in  $z$ -direction
- Initial time  $\tau_0 = 1/p_{sat}$
- Initial  $\epsilon(\tau_0)$  from  $\{E_T(p_{sat}), \tau_0\}$  and eBC/eWN profiles
- EoS: s95p-PCE from P. Huovinen and P. Petreczky [ref. Nucl. Phys. A837 (2010)]
- Freeze-out temperature  $T_f = 120$  MeV

⇒ Calculate particle multiplicities,  $p_T$  spectra for central Au+Au  $\sqrt{s} = 200$  GeV at RHIC, and Pb+Pb  $\sqrt{s} = 2760$  GeV at LHC (see Results)

## New EKRT model Results (preliminary)

Study the  $E_T \leftrightarrow N_{ch}$  systematics of the new EKRT model from RHIC to LHC (central collisions):

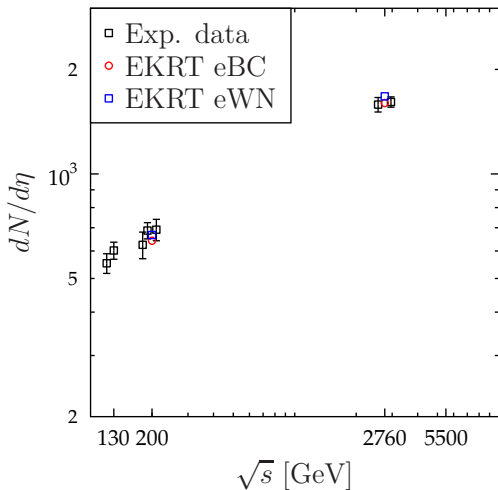
- effects of choosing  $C$
- effects of  $K_{sat}$
- map with Hydro (red bands): Hydro curves calculated by fixing multiplicity to the LHC and RHIC measurements



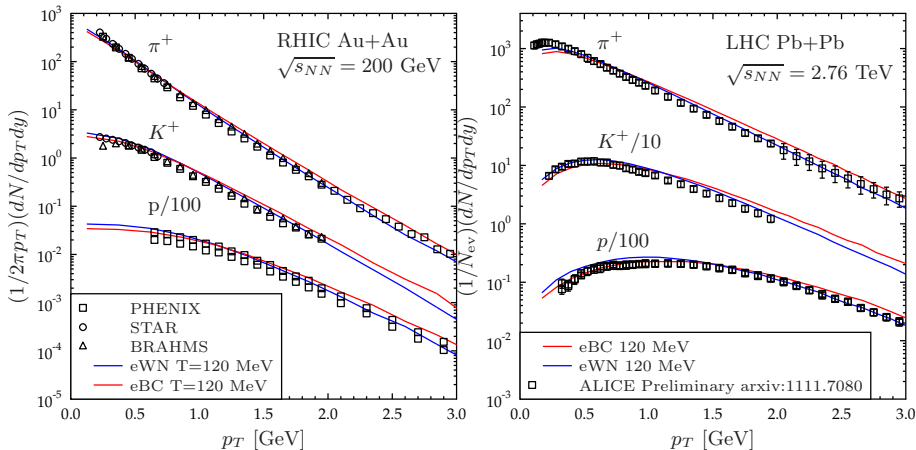
$E_T$  saturation for  $C = 0.75$

$\sqrt{s}$ GeV	$p_{sat}$	$\tau_0$
RHIC 200	1.31 GeV	0.15 fm
LHC 2760	1.97 GeV	0.10 fm

- Good agreement with the multiplicity data with  $C = 0.75$ , both for LHC and RHIC!



### $p_T$ spectra (ALICE preliminary data ref. arxiv:1111.7080)





## Improvements over the old EKRT

- new saturation criterion for IR/CL safe  $E_T$ , no  $K_{QCD}$  needed
- NLO nPDFs (CTEQ6 & EPS09)
- rigorous NLO pQCD computation for  $E_T$
- new measurement functions with  $0 \leq C \leq 1$

RHIC-LHC systematics of the new EKRT setup looks good!

## Next

- Centrality dependence: local saturation with impact parameter dependent nPDFs (EPS09s [ref. arXiv:1205.5359])
- Viscous hydrodynamics